

ECE330: Power Circuits & Electromechanics
Lecture 14. Co-Energy via line integrals

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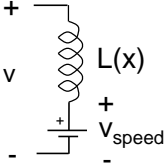
Schedule

- Mon 10/7: Voltage of mechanical origin
- Wed 10/9: Force of electrical origin
- Friday 10/11: Review + Quiz 5 (Mech. volt.)
- Mon 10/14: Energy via line integrals
- Wed 10/16: Co-energy via line integrals
- Friday 10/18: Review + Quiz 6 (Elec. force)

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Linear x-varying inductors

Inductance
 $\lambda(i, x) = L(x)i, \quad L(x) = \frac{N^2}{\mathcal{R}(x)}$



Voltage of mechanical origin
 $v = L(x)\frac{di}{dt} + [L'(x)i]\frac{dx}{dt}$

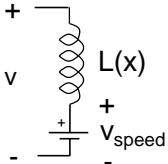
Force of electric origin (direction of increasing x)
 $f^e = +\frac{\partial}{\partial x} [\text{Co-energy}] = -\frac{\partial}{\partial x} [\text{Energy}]$

Energy & Co-energy (This week)
 Co-energy = $\frac{1}{2}L(x)i^2$
 Energy = $\frac{1}{2}L(x)^{-1}\lambda^2$

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Conservative electromagnetic systems

Flux & current relations
 $\lambda = \lambda(i, x), \quad i = i(\lambda, x)$



Voltage of mechanical origin
 $v = \underbrace{\frac{\partial \lambda}{\partial i}}_{L(x)} \frac{di}{dt} + \underbrace{\frac{\partial \lambda}{\partial x}}_{v_{\text{speed}}} \frac{dx}{dt}$

Force of electric origin (direction of increasing x)
 $f^e = +\frac{\partial}{\partial x} [\text{Co-energy}] = -\frac{\partial}{\partial x} [\text{Energy}]$

Energy & Co-energy (Today)
 Co-energy = $\int_0^i \lambda(i', x) di'$
 Energy = $\int_0^\lambda i(\lambda', x) d\lambda'$

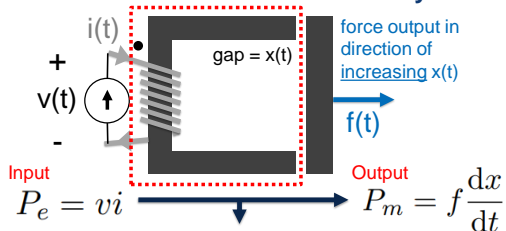
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Today

- (Review) Line integral definition of energy
- Line integral definition of co-energy
- Example: Forces from co-energy
- Example: Co-energy from flux
- Example: Co-energy from multi-port flux

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Conservative electromechanical system



Input: $P_e = vi$ Output: $P_m = f \frac{dx}{dt}$

Storage: $\frac{dE}{dt} = i \frac{d\lambda}{dt} - f \frac{dx}{dt}$

$\Delta E = i\Delta\lambda - f\Delta x$ Energy over Δt

Closed system = Conservative

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Line integral view of stored energy

$$\Delta E = i\Delta\lambda - f\Delta x$$

$$= (i, -f) \cdot (\Delta\lambda, \Delta x)$$

$$= \underline{F} \cdot \Delta \underline{r}$$

Position $\underline{r} = (\lambda, x)$
 Force $\underline{F} = (i, -f)$

Flux as spatial dim. Current as force.
 etymology for "mmf"

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Conservatism implies path independence

flux displacement

Integrate along any path \rightarrow same energy

$$\text{Energy}(\lambda_1, x_1) - \text{Energy}(\lambda_0, x_0) = \int_C \underline{F}(\underline{r}) \cdot d\underline{r} = \int_C i(\lambda, x) d\lambda + \int_C -f(\lambda, x) dx$$

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Line integral definition of energy

$$\text{Energy}(\lambda_1, x_1) - \text{Energy}(\lambda_0, x_0) = \int_C \underline{F}(\underline{r}) \cdot d\underline{r} = \int_C i(\lambda, x) d\lambda + \int_C -f(\lambda, x) dx$$

Energy From Electrical Energy From Mechanical

Path independent (via conservation) Δ Path dependent

Energy(0, x) = 0 (Can be any constant)

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Computing energy from current

$$\text{Energy}(\lambda, x) - \text{Energy}(0, 0) \stackrel{0 \text{ by def}}{=} \int_C i(\lambda', x') d\lambda' + \int_C -f(\lambda', x') dx'$$

Special integration path avoids the need to know the force

Draw magnetic circuit, compute reluctance & flux, compute flux linkage solve for current

$$i(\lambda, x) = \left(\frac{g + x}{L_0} \right) \lambda$$

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Computing energy from current

$$\text{Energy}(\lambda, x) - \text{Energy}(0, 0) \stackrel{0 \text{ by def}}{=} \int_C i(\lambda', x') d\lambda' + \int_C -f(\lambda', x') dx'$$

no change in flux no flux no force

no change in x

$$i(\lambda, x) = L(x)^{-1} \lambda$$

Path A: Move x into position with no flux
 Path B: Charge the inductor with fixed x

$$= \int_0^\lambda L(x)^{-1} \lambda' d\lambda' = L(x)^{-1} \left[\frac{(\lambda')^2}{2} \right]_0^\lambda = \frac{1}{2} L(x)^{-1} \lambda^2$$

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Energy in a conservative magnetic system

Theorem. If $L = \partial\lambda/\partial i > 0$, then

$$\text{Energy}(\lambda, x) = \int_0^\lambda i(\lambda', x) d\lambda' \quad [\text{J}]$$

Corollary. If $\lambda = L(x)i$ and $L(x) > 0$, then

$$\text{Energy}(\lambda, x) = \frac{1}{2} L(x)^{-1} \lambda^2 \quad [\text{J}]$$

Proof. (Graphical)

By definition, energy is the area of the triangle

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Force of electric origin in a conservative magnetic system

Theorem. If energy conserving, then force in the direction of increasing x is

$$f^e = -\frac{\partial}{\partial x} [\text{Energy}]$$

and the current into the dot is

$$i = \frac{\partial}{\partial \lambda} [\text{Energy}] \quad \text{By viewing current as force}$$

Proof. By the gradient theorem

$$\underline{F}(\underline{r}) = (i, -f) = \nabla \text{Energy}(\underline{r})$$

$$\underline{r} = (\lambda, x)$$

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- (Review) Line integral definition of energy
- Line integral definition of co-energy
- Example: Forces from co-energy
- Example: Co-energy from flux
- Example: Co-energy from multi-port flux

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Main issue with energy computations

Energy(λ, x) - Energy(0, 0) by def

$$= \int_C i(\lambda', x') d\lambda' + \int_C -f(\lambda', x') dx'$$

no change in flux no flux no force

$$= \left(\int_A i(\lambda', x') d\lambda' + \int_A -f(\lambda', x') dx' \right)$$

no change in x

$$+ \left(\int_B i(\lambda', x') d\lambda' + \int_B -f(\lambda', x') dx' \right)$$

Unnatural to express current as function of flux

Can require inverting a nonlinear function

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Definition of co-energy

Co-energy(i, x) = λi - Energy(λ, x)

$$\frac{dE_{co}}{dt} = \lambda \frac{di}{dt} + i \frac{d\lambda}{dt} - \frac{dE}{dt}$$

Swaps the roles of current and flux positive force

$$= \lambda \frac{di}{dt} + i \frac{d\lambda}{dt} - \left(i \frac{d\lambda}{dt} - f \frac{dx}{dt} \right) = \lambda \frac{di}{dt} + f \frac{dx}{dt}$$

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Line integral view of co-energy

$$\frac{dE_{co}}{dt} = \lambda \frac{di}{dt} + f \frac{dx}{dt}$$

$$\Delta E_{co} = \lambda \Delta i + f \Delta x$$

$$= (\lambda, f) \cdot (\Delta i, \Delta x)$$

$$= \underline{F} \cdot \Delta \underline{r}$$

Position $\underline{r} = (i, x)$

Force $\underline{F} = (\lambda, f)$

Current as spatial dim. Flux as force.

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Co-energy inherits conserv. of energy

current displacement

Integrate along any path → same co-energy

$$\text{Co-energy}(i_1, x_1) - \text{Co-energy}(i_0, x_0) = \int_C \underline{F} \cdot d\underline{r} = \int_C \lambda(i, x) di + \int_C f(i, x) dx$$

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Line integral definition of co-energy

$$\text{Co-energy}(i_1, x_1) - \text{Co-energy}(i_0, x_0) = \int_C \underline{F} \cdot d\underline{r} = \int_C \lambda(i, x) di + \int_C f(i, x) dx$$

$\int_C \lambda(i, x) di$ is **Path independent (via conservation)**
 $\int_C f(i, x) dx$ is **Path dependent**
 Δ **Not necessarily physical**

Co-energy(0, x) = 0 (Can be any constant)
 Co-energy(i, x) = λi - Energy(λ, x)

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Computing co-energy from flux

$$\begin{aligned} \text{Co-energy}(i, x) - \text{Co-energy}(0, 0) &= \int_C \lambda(i', x') di' + \int_C f(i', x') dx' \\ &= \left(\int_A \lambda(i', x') di' + \int_A f(i', x') dx' \right) + \left(\int_B \lambda(i', x') di' + \int_B f(i', x') dx' \right) \\ &= \int_0^i L(x) i' di' = \left[L(x) \frac{(i')^2}{2} \right]_0^i = \frac{1}{2} L(x) i^2 \end{aligned}$$

$\lambda(i, x) = L(x) i$
 $\Gamma = (i, x)$
 $\underline{E} = (\lambda, f)$
 $i = 0$
 $E_{co}(0, x) = 0$

Path A: Move x into position with no current
 Path B: Charge the inductor with fixed x

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Co-energy in a conservative magnetic system

Theorem. If $L = \partial N / \partial i > 0$, then

$$\text{Co-energy}(i, x) = \int_0^i \lambda(i', x) di' \text{ [J]}$$

Corollary. If $\lambda = L(x)i$ and $L(x) > 0$, then

$$\text{Co-energy}(i, x) = \frac{1}{2} L(x) i^2 \text{ [J]}$$

Proof. (Graphical)

By definition, co-energy is the area of the triangle

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Graphical relationship

$$\begin{aligned} \text{Energy}(\lambda, x) &= \int_0^\lambda i(\lambda', x) d\lambda' \\ \text{Co-energy}(i, x) &= \int_0^i \lambda(i', x) di' \\ &= \lambda i - \text{Energy}(\lambda, x) \end{aligned}$$

$\lambda = L(x)i$
 Energy = co-energy if and only if linear λ - i relation
 Δ λ vs i (space-force) **not** λ vs x (space-space).

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Force of electric origin in a conservative magnetic system

Theorem. If **energy conserving**, then force in the direction of increasing x is

$$f^e = + \frac{\partial}{\partial x} [\text{Co-energy}]$$

and the current into the dot is

$$\lambda = + \frac{\partial}{\partial i} [\text{Co-energy}]$$

Again, by viewing flux as force

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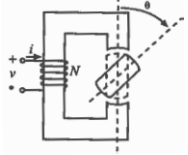
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Warm-up: Forces from co-energy



$W'_m(i, \theta) = \log(1 + i) \cos \theta$

Recall $\lambda = + \frac{\partial}{\partial i}$ [Co-energy]

$f^e = + \frac{\partial}{\partial x}$ [Co-energy] ← Treat θ as x

$\lambda = ?$ [At]

$f^e = ?$ [N]

A) $+ \frac{\cos \theta}{1 + i}$
 B) $-\log(1 + i) \sin \theta$
 C) $+ \frac{\cos \theta}{1 + i} - \log(1 + i) \sin \theta$
 D) $-\log(1 + i) \sin \theta$

A) $+ \log(1 + i) \sin \theta$
 B) $-\log(1 + i) \sin \theta$
 C) $+ \frac{\cos \theta}{1 + i} - \log(1 + i) \sin \theta$
 D) $-\log(1 + i) \sin \theta$

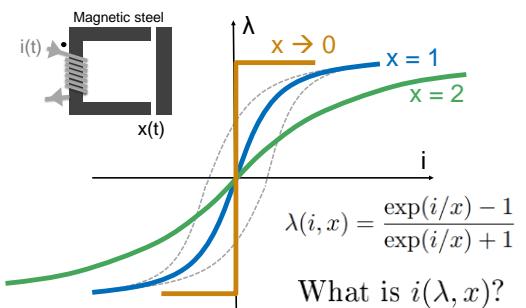
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Realistic inductor (without hysteresis)



$\lambda(i, x) = \frac{\exp(i/x) - 1}{\exp(i/x) + 1}$

What is $i(\lambda, x)$?

How does saturation affect force?

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Computing co-energy from flux

$\lambda(i, x) = \frac{\exp(i/x) - 1}{\exp(i/x) + 1}$ Co-energy(i, x) - Co-energy($0, x$)

$\text{Co-energy}(i, x) = \int_C \lambda(i', x') di' + \int_{e^{-i/x}}^1 f(i', x') dx'$

$\text{Co-energy}(0, x) = \int_C \lambda(i', x') di' + \int_{e^{-i/x}}^1 f(i', x') dx'$

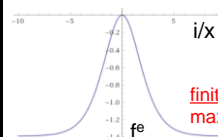
$\text{Co-energy}(i, x) - \text{Co-energy}(0, x) = \int_0^i \frac{\exp(i'/x) - 1}{\exp(i'/x) + 1} di'$

$= 2x \log \left(\frac{\exp(i/x) + 1}{2} \right) - i$

Path C: Charge the inductor with fixed x

$f^e = + \frac{\partial}{\partial x}$ [Co-energy]

$= \frac{-2i/x}{1 + \exp(-i/x)} + 2 \log \left(\frac{\exp(i/x) + 1}{2} \right)$



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Computing co-energy from flux

$\lambda(i, x) = \frac{\exp(i/x) - 1}{\exp(i/x) + 1}$ Co-energy(i, x) - Co-energy($0, x$)

$\text{Co-energy}(i, x) = \int_C \lambda(i', x') di' + \int_{e^{-i/x}}^1 f(i', x') dx'$

$\text{Co-energy}(0, x) = \int_C \lambda(i', x') di' + \int_{e^{-i/x}}^1 f(i', x') dx'$

$\text{Co-energy}(i, x) - \text{Co-energy}(0, x) = \int_0^i \frac{\exp(i'/x) - 1}{\exp(i'/x) + 1} di'$

$= 2x \log \left(\frac{\exp(i/x) + 1}{2} \right) - i$

Path C: Charge the inductor with fixed x

$\frac{\partial}{\partial i}$ [Co-energy] = $2x \cdot \frac{2}{\exp(i/x) + 1} \cdot \frac{\exp(i/x)}{2} \cdot \frac{1}{x} - 1$

$= \frac{2 \exp(i/x)}{\exp(i/x) + 1} - 1 = \lambda(i, x)$

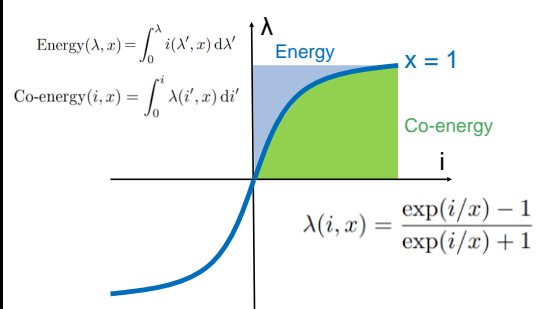
Partial derivative wrt current recovers flux, as expected.

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E and co-E in a realistic inductor

Energy(λ, x) = $\int_0^\lambda i(\lambda', x) d\lambda'$

Co-energy(i, x) = $\int_0^i \lambda(i', x) di'$



$\lambda(i, x) = \frac{\exp(i/x) - 1}{\exp(i/x) + 1}$

Δ In general, Energy \neq Co-energy

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Computing co-energy from multipoint flux

$$\lambda_1 = L_1 i_1 + M(\theta) i_2$$

$$\lambda_2 = M(\theta) i_1 + L_2 i_2$$

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Integration path in two-port system

$\underline{F} = (\lambda_1, \lambda_2, f)$
 $\underline{r} = (i_1, i_2, x)$

$\lambda_1 = L_1 i_1 + M(\theta) i_2$
 $\lambda_2 = M(\theta) i_1 + L_2 i_2$

Co-energy(i_1, i_2, θ)

$$= \int_0^{i_1} \lambda_1(i_1', 0, \theta) di_1' + \int_0^{i_2} \lambda_2(i_1, i_2', \theta) di_2'$$

Charge coil 1 with coil 1 charged Charge coil 2 with coil 1 charged

$$= \int_0^{i_1} [L_1 i_1' + M(\theta)] di_1' + \int_0^{i_2} [M(\theta) i_1 + L_2 i_2'] di_2'$$

$$= L_1 \frac{i_1^2}{2} + M(\theta) i_1 i_2 + L_2 \frac{i_2^2}{2}$$

$$= \frac{1}{2} \begin{bmatrix} i_1 \\ i_2 \end{bmatrix}^T \begin{bmatrix} L_1 & M(\theta) \\ M(\theta) & L_2 \end{bmatrix} \begin{bmatrix} i_1 \\ i_2 \end{bmatrix}$$

B: Charge coil 1 with fixed x and de-energized coil 2
 C: Charge coil 2 with fixed x and energized coil 1

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Co-energy in a conservative magnetic system with multiple ports

Theorem. If $\underline{\underline{L}} = \nabla_{\underline{i}} \lambda > 0$, then always fixed

Co-energy(i_1, i_2, i_3, \dots, x) = $\int_0^{i_1} \lambda_1(i_1', 0, 0, \dots, x) di_1'$
 + $\int_0^{i_2} \lambda_1(i_1, i_2', 0, \dots, x) di_2'$
 + $\int_0^{i_3} \lambda_1(i_1, i_2, i_3', \dots, x) di_3'$
 + ...

Corollary. If $\underline{\underline{L}} = \underline{\underline{L}}(x) \underline{i}$ and $\underline{\underline{L}}(x) \succ 0$, then

Co-energy(\underline{i}, x) = $\frac{1}{2} \underline{i}^T \underline{\underline{L}}(x) \underline{i}$ [J]

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Do homework 6
 For example answers, look at Problems 1 & 2 from previous Exam 2

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