

1. (a) (1 point) The vector wave field $\mathbf{E}(x, t)$ is given by

$$\mathbf{E}(y, t) = 80\Delta\left(\frac{t + y/c}{\tau}\right) \hat{z} \left[\frac{\text{V}}{\text{m}}\right].$$

- (b) (1 point) The associated wave field $\mathbf{H}(x, t)$ is

$$\mathbf{H}(y, t) = -\frac{80}{\eta_o} \Delta\left(\frac{t + y/c}{\tau}\right) \hat{x} \left[\frac{\text{A}}{\text{m}}\right].$$

- (c) (2 points) The Poynting vector is

$$\mathbf{S} = \mathbf{E} \times \mathbf{H} = -\frac{6400}{\eta_o} \Delta^2\left(\frac{t + y/c}{\tau}\right) \hat{y} \left[\frac{\text{W}}{\text{m}^2}\right],$$

and its maximum value is

$$\max(|\mathbf{E} \times \mathbf{H}|) = \frac{6400}{\eta_o} \approx \frac{160}{3\pi} \left[\frac{\text{W}}{\text{m}^2}\right].$$

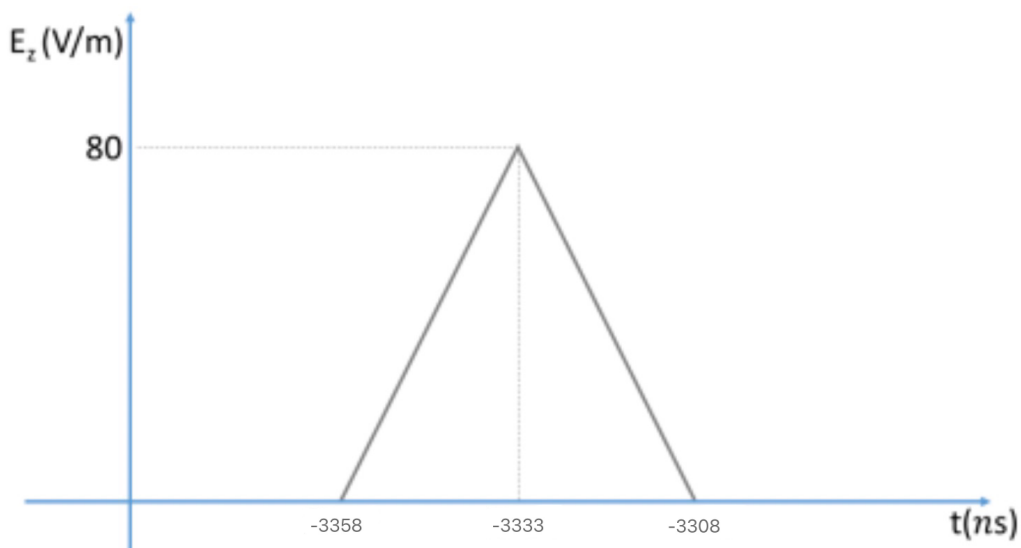
- (d) (2 points) The location of the peak of $\mathbf{E} \times \mathbf{H}$ evolves according to

$$\frac{t + y/c}{\tau} = 0 \quad \rightarrow \quad y = -ct.$$

- (e) (2 points) The field E_z given by

$$E_z(y, t)|_{y=1000\text{m}} = 80\Delta\left(\frac{t + 3.33 \times 10^{-6} \text{s}}{50 \text{ ns}}\right) \left[\frac{\text{V}}{\text{m}}\right],$$

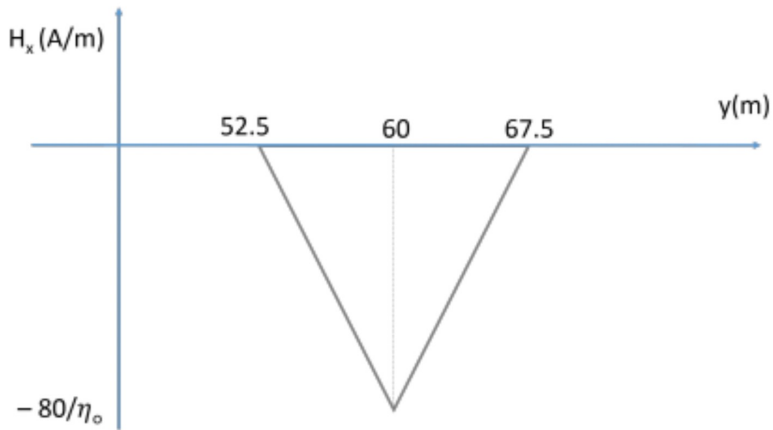
It is plotted in the following figure:



(f) (2 points) The field H_y given by

$$H_x(y, t)|_{t=-200 \text{ ns}} = -\frac{80}{\eta_o} \Delta\left(\frac{-60 + y}{15}\right) \left[\frac{\text{A}}{\text{m}}\right],$$

It is plotted in the following figure:



2. (a) (2 points) The propagation velocity is given by

$$v = \frac{c}{4} = 7.5 \times 10^7 \left[\frac{\text{m}}{\text{s}} \right].$$

- (b) (4 points) If $\epsilon = \epsilon_r \epsilon_o$ and $\mu = \mu_r \mu_o$, then the intrinsic impedance η can take the following form

$$\eta = \sqrt{\frac{\mu}{\epsilon}} = \sqrt{\frac{\mu_r}{\epsilon_r}} \eta_o$$

while the wave propagation velocity v can be written as

$$v = \frac{1}{\sqrt{\epsilon\mu}} = \frac{c}{\sqrt{\epsilon_r\mu_r}}.$$

Using the results of parts (a) and (b) we find that

$$\epsilon_r = 16, \mu_r = 1.$$

- (c) (2 points) Since $E_x = \eta H_y$ we have that

$$u\left(t - \frac{z}{c/4}\right) = \eta H_{y_o} u\left(t - \frac{4z}{c}\right)$$

with $\eta = \frac{1}{4}\eta_o$, then $H_{y_o} = \frac{1}{30\pi}$

- (d) (2 points) Finally, since $E_y = -\eta H_x$ (recall that the propagation direction is $\hat{z} = -\hat{y} \times \hat{x}$, hence the minus sign), we have that

$$g\left(t - \frac{z}{c/4}\right) = -10 \times \left(\frac{c}{4}t - z\right) \text{rect}\left(\frac{t - \frac{4z}{c} + 2ns}{4ns}\right),$$

then

$$g(t) = -\frac{5c}{2}t \text{rect}\left(\frac{t + 2ns}{4ns}\right).$$

3. The pulse of sheet current $\mathbf{J}_s(t) = \hat{y}8t\text{rect}\left(\frac{t}{\tau}\right) \frac{\text{A}}{\text{m}}$ will produce magnetic and electric fields. The magnetic field is: (it's direction can be verify using the right-hand-rule for Ampere's law $\oint_C \mathbf{H} \cdot d\mathbf{l} = I_C$)

$$\begin{aligned} \mathbf{H}^\pm(z, t) &= \pm \frac{1}{2} J_{so} \left(t \pm \frac{z}{c} \right) \hat{x} \left[\frac{\text{A}}{\text{m}} \right] \\ &= \pm 4 \left(t \mp \frac{z}{c} \right) \text{rect} \left(\frac{t \mp \frac{z}{c}}{\tau} \right) \hat{x} \left[\frac{\text{A}}{\text{m}} \right] \quad \text{for } z \geq 0, \end{aligned}$$

whereas the electric field is given by

$$\mathbf{E}^\pm(x, t) = -\frac{\eta_o}{2} J_{so} \left(t \mp \frac{z}{c} \right) \hat{y} = -4\eta_o \left(t \mp \frac{z}{c} \right) \text{rect} \left(\frac{t \mp \frac{z}{c}}{\tau} \right) \hat{y} \left[\frac{\text{V}}{\text{m}} \right] \quad \text{for } z \geq 0,$$

where $\eta_o = 120\pi \Omega$ is the intrinsic impedance of free-space.

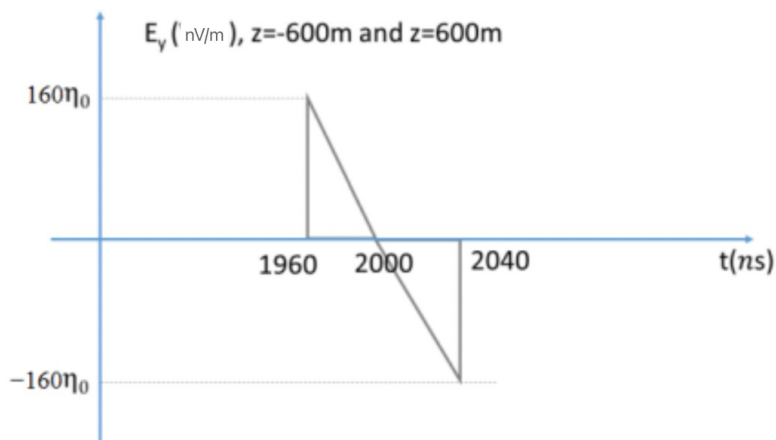
- (a) (8 points) The fields are given by

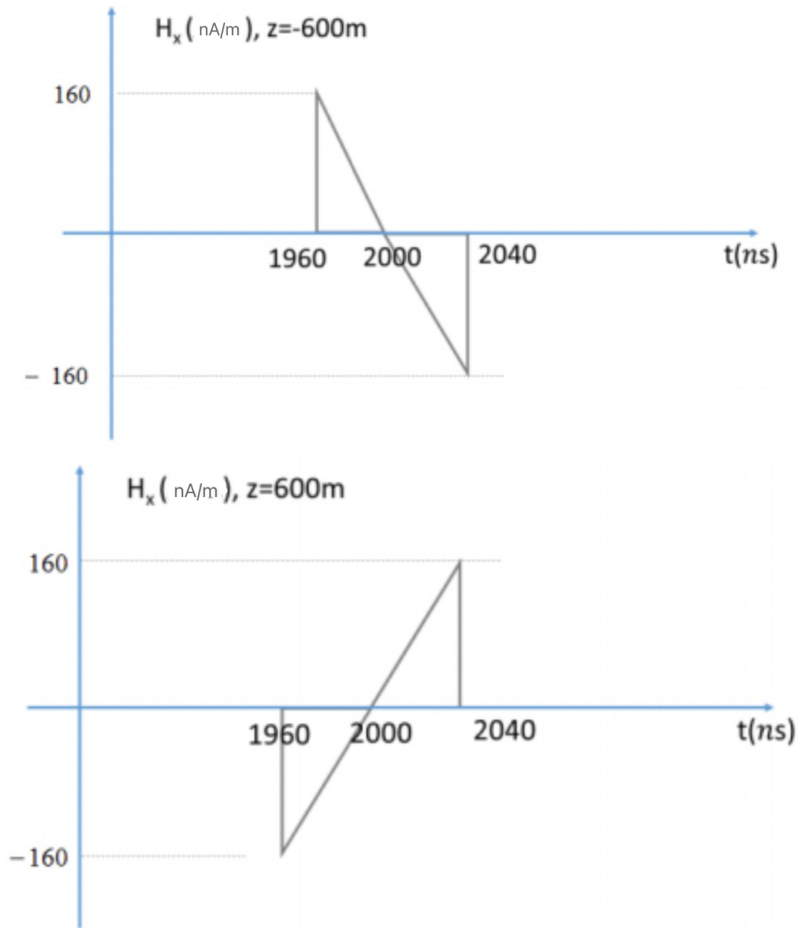
$$E_y(z, t)|_{z=-600\text{m}} = -4\eta_o (t - 2000\text{ns}) \text{rect} \left(\frac{t - 2000\text{ns}}{80\text{ns}} \right) \left[\frac{\text{V}}{\text{m}} \right]$$

$$E_y(z, t)|_{z=600\text{m}} = -4\eta_o (t - 2000\text{ns}) \text{rect} \left(\frac{t - 2000\text{ns}}{80\text{ns}} \right) \left[\frac{\text{V}}{\text{m}} \right]$$

$$H_x(z, t)|_{z=-600\text{m}} = -4 (t - 2000\text{ns}) \text{rect} \left(\frac{t - 2000\text{ns}}{80\text{ns}} \right) \left[\frac{\text{A}}{\text{m}} \right]$$

$$H_x(z, t)|_{z=600\text{m}} = 4 (t - 2000\text{ns}) \text{rect} \left(\frac{t - 2000\text{ns}}{80\text{ns}} \right) \left[\frac{\text{A}}{\text{m}} \right]$$

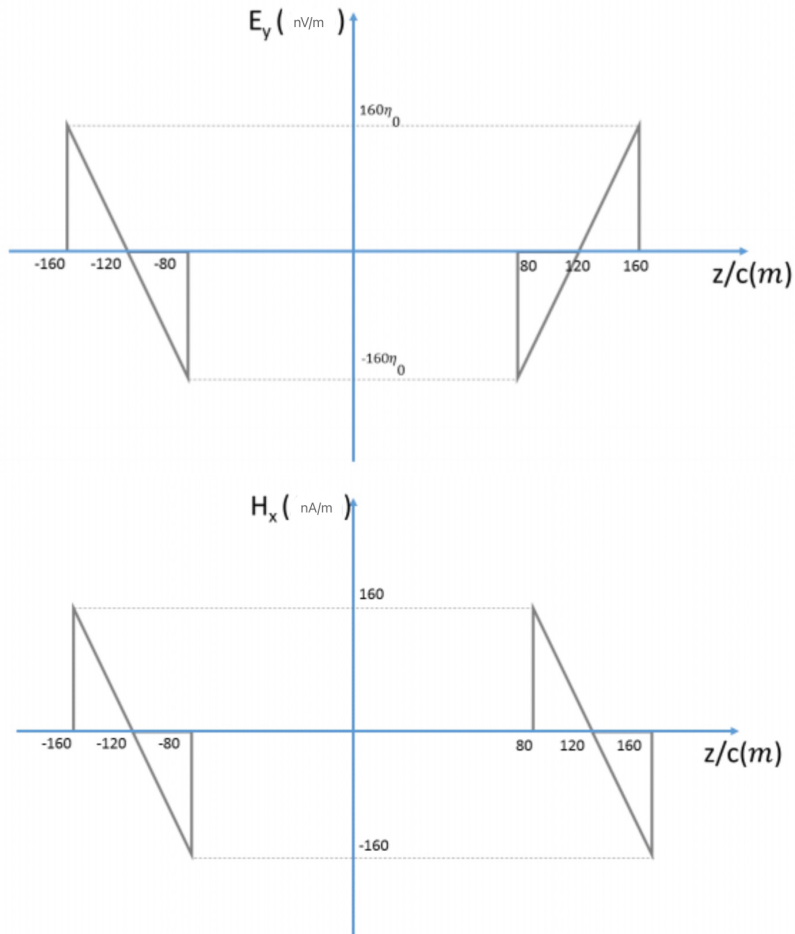




(b) (4 points) The fields are given by

$$E_y(z, t)|_{t=120\text{ns}} = -4\eta_0 \left(120\text{ns} \mp \frac{z}{c}\right) \text{rect}\left(\frac{120\text{ns} \mp \frac{z}{c}}{80\text{ns}}\right) \left[\frac{\text{V}}{\text{m}}\right]$$

$$H_x(z, t)|_{t=120\text{ns}} = \pm 4 \left(120\text{ns} \mp \frac{z}{c}\right) \text{rect}\left(\frac{120\text{ns} \mp \frac{z}{c}}{80\text{ns}}\right) \left[\frac{\text{A}}{\text{m}}\right]$$



(c) (2 points) Following the hint given in the problem, we can write

$$\begin{aligned} -\mathbf{J}_s \cdot \mathbf{E} &= -\left(\hat{y} 8t \operatorname{rect}\left(\frac{t}{\tau}\right)\right) \cdot \left(-4\eta_0 t \operatorname{rect}\left(\frac{t}{\tau}\right) \hat{y}\right) \\ &= 32\eta_0 t^2 \operatorname{rect}^2\left(\frac{t}{\tau}\right) \left[\frac{\text{W}}{\text{m}^2}\right]. \end{aligned}$$

Then, the TEM wave density energy is

$$\begin{aligned} \int -\mathbf{J}_s \cdot \mathbf{E} dt &= \int 32\eta_0 t^2 \operatorname{rect}^2\left(\frac{t}{\tau}\right) dt \\ &= \int_{-\tau/2}^{\tau/2} 32\eta_0 t^2 dt = \frac{32}{3}\eta_0 [t^3]_{-\tau/2}^{\tau/2} \\ &= \frac{8}{3}\eta_0 \tau^3 \approx 5.14 \times 10^{-19} \left[\frac{\text{J}}{\text{m}^2}\right]. \end{aligned}$$

4. (a) (3 points) For the plane wave described by $\mathbf{E}_1 = 3 \cos(\omega t - \beta z) \hat{y} \left[\frac{\text{V}}{\text{m}} \right]$:

- i. The magnetic field \mathbf{H} should satisfy $\mathbf{H} = -\frac{\mathbf{E} \times \hat{\beta}}{\eta}$, where $\hat{\beta}$ is the unit vector parallel to the propagation direction. Then, we can find the expressions for \mathbf{H} field of the given plane wave as

$$\mathbf{H}_1 = -\frac{3}{\eta_o} \cos(\omega t - \beta z) \hat{x} \left[\frac{\text{A}}{\text{m}} \right].$$

- ii. The instantaneous power flow density is given by the Poynting vector $\mathbf{S} = \mathbf{E} \times \mathbf{H}$. Therefore, the instantaneous power that crosses some surface A is given by $P = \int_A \mathbf{S} \cdot d\mathbf{A}$, where $d\mathbf{A} \equiv \hat{n} \cdot A$. Therefore, the Poynting vector is found as

$$\mathbf{S}_1 = \mathbf{E}_1 \times \mathbf{H}_1 = \frac{9}{\eta_o} \cos^2(\omega t - \beta z) \hat{z} \left[\frac{\text{W}}{\text{m}^2} \right],$$

and expression for instantaneous power that crosses a 1 m^2 area (i.e. $A = 1 \text{ m}^2$) in the xy -plane from $-z$ to $+z$ may be written as

$$P_1 = \frac{9}{\eta_o} \cos^2(\omega t - \beta z) [\text{W}].$$

- iii. We can calculate the time-average of the Poynting vector using the trigonometric identity: $\cos^2 \theta = \frac{1}{2} (1 + \cos(2\theta))$. Based on the fact that the time average of the cosine wave is zero ($\frac{1}{T} \int_T \cos(\omega t) dt = 0$), we can write

$$\langle \mathbf{S}_1 \rangle = \left\langle \frac{9}{\eta_o} \cos^2(\omega t - \beta z) \hat{z} \frac{\text{W}}{\text{m}^2} \right\rangle = \frac{9}{2\eta_o} \hat{z} \left[\frac{\text{W}}{\text{m}^2} \right].$$

Therefore, the average power that crosses some surface A is given by

$$\langle P_1 \rangle = \langle \mathbf{S}_1 \rangle \cdot \hat{n} A = \frac{9}{2\eta_o} [\text{W}].$$

(b) (3 points) For the plane wave described by $\mathbf{E}_2 = E_o (\cos(\omega t - \beta z) \hat{x} + \sin(\omega t - \beta z) \hat{y}) \left[\frac{\text{V}}{\text{m}} \right]$:

- i. The propagation direction is $\hat{\beta}_2 = \hat{z}$. Thus, the magnetic field \mathbf{H}_2 is given by

$$\mathbf{H}_2 = \frac{E_o}{\eta_o} (\cos(\omega t - \beta z) \hat{y} - \sin(\omega t - \beta z) \hat{x}) \left[\frac{\text{A}}{\text{m}} \right].$$

- ii. The Poynting vector is given by

$$\begin{aligned} \mathbf{S}_2 &= \mathbf{E}_2 \times \mathbf{H}_2 \\ &= E_o (\cos(\omega t - \beta z) \hat{x} + \sin(\omega t - \beta z) \hat{y}) \times \frac{E_o}{\eta_o} (\cos(\omega t - \beta z) \hat{y} - \sin(\omega t - \beta z) \hat{x}) \\ &= \frac{E_o^2}{\eta_o} (\cos^2(\omega t - \beta z) \hat{z} + \sin^2(\omega t - \beta z) \hat{z}) = \frac{E_o^2}{\eta_o} \hat{z} \left[\frac{\text{W}}{\text{m}^2} \right]. \end{aligned}$$

Therefore, the instantaneous power crossing the area $A = 1 \text{ m}^2$ is

$$P_2 = \frac{E_o^2}{\eta_o} [\text{W}].$$

iii. The Poynting vector is constant in time, thus the time-average power is

$$\langle P_2 \rangle = \frac{E_o^2}{\eta_o} [\text{W}].$$

(c) (3 points) For the plane wave described by $\mathbf{H}_3 = \cos(\omega t + \beta z + \frac{\pi}{3})\hat{x} - \sin(\omega t + \beta z - \frac{\pi}{6})\hat{y} [\frac{\text{A}}{\text{m}}]$:

i. The electric field \mathbf{E} should satisfy $\mathbf{E} = \eta(\mathbf{H} \times \hat{\beta})$ where $\hat{\beta} = -\hat{z}$ in this case. Then, we can find the expressions for the \mathbf{E} field of the given plane wave as

$$\mathbf{E}_3 = \eta_o \left(\cos(\omega t + \beta z + \frac{\pi}{3})\hat{y} + \sin(\omega t + \beta z - \frac{\pi}{6})\hat{x} \right) [\frac{\text{V}}{\text{m}}].$$

ii. The Poynting vector is given by

$$\begin{aligned} \mathbf{S}_3 &= \mathbf{E}_3 \times \mathbf{H}_3 \\ &= -\eta_o \left(\cos^2(\omega t + \beta z + \frac{\pi}{3})\hat{z} + \sin^2(\omega t + \beta z - \frac{\pi}{6})\hat{z} \right) \\ &= -\eta_o \left(\cos^2(\omega t + \beta z + \frac{\pi}{3})\hat{z} + \cos^2(\omega t + \beta z + \frac{\pi}{3})\hat{z} \right) \\ &= -2\eta_o \cos^2(\omega t + \beta z + \frac{\pi}{3})\hat{z} [\frac{\text{W}}{\text{m}^2}]. \end{aligned}$$

Therefore, the instantaneous power crossing the area $A = 1 \text{ m}^2$ is

$$P_3 = -2\eta_o \cos^2(\omega t + \beta z + \frac{\pi}{3}) [\text{W}].$$

iii. The time-average power crossing a 1 m^2 area is

$$\langle P_3 \rangle = -\eta_o [\text{W}].$$

(d) (3 points) For the plane wave described by $\mathbf{H}_4 = \cos(\omega t - \beta x)\hat{z} + \sin(\omega t - \beta x)\hat{y} [\frac{\text{A}}{\text{m}}]$:

i. The propagation direction is $\hat{\beta}_4 = \hat{x}$. Thus, the electric field \mathbf{E}_4 is given by

$$\mathbf{E}_4 = \eta_o (\cos(\omega t - \beta x)\hat{y} - \sin(\omega t - \beta x)\hat{z}) [\frac{\text{V}}{\text{m}}].$$

ii. The wave is propagating in the $+x$ direction, therefore there is no flux of energy flowing into the z direction. Therefore, the instantaneous power crossing a 1 m^2 area in the xy -plane from $-z$ to z is

$$P_4 = 0 [\text{W}].$$

iii. The time-average power crossing a 1 m^2 area in the xy -plane from $-z$ to z is also

$$\langle P_4 \rangle = 0 [\text{W}].$$

5. (10 points) The current produced is proportional to $|\mathbf{E}_1 + \mathbf{E}_2|^2$, hence it can be described by

$$I = A |\mathbf{E}_1 + \mathbf{E}_2|^2 = AE_o (\cos(\omega_1 t - \beta_1 z) + \cos(\omega_2 t - \beta_2 z))^2.$$

Since we are not interested in the z dependence, we may very well just ignore it by setting $z=0$. Also remember:

$$\begin{aligned} \cos(a)^2 &= \frac{\cos(2a) + 1}{2} \\ \cos(a+b) + \cos(a-b) &= 2 \cos(a) \cos(b) \end{aligned}$$

Therefore:

$$\begin{aligned} I|_{z=0} &= AE_o [\cos(\omega_1 t - \beta_1 z) + \cos(\omega_2 t - \beta_2 z)]^2|_{z=0} \\ &= AE_o [\cos(\omega_1 t) + \cos(\omega_2 t)]^2 \\ &= AE_o [\cos^2(\omega_1 t) + \cos^2(\omega_2 t) + 2 \cos(\omega_1 t) \cos(\omega_2 t)] \\ &= AE_o \left[\frac{2 + \cos(2\omega_1 t) + \cos(2\omega_2 t)}{2} \right] \\ &+ AE_o [\cos((\omega_1 - \omega_2)t) + \cos((\omega_1 + \omega_2)t)] \end{aligned}$$

Hence, the frequencies produced are $2\omega_1$, $2\omega_2$, $\omega_1 + \omega_2$ and $|\omega_1 - \omega_2|$.