

1. (a) (4 points) Since both ends of the transmission line (TL) are open, each resonant mode will have resonance frequency

$$f = \frac{v}{2l}n, \quad n \leq 1.$$

Plugging in numbers, we get

$$f = \frac{1 \times 10^8}{2 \times 5}n = 10n \text{ [MHz]}$$

- (b) (3 points) The general solution to TL problems are

$$\begin{aligned}\tilde{V}(z) &= V^+ e^{-j\beta z} + V^- e^{j\beta z} \\ \tilde{I}(z) &= \frac{V^+}{Z_o} e^{-j\beta z} - \frac{V^-}{Z_o} e^{j\beta z}\end{aligned}$$

Since both ends are open, the current at the TL ends must be zero. At $z = 0$ end, plug in the current expression, we get:

$$V^+ = V^-.$$

Therefore, the voltage and the current expressions are simplified to

$$\begin{aligned}\tilde{V}(z) &= 2V^+ \cos(\beta z) \\ \tilde{I}(z) &= -\frac{V^+}{Z_o} 2j \sin(\beta z).\end{aligned}$$

Using the relations $\beta = \frac{2\pi}{\lambda}$ and $\lambda = \frac{v}{f} = \frac{2l}{n}$, we can get the expression for current phasor:

$$\tilde{I}(z) = -\frac{V^+}{Z_o} 2j \sin\left(\frac{n\pi}{l}z\right).$$

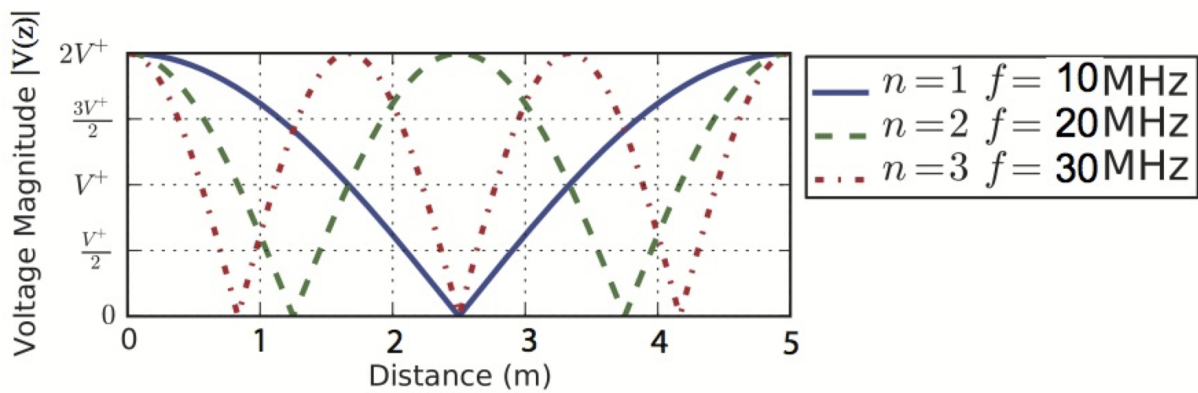
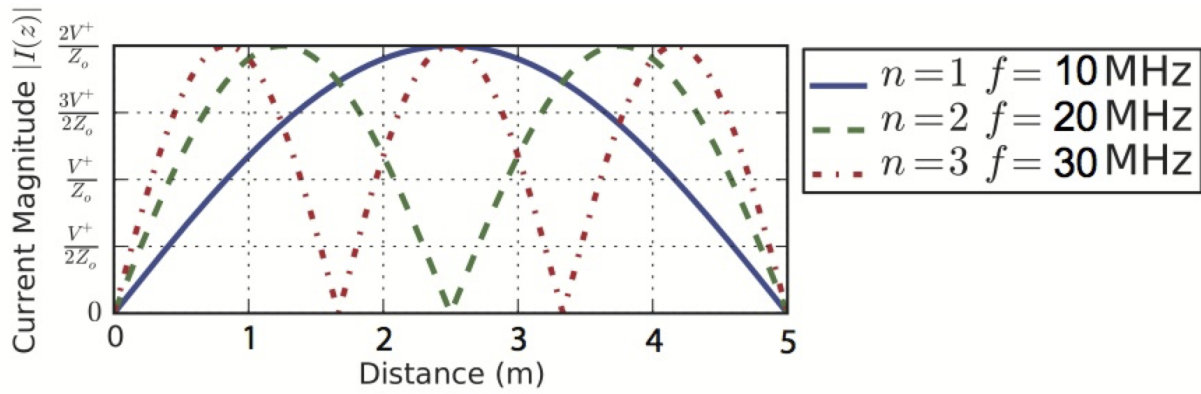
The current plot is shown in the following figure. Notice that each mode n has n half wavelengths fitted into the length of TL, and higher modes oscillates with higher frequencies.

- (c) (3 points) We repeat the same step from part (b) to the voltage expression. We get

$$\tilde{V}(z) = 2V^+ \cos\left(\frac{n\pi}{l}z\right),$$

and the voltage is printed in the following figure. We can see that there is also n half wavelengths for mode n . Additionally, we observe that the null of voltage and current on the TL are separately by a quarter wavelength.

Note: Keep in mind the convention used here is to place the load at $z = 0$ and source at $z = -L$. Alternatively, we can say the load is placed at $d = 0$ and source is $d = +L$. Remember to stay consistent!



2. (a) (2 points) From previous question, we know the voltage and current expressions on TL with load end open are:

$$\begin{aligned}\tilde{V}(z) &= 2V^+ \cos(\beta z) \\ \tilde{I}(z) &= -\frac{V^+}{Z_o} 2j \sin(\beta z).\end{aligned}$$

Therefore we have the input impedance of the open TL as

$$Z_{in} = Z(-l) = \frac{\tilde{V}(-l)}{\tilde{I}(-L)} = -jZ_o \cot(\beta l)$$

- (b) (3 points) By the steady state analysis, we can regard all circuit elements as general impedance. Thus, the impedance of the capacitor will be

$$Z_C = \frac{1}{j\omega C}.$$

Using the voltage division rule, we write

$$V_L = V_g \frac{R_L}{R_L + R_g + Z_{in} + Z_C}.$$

Given that $R_g = R_L$ and $V_L = \frac{1}{2}V_g$, we find $Z_{in} + Z_C = 0$. Therefore, we can write

$$-jZ_o \cot(\beta L) = \frac{1}{j\omega C}$$

from which we obtain

$$C = -\frac{1}{\omega Z_o \cot(\beta L)} = -\frac{1}{\omega Z_o \cot(\frac{\omega}{v}L)} = -63.66 \text{ pF}.$$

Note: Capacitance should be positive. The length in the problem is not practical.

3. (a) (1 point) The signal wavelength λ is $\lambda = \frac{v}{f} = \frac{1 \times 10^8}{\pi \times 10^8 / 2\pi} = 2$ m.
- (b) (1 point) Since the end of the line (i.e. at $d = 0$) is open circuited, $I(0) = 0$.
- (c) (2 points) Yes, because the magnitude of the reflection coefficient at the open termination is one (i.e., $|\Gamma_L| = 1$), then, the reflected wave has the same amplitude as the incident wave, which is enough to set up a standing wave.
- (d) (2 points) If $I_R = 2$ A then the current in the transmission line at $d = l$ is zero (i.e., $I(l) = 0$). Since successive current nulls are $\lambda/2$ apart, the smallest non-zero transmission line length is $l = \lambda/2 = 1$ m.
- (e) (2 points) If $l = \frac{\lambda}{2}$ then the voltage at $d = l/2$ is zero (i.e., $V(\frac{l}{2}) = 0$), because half-way between successive current nulls there is a voltage null.
- (f) (2 points) If $l = \frac{\lambda}{2}$ then the current at $d = l/2$ can not be zero (i.e., $I(\frac{l}{2}) \neq 0$), because at this point the current is a maximum.
- (g) (1 point) If $I_R = 0$ then the voltage at $d = l$ is also zero (i.e., $V(l) = 0$). Since voltage and current nulls are $\lambda/4$ apart, the smallest non-zero transmission line length is $l = \frac{\lambda}{4} = 0.5$ m.
- (h) (2 points) If $I(\frac{l}{2}) = 0$ and $I(0) = 0$ then $\frac{l}{2} = \frac{\lambda}{2}$ because the separation between two successive current nulls is $\lambda/2$ (non-trivial case). Therefore, $l = \lambda = 2$ m.
- (i) (2 points) In part (h), $I(\frac{l}{2}) = 0$ and $l = \lambda$ which implies that $I(l) = 0$. Since $I_R = 2 \angle 0^\circ$ A then $V(l) = I_R \times 50 \Omega = 100 \angle 0^\circ$ V.

4. (a) (2 points) Referring to page 3 of Lecture 33, the input current is given by

$$I_{in} \equiv I\left(\frac{\lambda}{4}\right) = \frac{jV^+ + jV^-}{Z_o} = j\frac{V(0)}{Z_o} = j\frac{V_L}{Z_o}$$

Therefore

$$V_L = -jI_{in}Z_o = -j50 \text{ V.}$$

- (b) (2 points) If the load end is shorted then $V_L = 0$. This implies that the current input must also be zero ($I_{in} = j\frac{V_L}{Z_o} = 0$), and therefore, $I_{in} = 0.5\angle 0^\circ \text{ A}$ is not a possibility in this situation.

- (c) (2 points) Since the voltage measured at any distance along the TL is given by

$$V(d) = V^+(e^{j\beta d} + \Gamma_L e^{-j\beta d}),$$

where $\Gamma_L = \frac{Z_L - Z_o}{Z_L + Z_o}$, the voltages at one end is

$$V_L = V(0) = V^+(1 + \Gamma_L)$$

whereas the input voltage of the quarter-wave section at the other end is

$$V_{in} = V\left(\frac{\lambda}{4}\right) = jV^+(1 - \Gamma_L).$$

Therefore, we can write

$$\frac{V_L}{V_{in}} = -j\frac{1 + \Gamma_L}{1 - \Gamma_L} = -j\frac{Z_L}{Z_o},$$

from which we can obtain

$$V_{in} = j\frac{V_L Z_o}{Z_L} = 100 \text{ V.}$$

5. (a) (1 point) Voltage input at the load end of the half-wave transformer

$$V_{L_1} = -V_{in} = -j10 \text{ V.}$$

- (b) (1 point) Voltage input at the load end of the quarter-wave transformer

$$V_{L_2} = -j \frac{V_{in} Z_{L_2}}{Z_o} = 5 \text{ V.}$$

- (c) (1 point) Current input at the load end of the half-wave transformer

$$I_{L_1} = \frac{V_{L_1}}{Z_{L_1}} = \frac{-j}{20} \text{ A.}$$

- (d) (1 point) Current input at the load end of the quarter-wave transformer

$$I_{L_2} = -j \frac{V_{in}}{Z_o} = \frac{V_{L_2}}{Z_{L_2}} = \frac{1}{10} \text{ A.}$$

- (e) (2 points) Input impedance of the combined network is given as

$$Z_{in} = Z_{in_1} || Z_{in_2} = Z_{L_1} || \frac{Z_o^2}{Z_{L_2}} = 100 \Omega.$$

- (f) (2 points) Total time-averaged power absorbed is

$$\begin{aligned} P_{avg} &= \frac{1}{2} \text{Re}\{V_{L_1} I_{L_1}^*\} + \frac{1}{2} \text{Re}\{V_{L_2} I_{L_2}^*\} \\ &= \frac{1}{2} \text{Re}\{-j10 \times \frac{j}{20}\} + \frac{1}{2} \text{Re}\{(5) \times (\frac{1}{10})\} = 0.5 \text{ W.} \end{aligned}$$

6. (a) (4 points) Since the $\frac{\lambda}{2}$ -transformer inverts the sign of its voltage and current inputs at the load end, and two $\frac{\lambda}{2}$ transformers will be a double negative. We can write

$$V_{L1} = -V_g = -100 \text{ V}$$

which implies that

$$I_{L1} = \frac{V_{L1}}{R_{L1}} = -2 \text{ A.}$$

We know that a $\frac{\lambda}{4}$ -transformer has a load current $I_L = -j\frac{V_{in}}{Z_o}$. It follows that

$$\begin{aligned} V_{L2} &= I_{L2}R_{L2} \\ &= -j\frac{V_{in}}{Z_{o2}}R_{L2} = -j50 \text{ V,} \end{aligned}$$

which yields

$$I_{L2} = \frac{V_{L2}}{R_{L2}} = -j\frac{V_{in}}{Z_{o2}} = -j \text{ A.}$$

- (b) (2 points) The power dissipated in resistor R_{L1} is

$$P_{L1} = \frac{1}{2}\text{Re}\{V_{L1}I_{L1}^*\} = \frac{1}{2}\text{Re}\left\{\frac{|V_{L1}|^2}{R_{L1}}\right\} = 100 \text{ W.}$$

Then for the second line which is a $\frac{\lambda}{4}$ -transformer, the time average dissipated power in resistor R_{L2} is given by

$$P_{L2} = \frac{1}{2}\text{Re}\{V_{L2}I_{L2}^*\} = 25 \text{ W.}$$

7. **Bonus Problem** : An air filled (i.e. $\mu = \mu_o$ and $\epsilon = \epsilon_o$) parallel-plate transmission line with length $l = 1$ m, width $w = 5$ cm, and plate separation $d = 0.5$ cm is shorted at both ends.

- (a) (1 point) As both ends of the TL has been “short circuited”, and thus TL voltage $V(z, t)$ needs to vanish at $z = 0$ and $z = 1$. Since

$$V(z, t) = V^+(t - \frac{z}{v}) + V^-(t + \frac{z}{v})$$

from d'Alembert solutions, these boundary conditions

$$V(0, t) = V^+(t) + V^-(t) = 0$$

and

$$V(l, t) = V^+(t - \frac{l}{v}) + V^-(t + \frac{l}{v}) = 0$$

require that

$$V^+(t) = -V^-(t),$$

and

$$V^+(t - \frac{l}{v}) = -V^-(t + \frac{l}{v}) \Rightarrow V^+(t') = V^+(t' + \frac{2l}{v}),$$

where $t' = t - \frac{l}{v}$. Finally, we can conclude that period $T = \frac{2l}{v}$, and fundamental frequency $\omega_o = \frac{2\pi}{T} = \frac{\pi v}{l}$. At the lowest resonant frequency l should be half a wavelength long, therefore

$$f_o = \frac{\omega_o}{2\pi} = \frac{v}{2l} = 1.5 \times 10^8 \text{ Hz} = 150 \text{ MHz}.$$

- (b) (2 points) The phasor counterpart of a co-sinusoidal wave $V(z, t)$ is

$$\tilde{V}(z) = V^+ e^{-j\beta z} + V^- e^{j\beta z}.$$

Using the result from part (a) $V^+(t) = -V^-(t)$, we get

$$\tilde{V}(z) = V^+ e^{-j\beta z} - V^+ e^{j\beta z} = -j2V^+ \sin(\beta z),$$

which leads

$$V(z, t) = \text{Re} \left\{ \tilde{V}(z) e^{j\omega t} \right\} = 2V^+ \sin(\beta z) \sin(\omega t),$$

where $2V^+ = V_o$ (maximum voltage of the resonant mode), and therefore

$$V(z, t) = V_o \sin(\beta z) \sin(\omega t) \text{ V}.$$

- (c) (2 points) Utilizing the Telegrapher's equation $-\frac{\partial V}{\partial z} = \mathcal{L} \frac{\partial I}{\partial t}$ together with the expressions for $V(z, t)$ obtained in part (b), we can find

$$-\frac{\partial V}{\partial z} = -\beta V_o \cos(\beta z) \sin(\omega t),$$

which yields

$$\frac{\partial I}{\partial t} = -\frac{\beta V_o}{\mathcal{L}} \cos(\beta z) \sin(\omega t).$$

Next, finding the anti-derivative of the above, we get

$$I(z, t) = \frac{\beta V_o}{\omega \mathcal{L}} \cos(\beta z) \cos(\omega t) \text{ A}.$$

(d) (3 points) The stored energy is given by

$$W = \int_{x=0}^d \int_{y=0}^w \int_{z=0}^{\ell} \left(\frac{1}{2} \epsilon_o E_x^2 + \frac{1}{2} \mu_o H_y^2 \right) dx dy dz.$$

Given that $V = E_x d$ and $I = H_y w$ and substituting V and I with those expressions derived in parts (b) and (c), respectively, we obtain a more complicated expression for the stored energy

$$W = \int_{x=0}^d \int_{y=0}^w \int_{z=0}^{\ell} \left[\frac{1}{2} \epsilon_o \frac{V_o^2}{d^2} \sin^2(\beta z) \sin^2(\omega t) + \frac{1}{2} \mu_o \frac{\beta^2 V_o^2 w^2}{w^2 \omega^2 \mu_o^2 d^2} \cos^2(\beta z) \cos^2(\omega t) \right] dx dy dz,$$

where $\beta = \omega \sqrt{\epsilon_o \mu_o} \implies \epsilon_o = \frac{\beta^2}{\omega^2 \mu_o}$ which lets us work through the previous equation

$$W = \frac{1}{2} \epsilon_o \frac{V_o^2}{d^2} \left[\int_{x=0}^d \int_{y=0}^w \int_{z=0}^{\ell} \left\{ \sin^2(\beta z) \sin^2(\omega t) + \cos^2(\beta z) \cos^2(\omega t) \right\} dx dy dz \right],$$

$$W = \frac{1}{2} \epsilon_o \frac{w V_o^2}{d} \left[\int_{z=0}^{\ell} \left\{ \sin^2(\beta z) \sin^2(\omega t) + \cos^2(\beta z) \cos^2(\omega t) \right\} dz \right],$$

$$W = \frac{1}{2} C V_o^2 \left[\int_{z=0}^{\ell} \left\{ \sin^2(\beta z) \sin^2(\omega t) + \cos^2(\beta z) \cos^2(\omega t) \right\} dz \right],$$

$$W = \frac{1}{2} C V_o^2 \left[\sin^2(\omega t) \left(\frac{\ell}{2} - \frac{\sin(2\beta\ell)}{4\beta} \right) + \cos^2(\omega t) \left(\frac{\ell}{2} + \frac{\sin(2\beta\ell)}{4\beta} \right) \right].$$

Now, for the lowest resonant frequency $f_o = \frac{v}{2\ell}$ (from part (a)), we get $\lambda = 2\ell = 2$ (as $\ell = 1$ m). Also, $\sin(2\beta\ell) = \sin(2 \times 2\pi \times 2) = 0$. Therefore, the above equation will reduce to

$$W = \frac{1}{4} C \ell V_o^2 [\sin^2(\omega t) + \cos^2(\omega t)],$$

$$W = \frac{1}{4} C \ell V_o^2 \text{ J.}$$

The numerical value of proportionality constant is $\frac{1}{4}$.

(e) (2 points) From part (d), we know that $W = \frac{1}{4} C \ell V_o^2$. The expected value of V_o^2 will be

$$E[V_o^2] = \frac{E[W]}{\frac{1}{4} C \ell} = \frac{k_B T}{\frac{1}{4} C \ell} = \frac{4 \times 1.38 \times 10^{-23} \times 300}{10 \times 8.85 \times 10^{-12} \times 1},$$

which gives us

$$E[V_o^2] = 1.87 \times 10^{-10}.$$

Therefore, the square root of the expected value of V_o^2 will be $\sqrt{E[V_o^2]} = 13.68 \times 10^{-6}$ V.