20 Poynting theorem and monochromatic waves

- The magnitude of **Poynting vector**

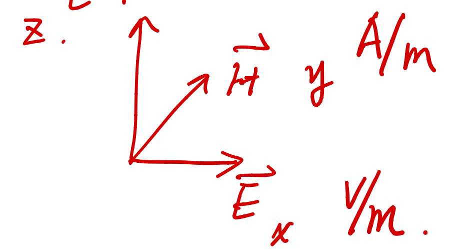
  pointing

  $$\mathbf{S} = \mathbf{E}\times\mathbf{H}$$

  W/m^2. $\frac{\partial \rho}{\partial t} + \nabla\cdot\vec{J} = 0$.

  represents the amount of power transported — often called energy flux — by electromagnetic fields **E** and **H** over a unit area transverse to the $\mathbf{E}\times\mathbf{H}$ direction.

$\vec{E}\times\vec{H} = \vec{S}$.



$$\frac{V}{m}\cdot\frac{A}{m} = \frac{VA}{m^2} \rightarrow \frac{W}{m^2}.$$

This interpretation of the Poynting vector is obtained from a conservation law extracted from Maxwell's equations (see margin) as follows:

$$\begin{aligned}
\nabla\cdot\mathbf{D} &= \rho\\
\nabla\cdot\mathbf{B} &= 0\\
\nabla\times\mathbf{E} &= -\frac{\partial\mathbf{B}}{\partial t}\\
\nabla\times\mathbf{H} &= \mathbf{J}+\frac{\partial\mathbf{D}}{\partial t}.
\end{aligned}$$

1. Dot multiply Faraday's law by **H**, dot multiply Ampere's law by **E**,

$$\begin{aligned}
(\nabla\times\mathbf{E} &= -\frac{\partial\mathbf{B}}{\partial t})\cdot\mathbf{H}\\
-(\nabla\times\mathbf{H} &= \mathbf{J}+\frac{\partial\mathbf{D}}{\partial t})\cdot\mathbf{E}
\end{aligned}$$

   and take their difference:

$$\underbrace{\mathbf{H}\cdot\nabla\times\mathbf{E}-\mathbf{E}\cdot\nabla\times\mathbf{H}}_{\nabla\cdot(\mathbf{E}\times\mathbf{H})} = \underbrace{-\frac{\partial\mathbf{D}}{\partial t}\cdot\mathbf{E}-\frac{\partial\mathbf{B}}{\partial t}\cdot\mathbf{H}}_{-\frac{\partial}{\partial t}(\frac{1}{2}\epsilon\mathbf{E}\cdot\mathbf{E}+\frac{1}{2}\mu\mathbf{H}\cdot\mathbf{H})}-\mathbf{J}\cdot\mathbf{E}.$$

$\vec{D} = \varepsilon\vec{E}$
$\vec{B} = \mu\vec{H}$.

2. After re-arrangements shown above, the result can be written as

storage. energy transportation power dissipation.

$$\frac{\partial}{\partial t}(\frac{1}{2}\epsilon \mathbf{E}\cdot\mathbf{E}+\frac{1}{2}\mu\mathbf{H}\cdot\mathbf{H})+\nabla\cdot(\mathbf{E}\times\mathbf{H})+\mathbf{J}\cdot\mathbf{E}=0.$$

Joule Heating. $A/m^2 \cdot V/m = VA/m^3$

Poynting theorem

- **Poynting theorem** derived above is a *conservation law* just like the continuity equation $\frac{\partial\rho}{\partial t}+\nabla\cdot\mathbf{J}=0$:

  charge transporter

  $E \to J$

  – The first term on the left,

$$\frac{\partial}{\partial t}(\frac{1}{2}\epsilon \mathbf{E}\cdot\mathbf{E}+\frac{1}{2}\mu\mathbf{H}\cdot\mathbf{H}),$$

  is time rate of change of total electric and magnetic **energy** density.

Hence, **Poynting theorem is the conservation law for electromagnetic energy**, just like continuity equation is the conservation law for electric charge.

  – The second term

$$\nabla\cdot(\mathbf{E}\times\mathbf{H})$$

  accounts for energy transport in Poynting theorem, just like $\nabla\cdot\mathbf{J}$ accounts for charge transport in the continuity equation. Therefore

$$\mathbf{S}\equiv\mathbf{E}\times\mathbf{H}$$

is **energy flux per unit area** measured in

$$\frac{\text{V}}{\text{m}}\frac{\text{A}}{\text{m}} = \frac{\text{W}}{\text{m}^2} = \frac{\text{J/s}}{\text{m}^2}$$

units, just like $\mathbf{J}$ is charge flux per unit area in $\frac{\text{C/s}}{\text{m}^2} = \frac{\text{A}}{\text{m}^2}$ units.

– Finally, the last term in Poynting theorem (repeated in the margin),

$$\mathbf{J} \cdot \mathbf{E}$$

is called **Joule heating**, and it represents power absorbed per unit volume (which can only be non-zero in the presence of $\mathbf{J}$).

If $\mathbf{J}\cdot\mathbf{E}$ is negative in any region, then $\mathbf{J}$ in that region is acting as a source of electromagnetic energy, just like any circuit component with negative vi is acting as an energy source in the electrical circuit.

Note that $\mathbf{J} \cdot \mathbf{E}$ had a negative value on the current sheet radiator examined in last lecture. We return to the current sheet radiator in the next example.

Poynting thm:

$\frac{\partial}{\partial t}(\frac{1}{2}\epsilon\mathbf{E} \cdot \mathbf{E} + \frac{1}{2}\mu\mathbf{H} \cdot \mathbf{H}) + \nabla \cdot (\mathbf{E} \times \mathbf{H}) + \mathbf{J} \cdot \mathbf{E} = 0$

Example 1: On $z = 0$ plane we have a time-harmonic surface current specified as

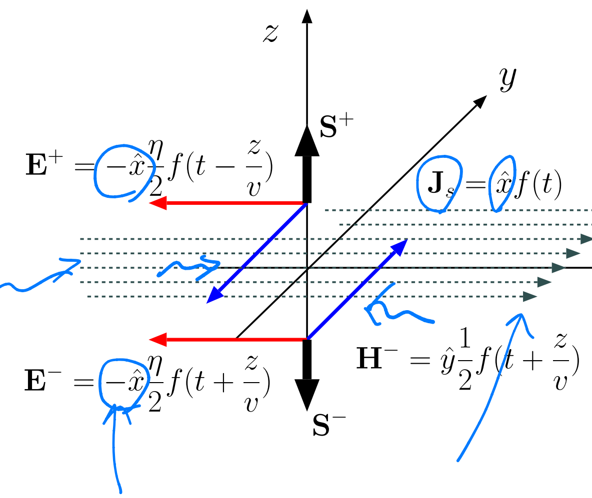
$$\mathbf{J}_s = \hat{x} f(t) = \hat{x} 2\cos(\omega t)\,\frac{\mathrm{A}}{\mathrm{m}}$$

where ω is some frequency of oscillation. $f(t)$.

(a) Determine the radiated TEM wave fields $\mathbf{E}(z,t)$ and $\mathbf{H}(z,t)$ in the regions $z \gtrless 0$,

(b) The associated Poynting vectors $\mathbf{E} \times \mathbf{H}$, and

(c) $\mathbf{J}_s \cdot \mathbf{E}$ on the current sheet.



$-\eta\cos(\omega(t - \frac{z}{c}))$

Solution:

$\checkmark$ $\cos(\omega t - \beta z)$. $\frac{\omega}{c}$.

$\sin(\omega t + \beta z)$. $-z$.

(a) With reference to the solution of the current sheet radiator depicted in the margin (from last lecture), we that an x-polarized surface current $f(t)$ produces the wave fields

$$E_x = -\frac{\eta}{2} f(t \mp \frac{z}{v}) \quad \text{and} \quad H_y = \mp\frac{1}{2} f(t \mp \frac{z}{v})$$

in the surrounding regions propagating away from the current sheet on both sides. Given that $f(t) = 2\cos(\omega t)$, this implies that

$$E_x = -\eta\cos(\omega t \mp \beta z) \quad \text{and} \quad H_y = \mp\cos(\omega t \mp \beta z)$$

where

$\frac{2\pi}{\beta} = \lambda$. wave number $\rightarrow$

$$\beta = \frac{\omega}{c} \quad \text{and} \quad \eta = \eta_o \approx 120\pi\,\Omega$$

$\rightarrow \omega = 2\pi f$. rad/s. $1/f = T. = \frac{2\pi}{\omega}$.

since the current sheet is surrounded by vacuum. Hence in vector form we have

$$\mathbf{E}(z,t) = -\eta\cos(\omega t \mp \beta z)\hat{x}\,\frac{\mathrm{V}}{\mathrm{m}} \quad \text{and} \quad \mathbf{H}(z,t) = \mp\cos(\omega t \mp \beta z)\hat{y}\,\frac{\mathrm{A}}{\mathrm{m}},$$

where the upper signs are for $z > 0$, and lower signs for $z < 0$.

(b) The associated Poynting vectors are

$$\mathbf{S} = \mathbf{E} \times \mathbf{H} = \pm\eta \cos^2(\omega t \mp \beta z)\hat{z} \frac{\mathrm{W}}{\mathrm{m}^2}.$$

120π

Note that the time-average value of vector $\mathbf{S}$ points in the direction of wave propagation on both sides of the current sheet.

(c) Since on $z = 0$ surface of the current sheet the electric field vector is

$$\mathbf{E}(0,t) = -\eta \cos(\omega t)\hat{x} \frac{\mathrm{V}}{\mathrm{m}},$$

it follows that $\mathbf{J}_s \cdot \mathbf{E}$ on the same surface is

$$\mathbf{J}_s(t) \cdot \mathbf{E}(0,t) = (\hat{x}2 \cos(\omega t) \frac{\mathrm{A}}{\mathrm{m}}) \cdot (-\eta \cos(\omega t)\hat{x} \frac{\mathrm{V}}{\mathrm{m}}) = -2\eta \cos^2(\omega t) \frac{\mathrm{W}}{\mathrm{m}^2}.$$

$\frac{1}{2}$.

$\rightsquigarrow -\eta \ \mathrm{W/m^2}$.

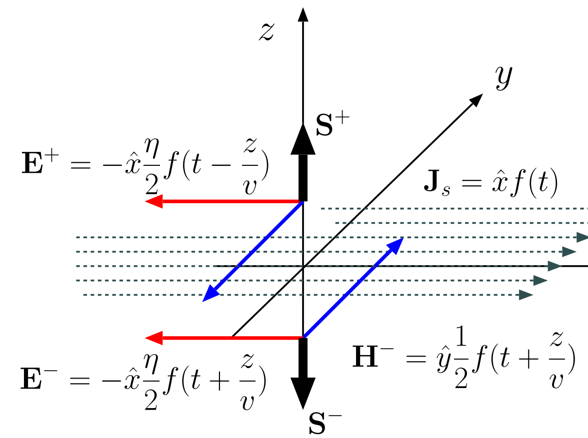
- In the above example, a time-harmonic source current oscillating at some frequency ω produced "monochromatic waves" of radiated fields propagating away from the current sheet on both sides.

  – The calculations showed time-varying Poynting vectors $\mathbf{E} \times \mathbf{H}$. The time-averaged values of these time-varying vectors can be easily determined by making use of the trig identity

$$\cos^2(\omega t + \phi) = \frac{1}{2}[1 + \cos(2\omega t + 2\phi)].$$

Since the time-average of the second term on the right is zero, we can express the time-average of this identity as

$$\langle\cos^2(\omega t+\phi)\rangle = \langle\frac{1}{2}[1+\cos(2\omega t+2\phi)]\rangle = \frac{1}{2},$$



where the angular brackets denote the time-averaging procedure.

- Consequently, the result

$$\mathbf{E}\times\mathbf{H} = \pm\eta\cos^2(\omega t\mp\beta z)\hat{z}\,\frac{\mathrm{W}}{\mathrm{m}^2}$$

from Example 1 implies that

$+\hat{z}$ $z>0$.

$$\langle\mathbf{E}\times\mathbf{H}\rangle = \pm\eta\frac{1}{2}\hat{z}\,\frac{\mathrm{W}}{\mathrm{m}^2} = \pm 60\pi\hat{z}\,\frac{\mathrm{W}}{\mathrm{m}^2},$$

$-\hat{z}$ $z<0$.

which represent the time-average power per unit area transported by the waves radiated by the current sheet.

- In Poynting theorem the Joule heating term $\mathbf{J}\cdot\mathbf{E}$ is **power absorbed per unit volume**, and, accordingly, $\boxed{-\mathbf{J}\cdot\mathbf{E}}$ is **power injected per unit volume**.

  – Likewise, $\pm\mathbf{J}_s\cdot\mathbf{E}$ can be interpreted as **power absorbed/injected per unit area** on a surface.

In Example 1 above we calculated an instantaneous injected power density of

$$-\mathbf{J}_s \cdot \mathbf{E} \;=\; 2\eta\cos^2(\omega t)\,\frac{\mathrm{W}}{\mathrm{m}^2}.$$

Clearly, its time-aveage works out as

$$\langle -\mathbf{J}_s \cdot \mathbf{E}\rangle \;=\; \eta\,\frac{\mathrm{W}}{\mathrm{m}^2} = 120\pi\,\frac{\mathrm{W}}{\mathrm{m}^2}.$$

– Note that $\langle -\mathbf{J}_s \cdot \mathbf{E}\rangle$ exactly matches the sum of $|\langle \mathbf{E} \times \mathbf{H}\rangle|$ calculated on both sides of the current sheet, in conformity with energy conservation principle (Poynting theorem).