

1. Given that

$$\mathbf{E} = \hat{x}\sin(y) + \hat{y}\cos(x)$$

Let us calculate the following,

a) Curl of \mathbf{E} ,

$$\nabla \times \mathbf{E} = \begin{pmatrix} \hat{x} & \hat{y} & \hat{z} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ \sin(y) & \cos(x) & 0 \end{pmatrix} = \left(\frac{\partial}{\partial x}(\cos(x)) - \frac{\partial}{\partial y}(\sin(y)) \right) \hat{z} = (-\sin(x) - \cos(y)) \hat{z}.$$

Curl of curl of \mathbf{E} ,

$$\nabla \times (\nabla \times \mathbf{E}) = \begin{pmatrix} \hat{x} & \hat{y} & \hat{z} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ 0 & 0 & -\sin(x) - \cos(y) \end{pmatrix} = \sin(y) \hat{x} + \cos(x) \hat{y}.$$

b) Applying Gauss law(in differential form)

$$\nabla \cdot \mathbf{D} = \epsilon_0 \nabla \cdot \mathbf{E} = \rho,$$

we find that

$$\rho = \epsilon_0 \nabla \cdot \mathbf{E} = 0.$$

2. Given an electrostatic potential $V(x,y,z) = x^2 - 2V$ in certain region of space, let us calculate the following,

a) Electrostatic field \mathbf{E} ,

$$\mathbf{E} = -\nabla V = -\frac{\partial}{\partial x}(x^2 - 2)\hat{x} - \frac{\partial}{\partial y}(x^2 - 2)\hat{y} - \frac{\partial}{\partial z}(x^2 - 2)\hat{z} = -2x\hat{x}\left[\frac{V}{m}\right].$$

b) Curl of \mathbf{E} ,

$$\nabla \times \mathbf{E} = \begin{pmatrix} \hat{x} & \hat{y} & \hat{z} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ -2x & 0 & 0 \end{pmatrix} = 0.$$

c) Divergence of \mathbf{E} ,

$$\nabla \cdot \mathbf{E} = \frac{\partial}{\partial x}(-2x) = -2.$$

d) Charge density ρ ,

$$\rho = \epsilon_0 \nabla \cdot \mathbf{E} = -2\epsilon_0 \left[\frac{C}{m^3}\right].$$

3. In electrostatics, we generate a curl-free vector field $\mathbf{E}(x,y,z)$ if we take the gradient of a scalar function $V(x,y,z)$. Therefore, $\nabla \times \mathbf{E} = \nabla \times (-\nabla V) = 0$.

The same answer can be reached by calculating:

$$\mathbf{E} = -\nabla V = -2x\hat{x} + z\hat{y} + y\hat{z}$$

$$\nabla \times \mathbf{E} = \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ -2x & z & y \end{vmatrix} = 0$$

4. Starting with the left-hand side of $\nabla \times (\nabla \times \mathbf{A}) = \nabla (\nabla \cdot \mathbf{A}) - \nabla^2 \mathbf{A}$, we write

$$\begin{aligned} \nabla \times (\nabla \times \mathbf{A}) &= \nabla \times \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ x-y & x+y & 2 \end{vmatrix} \\ &= \nabla \times (2\hat{z}) = 0. \end{aligned}$$

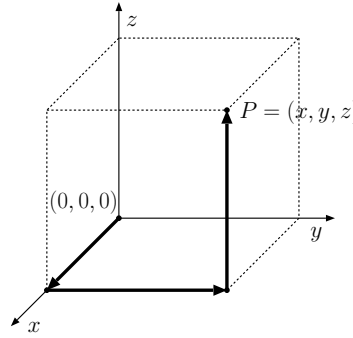
Solving the right-hand side of the equation, we obtain

$$\begin{aligned} \nabla (\nabla \cdot \mathbf{A}) - \nabla^2 \mathbf{A} &= \nabla (\nabla \cdot ((x-y)\hat{x} + (x+y)\hat{y} + 2\hat{z})) - \nabla^2 ((x-y)\hat{x} + (x+y)\hat{y} + 2\hat{z}) \\ &= \nabla (2) - \mathbf{0} = \mathbf{0}. \end{aligned}$$

Consequently, we can see that $\nabla \times (\nabla \times \mathbf{A}) = \nabla (\nabla \cdot \mathbf{A}) - \nabla^2 \mathbf{A}$ is verified since both sides have the same solutions.

5.

a) The electrostatic potential V at any point $P = (x, y, z)$ can be calculated by performing a vector line integral by using the path shown in the below figure.



Therefore, we can write

$$\begin{aligned} V(P) - V(0) &= - \int_0^P \mathbf{E} \cdot d\mathbf{l} \\ &= - \int_0^x E_x(x, 0, 0) dx - \int_0^y E_y(x, y, 0) dy - \int_0^z E_z(x, y, z) dz \\ &= -\frac{43}{2} \text{ V}. \end{aligned}$$

Given that $V(0) = 0 \text{ V}$, the electrostatic potential at $P = (1, 2, 3)$ is $V(1, 2, 3) = -\frac{43}{2} \text{ V}$.

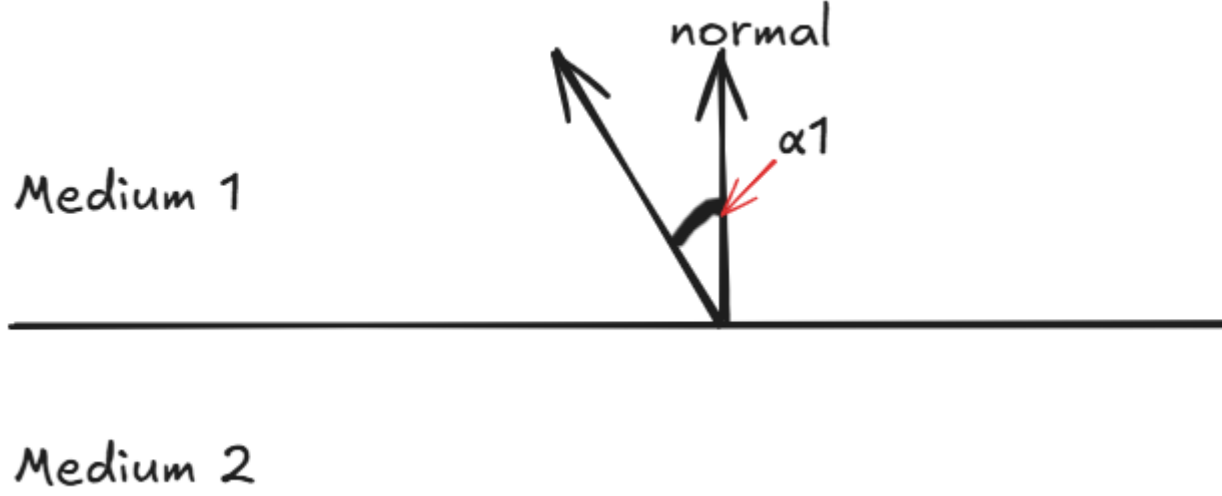
b) Using differential form of Gauss' Law,

$$\rho = \nabla \cdot \mathbf{D} = \epsilon_o(4 + 4z)$$

At $(0,0,0)$, $\rho = 4\epsilon_o \frac{C}{m^3}$ and at $(1,2,3)$, $\rho = 16\epsilon_o \frac{C}{m^3}$.

6.

a) In this part, consider the following geometry:



We know that the boundary conditions for the electric field at an interface are $\hat{n} \cdot (\mathbf{D}^+ - \mathbf{D}^-) = \rho_s$ and $\hat{n} \times (\mathbf{E}^+ - \mathbf{E}^-) = 0$. We also know that the electric field magnitude E_1 is 10 V/m at $\alpha_1 = 30^\circ$ from the normal. Therefore, we can compute the components of the field in medium 1.

$$E_{1n} = E_1 \cos(30^\circ) = 5\sqrt{3} \text{ V/m} \approx 8.66 \text{ V/m}$$

$$E_{1t} = E_1 \sin(30^\circ) = 5 \text{ V/m}$$

Since the tangential component of the electric field is continuous at a boundary, we know that $E_{1t} = E_{2t} = 5 \text{ V/m}$. From the boundary condition on the normal component of the displacement field, $E_{2n} = \frac{\epsilon_1}{\epsilon_2} E_{1n} = \frac{5\sqrt{3}}{4} \text{ V/m} \approx 2.165 \text{ V/m}$. Therefore, $E_2 = \sqrt{E_{2n}^2 + E_{2t}^2} \text{ V/m} \approx 5.449 \text{ V/m}$.

The direction of the field in medium 2, α_2 , is given by $\tan^{-1}(\frac{E_{2t}}{E_{2n}}) \approx 66.7^\circ$.

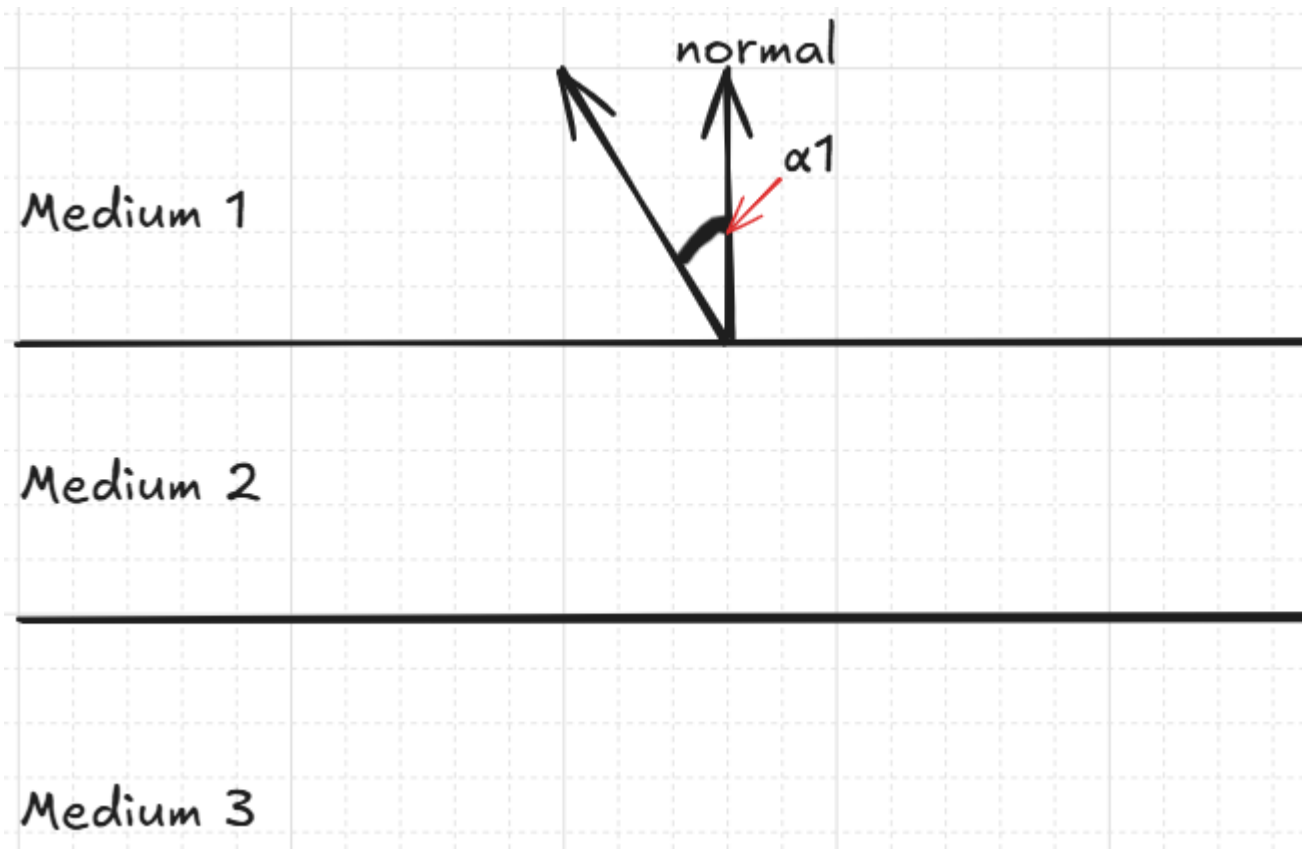
b) Again, consider the geometry from part a).

Now, $|\mathbf{E}_2| = 10 \text{ V/m}$, so we must compute its components in order to determine the charge density on the boundary.

Still tangential $E_{2t} = E_{1t} = 5 \text{ V/m}$, so normal component $E_{2n} = \sqrt{10^2 - 5^2} \approx 8.66 \text{ V/m}$

Evaluating the boundary condition $\rho_s = \epsilon_1 E_{1n} - \epsilon_2 E_{2n} \approx \epsilon_0 (10 \cos(30^\circ) - 4 \cdot 8.66) \approx -3 \times 8.66 \epsilon_0 \approx -25.98 \epsilon_0 \text{ C/m}^2$.

c) Now, the geometry has changed from the previous two parts.



It is stated that the boundaries between the media are charge-free, and that medium 2 is now a slab of finite thickness, while medium 3 is free space. To find α_3 , the angle from normal of the electric field in medium 3, we must again use boundary conditions.

Since $\rho_{s,23} = 0$, we know that $\epsilon_2 E_{2n} = \epsilon_3 E_{3n}$ and $E_{2t} = E_{3t}$ due to the continuity of tangential electric fields at boundaries. Since $\epsilon_2 E_{2n} = \epsilon_3 E_{3n}$, $E_{3n} = \frac{\epsilon_2}{\epsilon_3} E_{2n} = 4 \cdot E_{2n}$. From part a), $E_{2n} = \frac{5\sqrt{3}}{4} V/m$ and $E_{2t} = 5 V/m$.

Then we can find E_{3n} as $4 \cdot \frac{5\sqrt{3}}{4} V/m = 5\sqrt{3} V/m$, which is also the value of E_{1n} !

Finally, since $E_{3t} = E_{2t} = 5 V/m$, $\alpha_3 = \tan^{-1}\left(\frac{E_{3t}}{E_{3n}}\right) = 30^\circ$.