

Generalized Quantum States

via Mixed States and Density Operators

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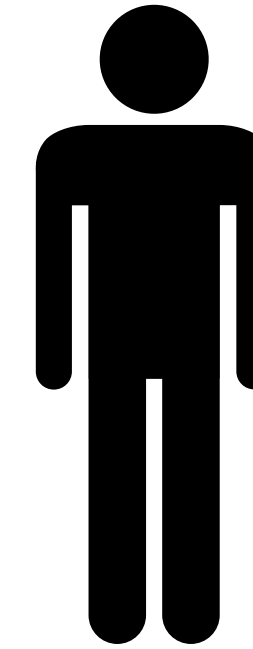
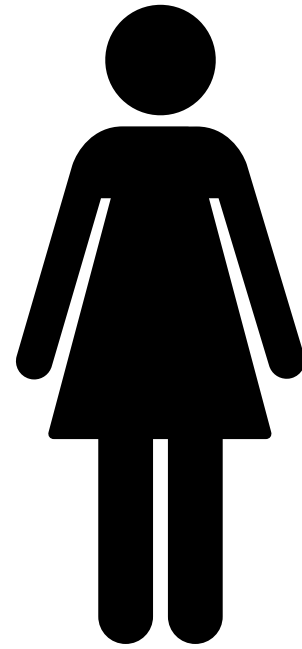
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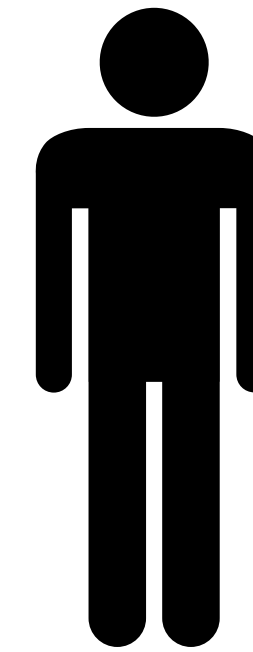
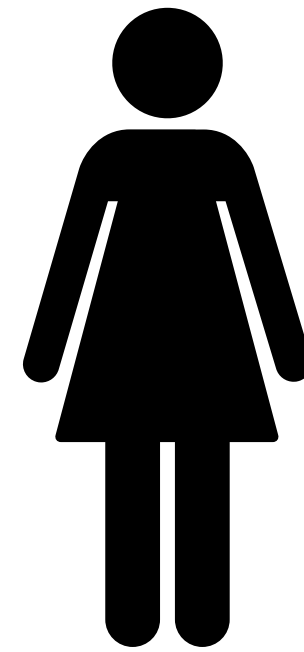
Is this kind of description sufficient?

Scenario 1

Suppose Alice and Bob share an EPR pair

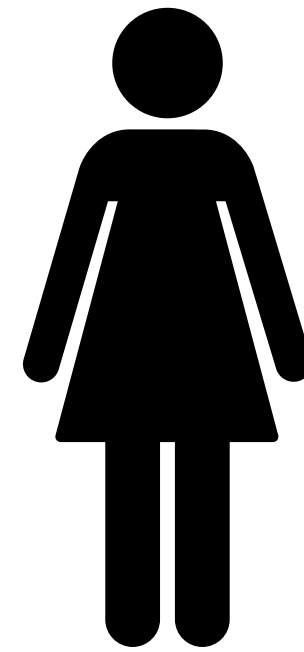


Suppose Alice and Bob share an EPR pair



$$|\mathbf{EPR}\rangle = \frac{1}{\sqrt{2}}|0\rangle^A|0\rangle^B + \frac{1}{\sqrt{2}}|1\rangle^A|1\rangle^B$$

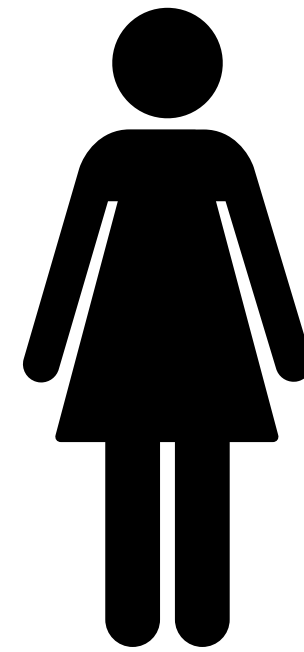
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Goodbye, I have to go

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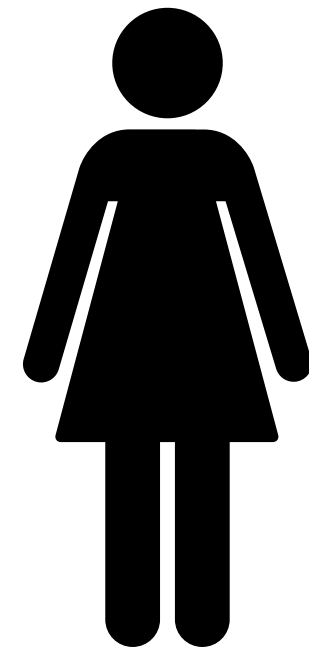


What is the state of Alice's quantum system?

$$|\mathbf{EPR}\rangle = \frac{1}{\sqrt{2}}|0\rangle^A|0\rangle^B + \frac{1}{\sqrt{2}}|1\rangle^A|1\rangle^B \neq |\psi_A\rangle \otimes |\psi_B\rangle$$

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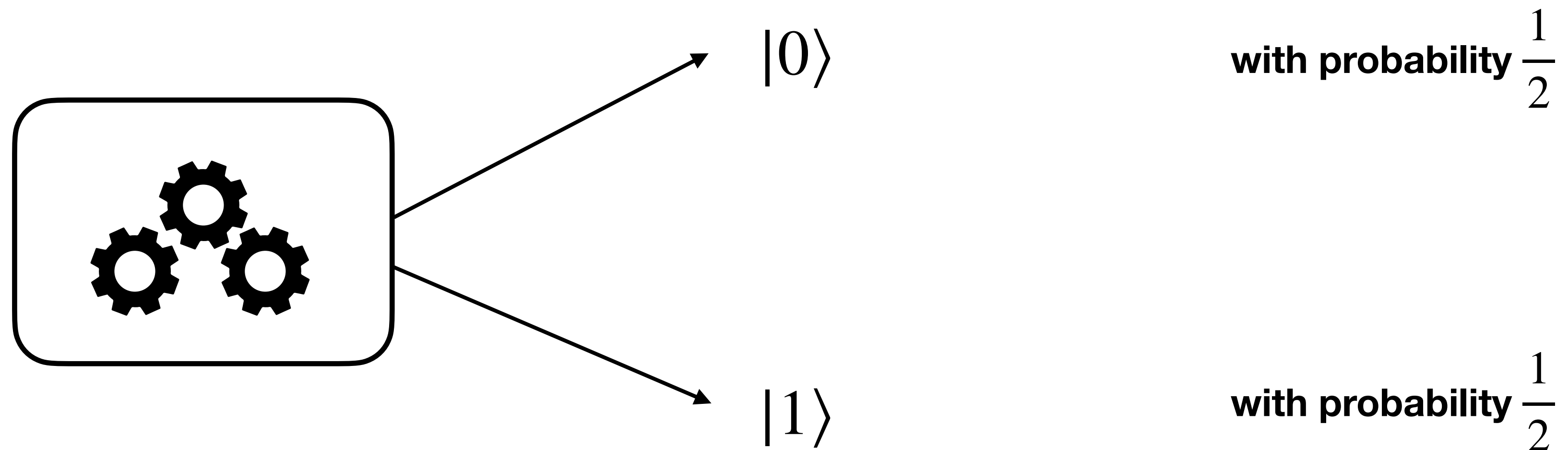
It is not a pure state!

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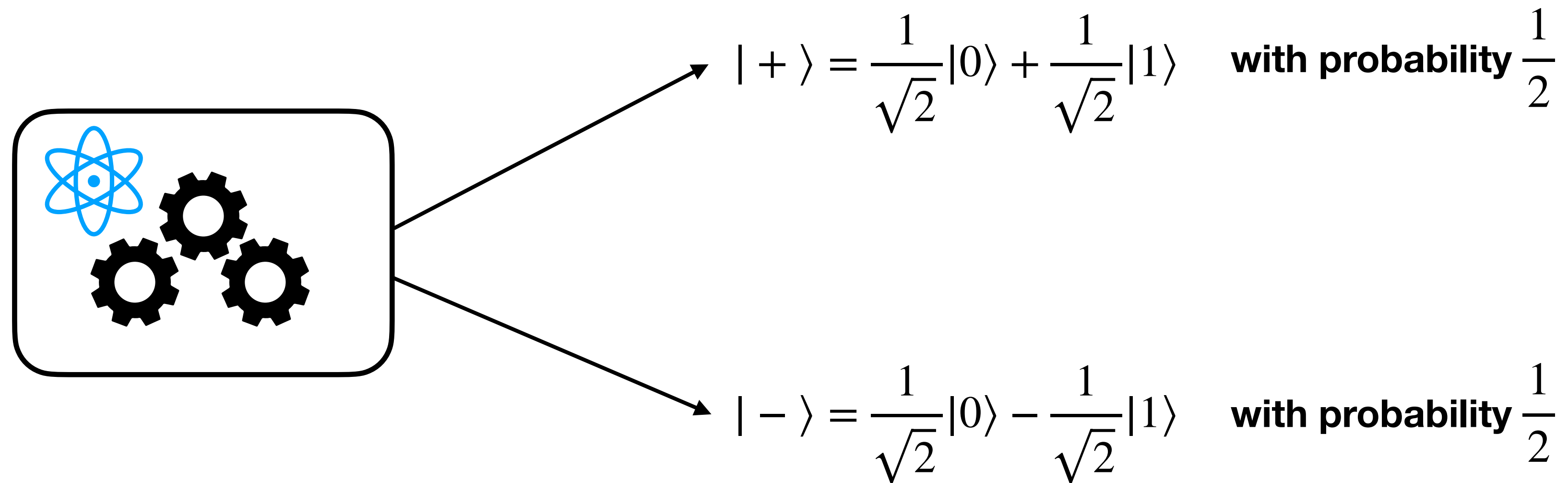
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Scenario 2

Suppose a Classical Machine Produces States

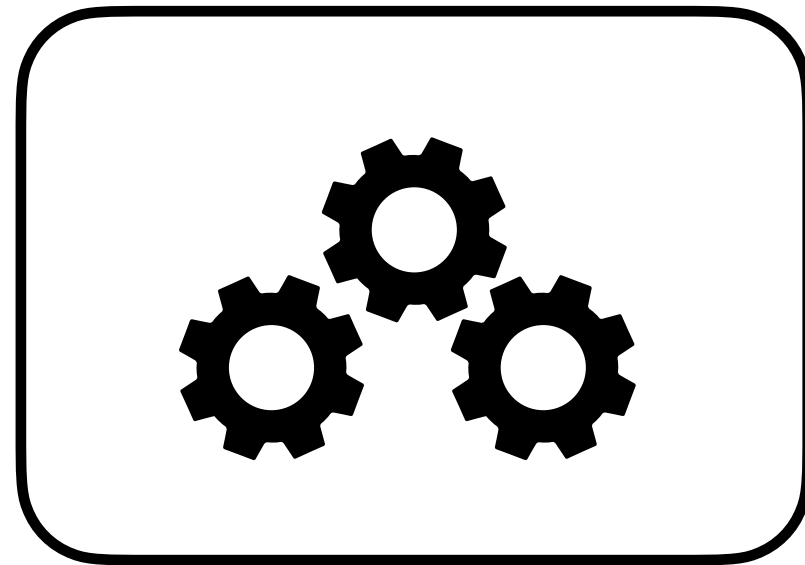


Suppose a Quantum Machine Produces States

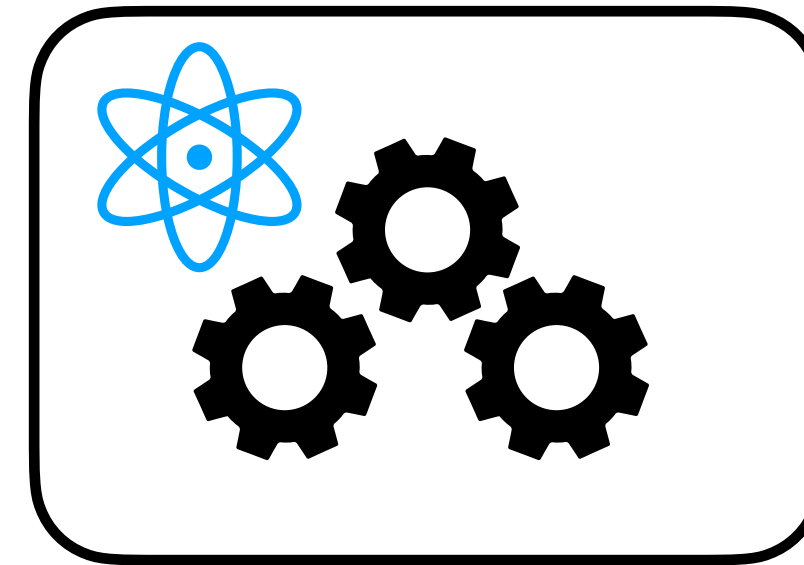


How can we describe the output quantum state?

How different are the outputs of these machines?



Classical



Quantum

**We need a more general formalism to
describe quantum states!**

The Ensemble Formalism

$$[m] = \{1, 2, \dots, m\}$$

Given

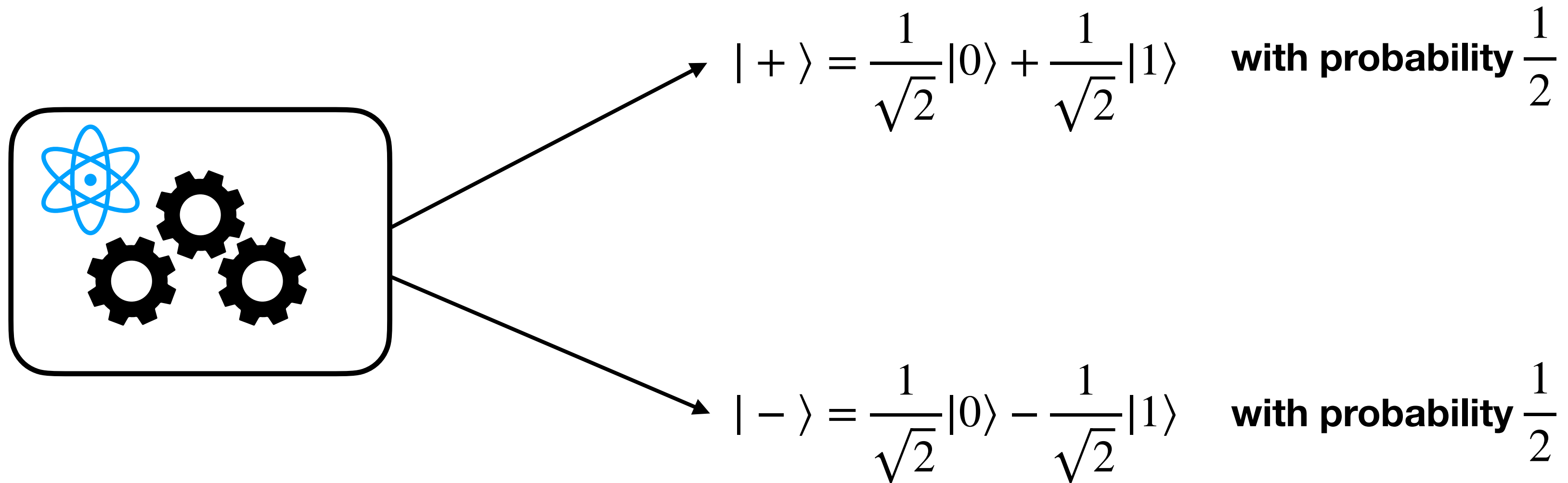
(1) $\{p_i\}_{i \in [m]}$ **probability distribution** $\left(\sum_{i=1}^m p_i = 1, p_i \geq 0 \forall i \in [m]\right)$ **and**

(2) $\{|\psi_i\rangle \in \mathbb{C}^d\}_{i \in [m]}$ **pure quantum states,**

we define the *ensemble* of quantum states:

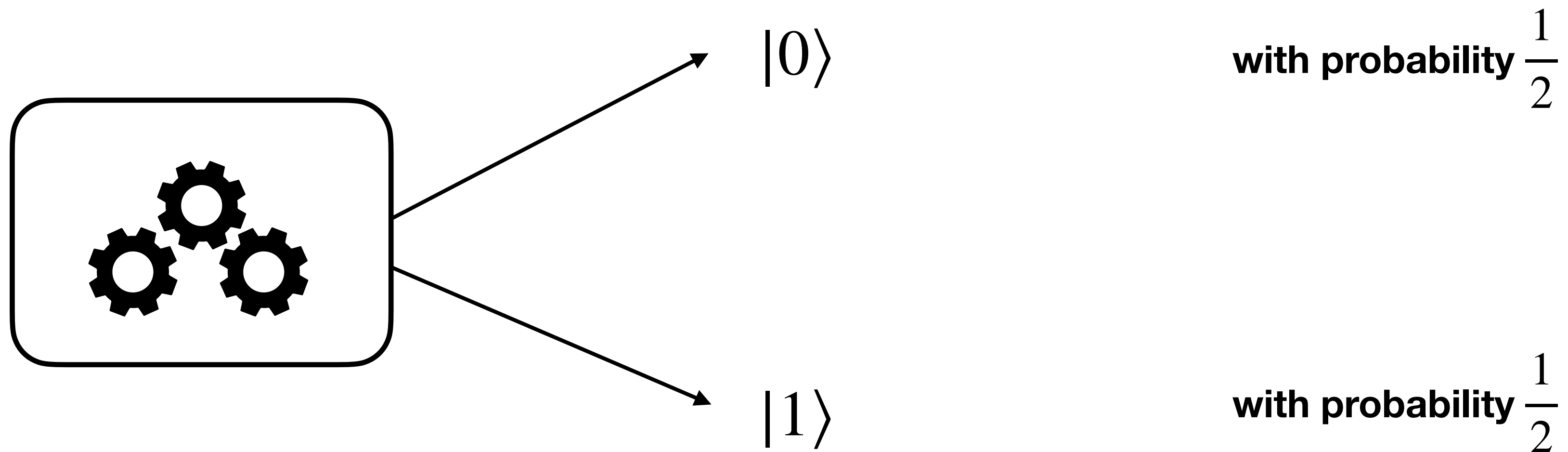
$$\left\{ (p_i, |\psi_i\rangle) \right\}_{i \in [m]}$$

Example of Ensembles



$$\left\{ \left(\frac{1}{2}, |+\rangle \right), \left(\frac{1}{2}, |-\rangle \right) \right\}$$

Example of Ensembles



$$\left\{ \left(\frac{1}{2}, |0\rangle \right), \left(\frac{1}{2}, |1\rangle \right) \right\}$$

Example of Ensembles

$|\psi\rangle \in \mathbb{C}^d$ **(a pure quantum state)**

$$\{(1, |\psi\rangle)\}$$

Density Operator Formalism

A quantum state is a matrix $\rho \in \mathbb{C}^{d \times d}$ satisfying

(1) **$\text{Tr}(\rho) = 1$**

(2) **$\rho \succeq 0$ positive semi-definite (PSD), i.e., $\rho = \rho^\dagger$ and $\langle \phi | \rho | \phi \rangle \geq 0, \forall |\phi\rangle \in \mathbb{C}^d$**

From Ensembles to Density Operators

Given an ensemble $\left\{ (p_i, |\psi_i\rangle) \right\}_{i \in [m]}$, what is the corresponding density operator?

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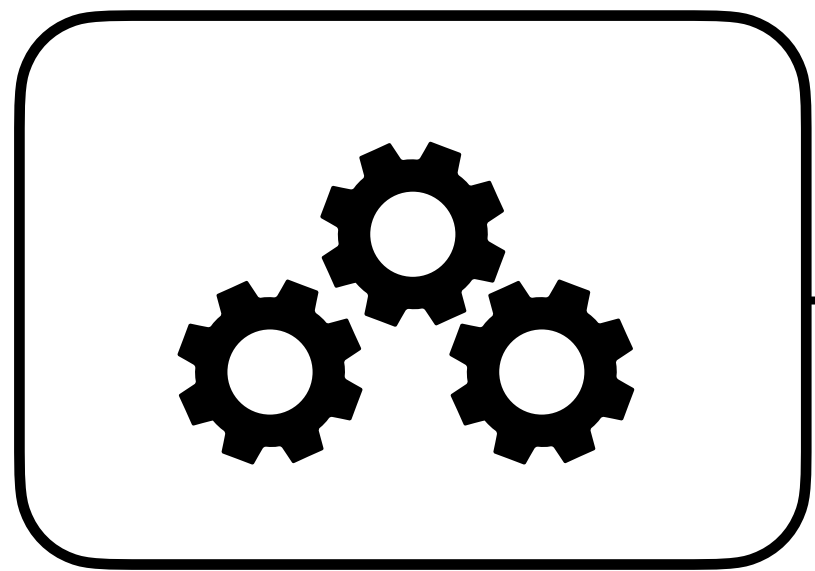
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(3) $\langle \phi | \rho | \phi \rangle = \sum_{i=1}^m p_i \langle \phi | \psi_i \rangle \langle \psi_i | \phi \rangle = \sum_{i=1}^m p_i |\langle \phi | \psi_i \rangle|^2 \geq 0, \forall |\phi\rangle \in \mathbb{C}^d$ ✓

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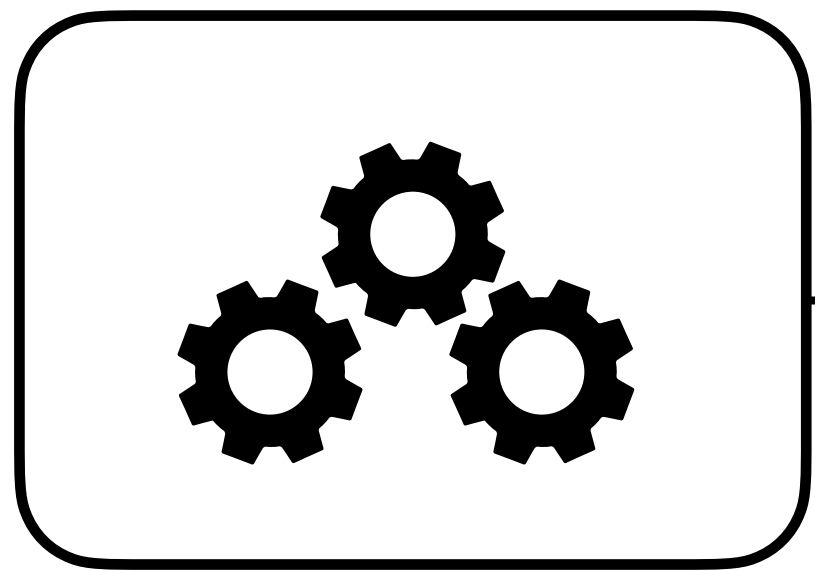


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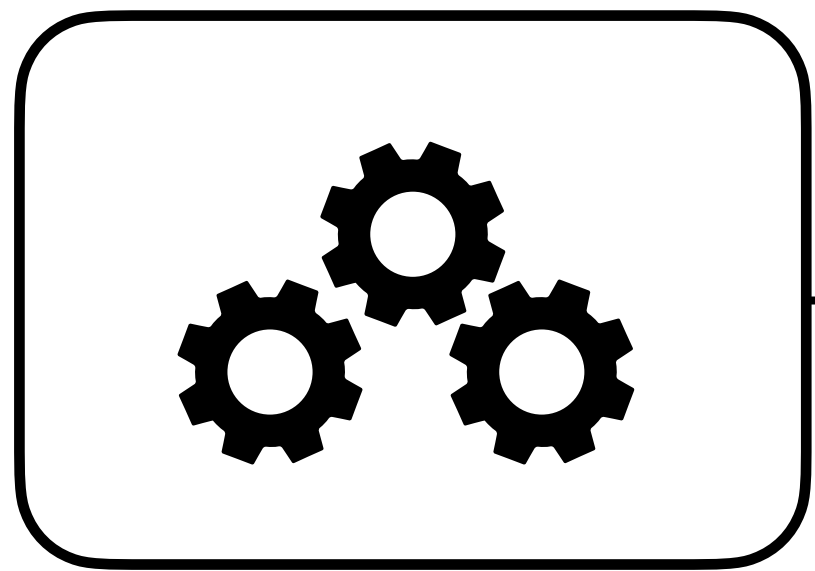


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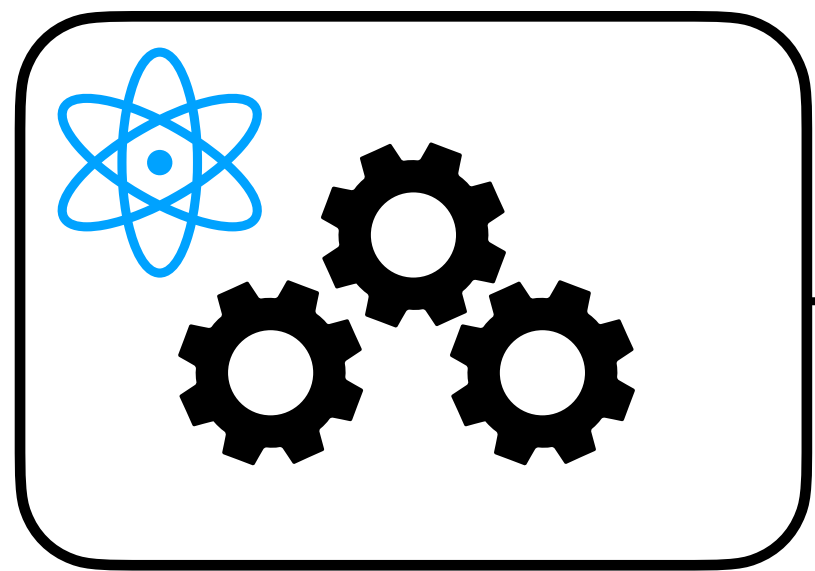
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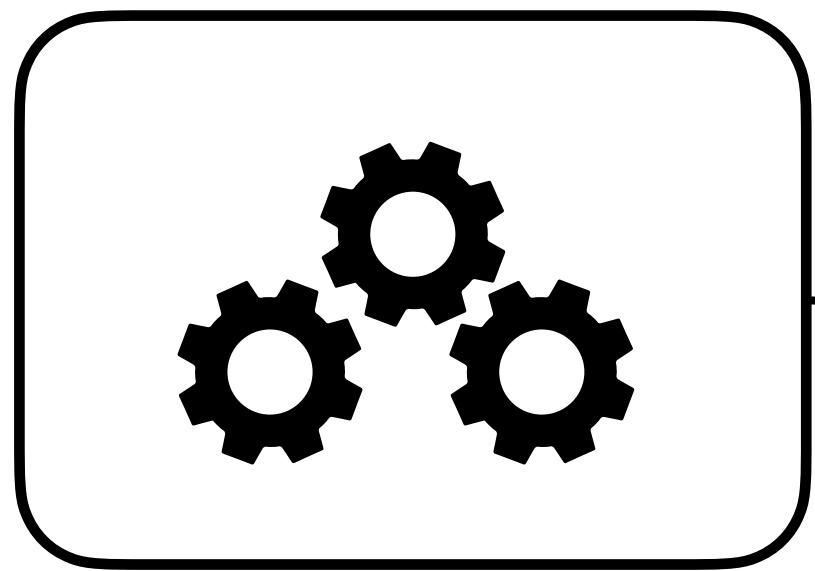


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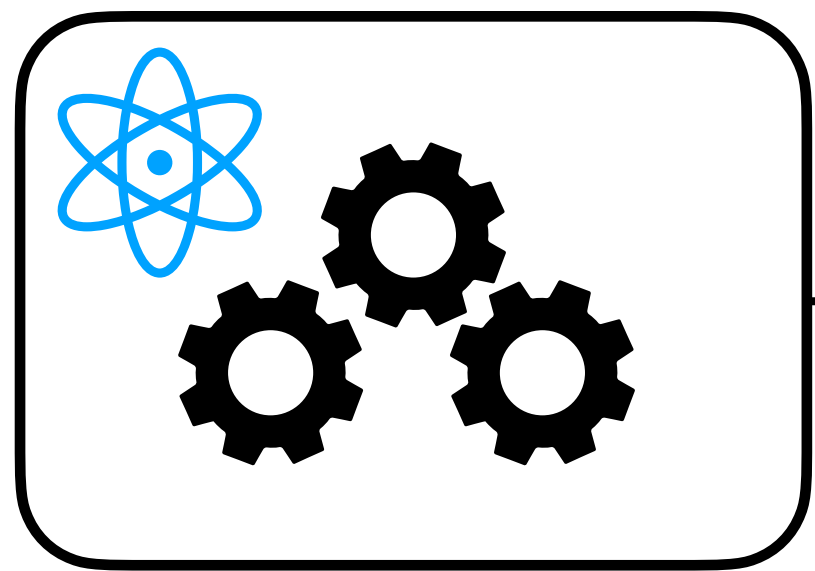
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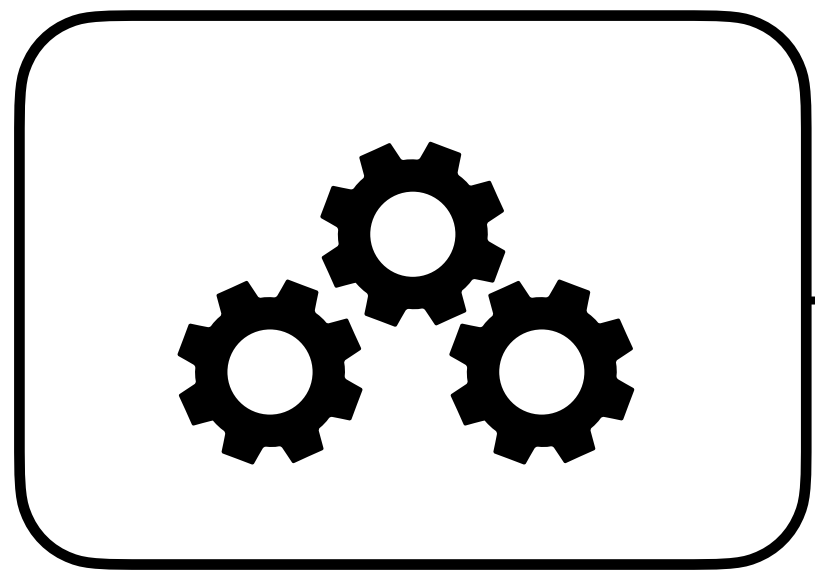
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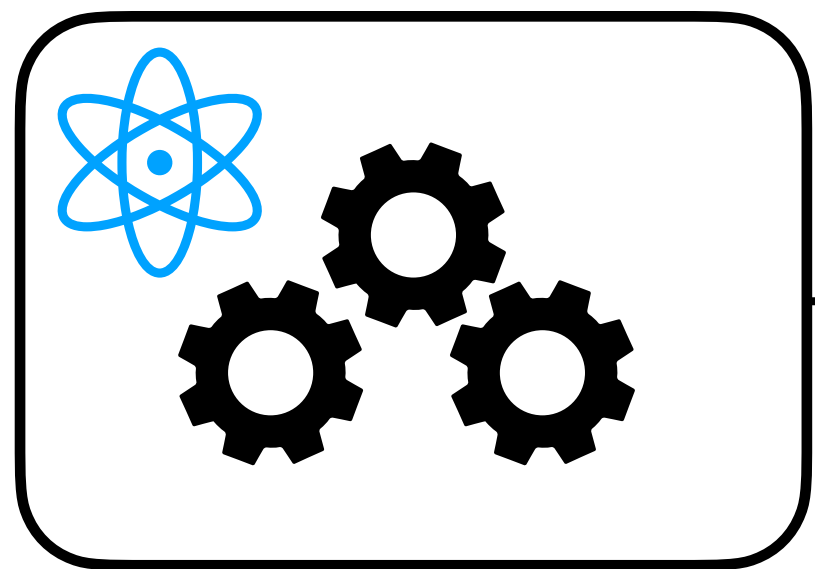
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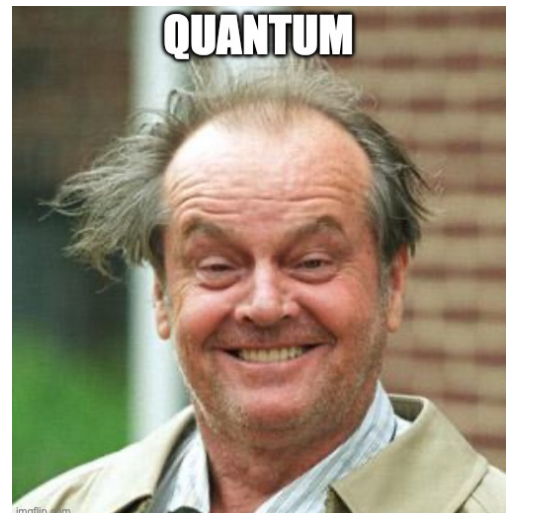
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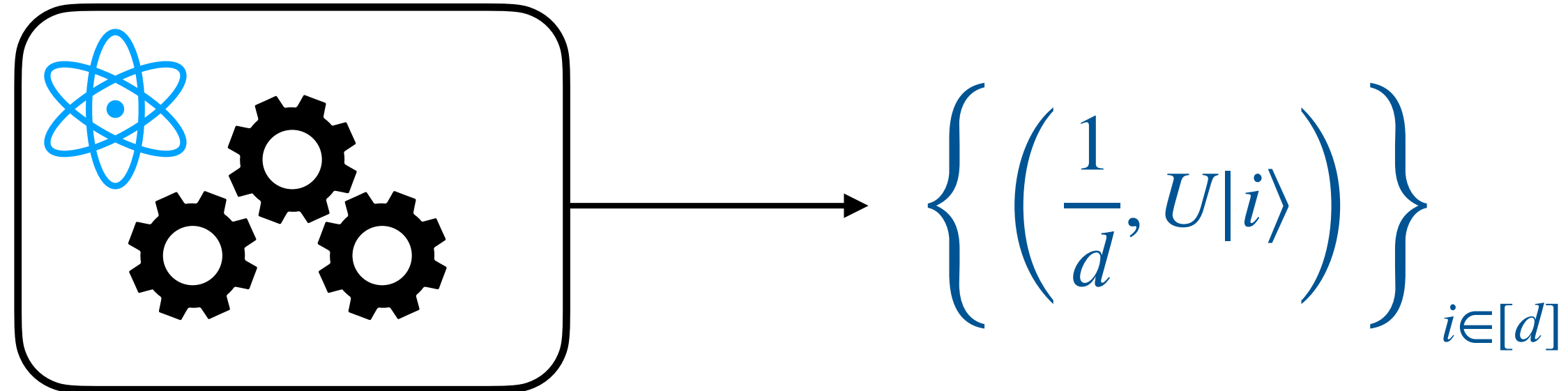
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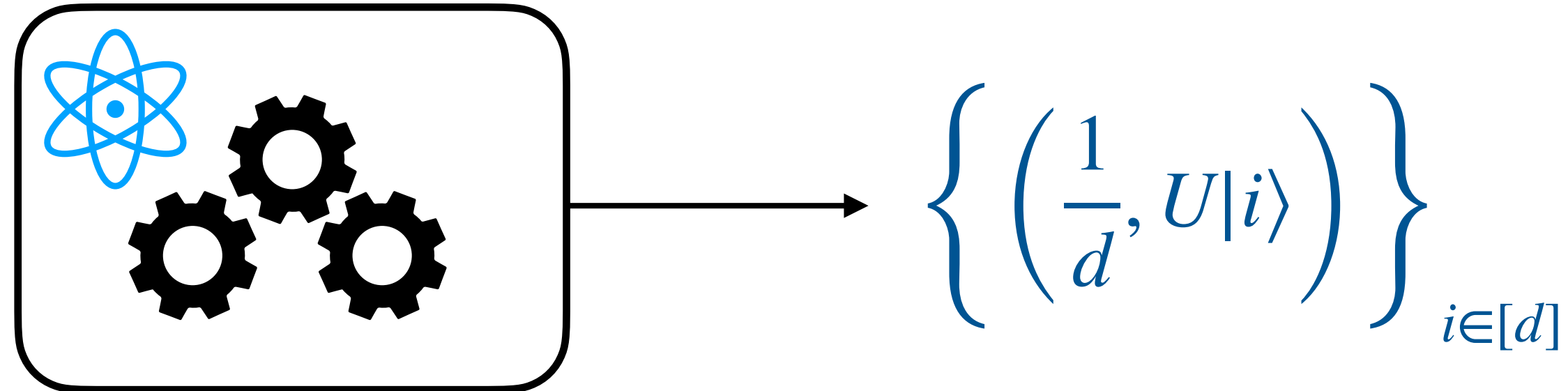


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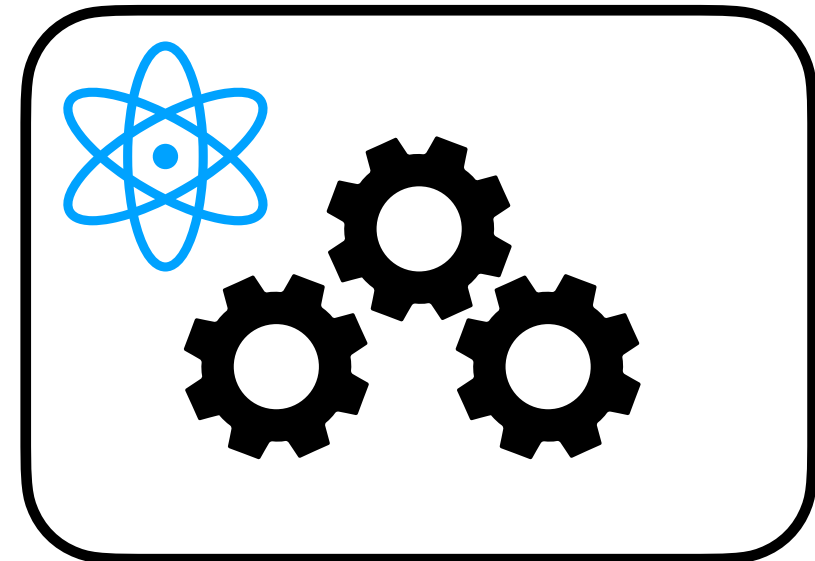


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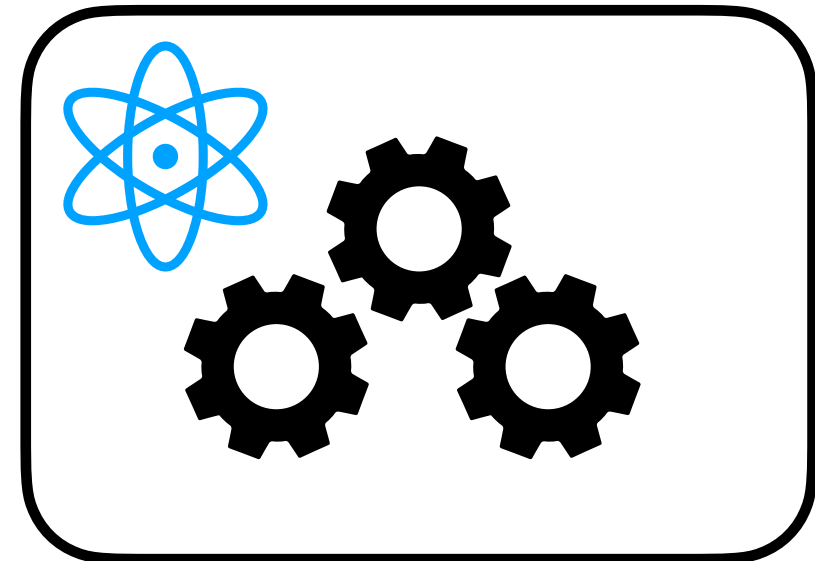
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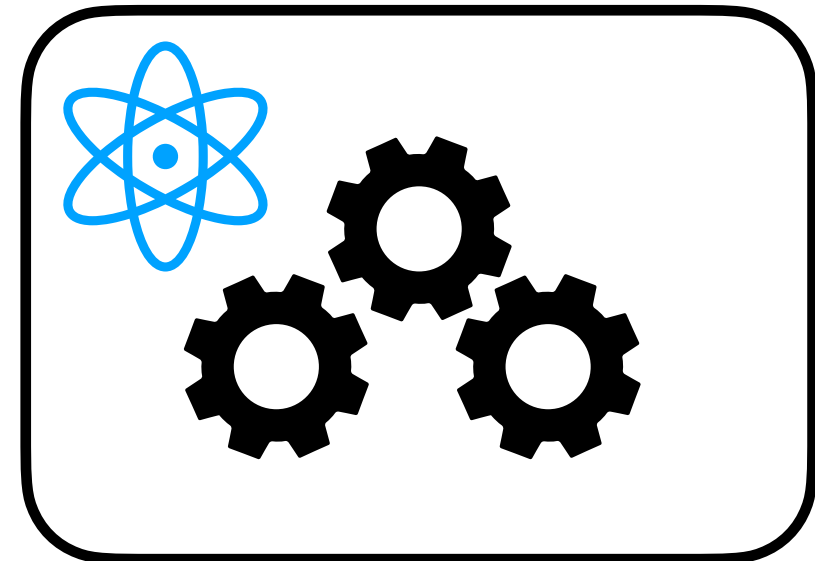
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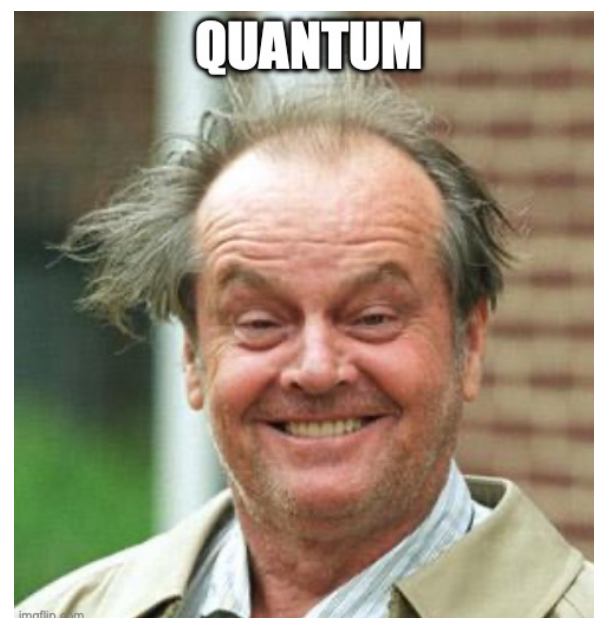
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For this, we will invoke the **Spectral Theorem from Linear Algebra!**

Refresher on the Spectral Theorem

[Spectral Theorem]

Given a normal matrix $A \in \mathbb{C}^{d \times d}$ (i.e., $AA^\dagger = A^\dagger A$), there exist

- (1) Orthonormal basis of eigenvectors $\{|\psi_i\rangle\}_{i \in [d]}$, and**
- (2) Corresponding Eigenvalues $\{\lambda_i\}_{i \in [d]}$ (i.e., $A|\psi_i\rangle = \lambda_i|\psi_i\rangle$), such that**

$$A = \sum_{i=1}^d \lambda_i |\psi_i\rangle \langle \psi_i|$$

Refresher on the Spectral Theorem

In particular, this theorem holds for Hermitian and unitary matrices

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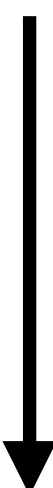
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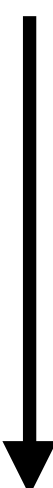
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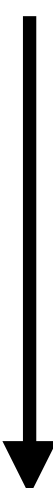
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$$(3) \quad \rho \geq 0 \implies \lambda_i \geq 0, \forall i \in [d]$$


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Unitary Evolution

Given unitary U , evolution is given by

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$$\rho \mapsto_U U\rho U^\dagger \quad \text{(Density Operator)}$$

Measurement

A measurement is defined by a collection $\{M_i = E_i^\dagger E_i \in \mathbb{C}^{d \times d}\}_{i \in [m]}$ of positive semi-definite (PSD) matrices satisfying $\sum_{i=1}^m M_i = I$

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$$(2) \quad \sum_{i=1}^m M_i = I, \text{Tr}(\rho) = 1 \implies \sum_{i=1}^m p_i = \text{Tr}\left(\sum_{i=1}^m M_i \rho\right) = 1 \quad \checkmark$$

Measurement

A measurement is defined by a collection $\{M_i = E_i^\dagger E_i \in \mathbb{C}^{d \times d}\}_{i \in [m]}$ of positive semi-definite (PSD) matrices satisfying $\sum_{i=1}^m M_i = I$

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(we require an explicit factorization $M_i = E_i^\dagger E_i$ since $M_i = E_i^\dagger U^\dagger U E_i$ for any unitary U)

POVM Measurement

The collection $\{M_i \in \mathbb{C}^{d \times d}\}_{i \in [m]}$ of positive semi-definite (PSD) matrices satisfying $\sum_{i=1}^m M_i = I$ is called Positive Operator-Valued Measure (POVM)

(In this case we do not require an explicit factorization $M_i = E_i^\dagger E_i$)

Example of Measurement

Measuring in the computational basis $\{M_x = |x\rangle\langle x| \in \mathbb{C}^{2^n \times 2^n}\}_{x \in \{0,1\}^n}$

Suppose $\rho = |\psi\rangle\langle\psi|$, **where** $|\psi\rangle = \sum_{x \in \{0,1\}^n} \alpha_x |x\rangle$

Then, $p_x = \text{Tr}(M_x \rho) = \text{Tr}(|x\rangle\langle x| \rho) = \langle x | \rho | x \rangle = \langle x | \psi \rangle \langle \psi | x \rangle = |\alpha_x|^2$

Example of Measurement

Measuring Hamming weight parity

$$\left\{ M_{\text{even}} = \sum_{x \in \{0,1\}^n: |x| \equiv 0 \pmod{2}} |x\rangle\langle x|, M_{\text{odd}} = I - M_{\text{even}} \right\}$$

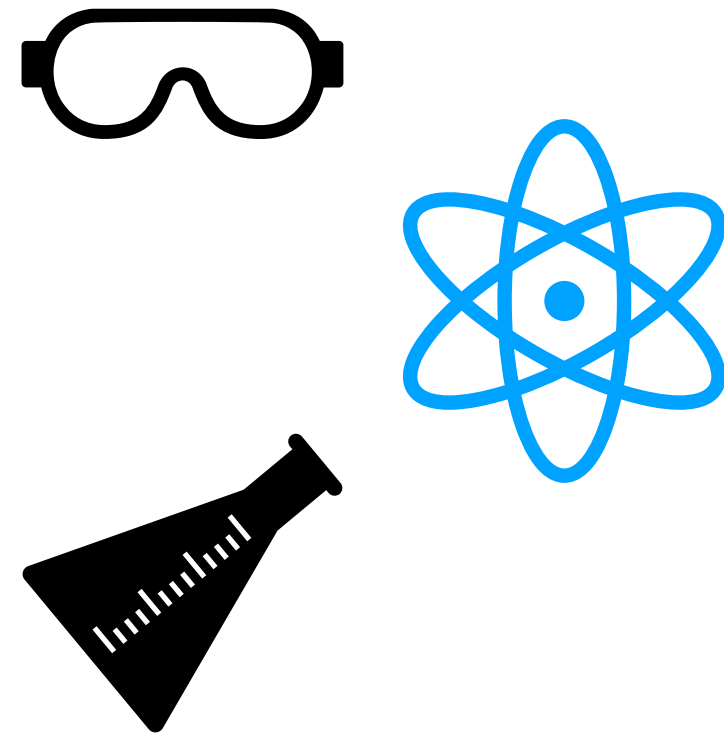
Suppose $\rho = |\psi\rangle\langle\psi|$, **where** $|\psi\rangle = \sum_{x \in \{0,1\}^n} \alpha_x |x\rangle$

$$\text{Then, } p_{\text{even}} = \text{Tr}(M_{\text{even}}\rho) = \sum_{x \in \{0,1\}^n: |x| \equiv 0 \pmod{2}} |\alpha_x|^2$$

How can we learn a quantum state?

Learning Scenario

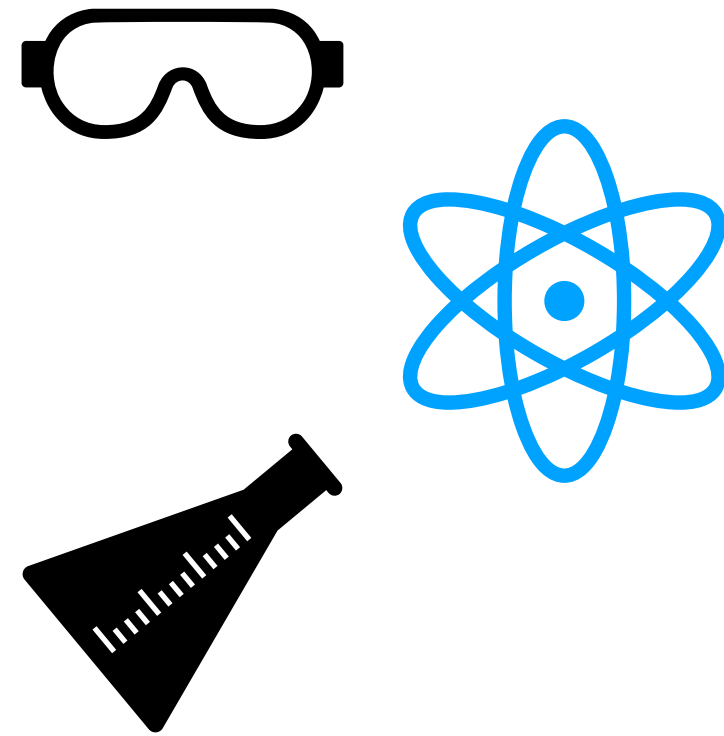
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Learning Scenario

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This task is called Quantum State Tomography

Hilbert-Schmidt Inner Product

Let $A, B \in \mathbb{C}^{d \times d}$. The Hilbert-Schmidt inner product between A and B is defined as

$$\langle A, B \rangle = \mathbf{Tr}(A^\dagger B) = \sum_{i=1}^d \sum_{j=1}^d \overline{A_{i,j}} B_{i,j}$$

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(this is the entrywise inner product of the matrices)

Pauli Matrices

Widely used matrices in Quantum Information Sciences in many contexts

$$X = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$Z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$Y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$$

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$$X = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

bit flip error

$$Z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

phase flip error

$$Y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$$

bit and phase flip errors

Pauli Matrices

Some Properties

$$X = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

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Pauli Matrices

Given an arbitrary Hermitian matrix $H \in \mathbb{C}^{2 \times 2}$, $H = \begin{pmatrix} a & b - ic \\ b + ic & d \end{pmatrix}$, we have

$$H = \frac{(a + d)}{2}I + bX + cY + \frac{(a - d)}{2}Z$$

$$I = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \quad X = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \quad Z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \quad Y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$$

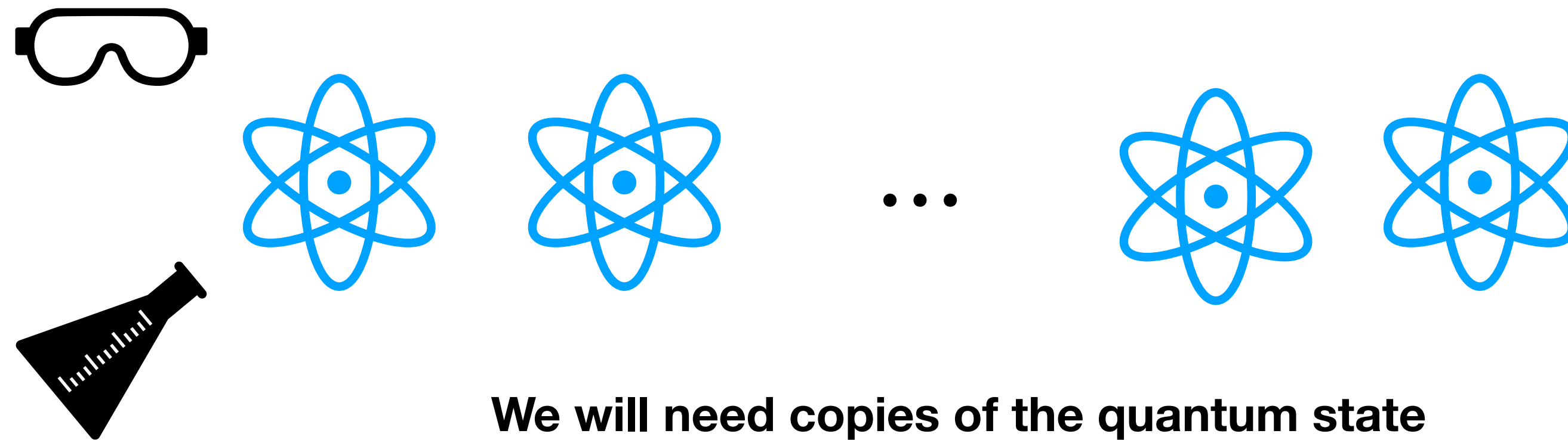
Pauli Matrices

The matrices $\{I, X, Y, Z\}$ form an orthogonal basis for the space of Hermitian matrices in $\mathbb{C}^{2 \times 2}$

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Back to our Learning Scenario

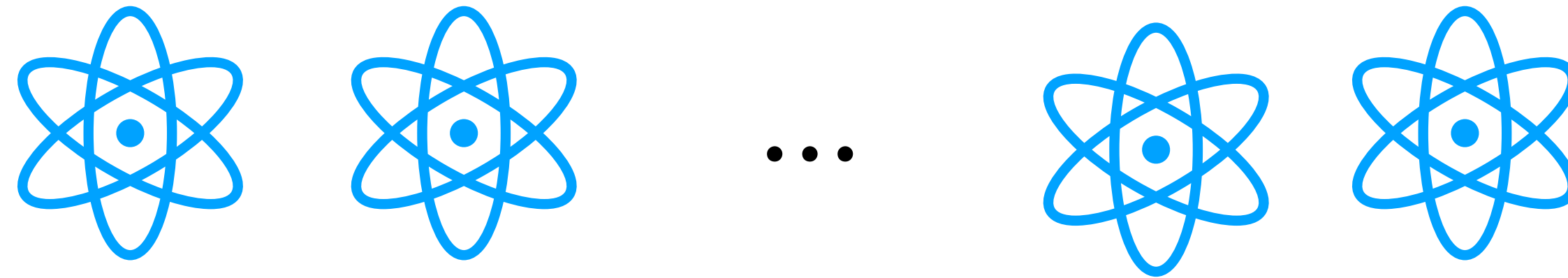
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Back to our Learning Scenario



We will need copies of the quantum state

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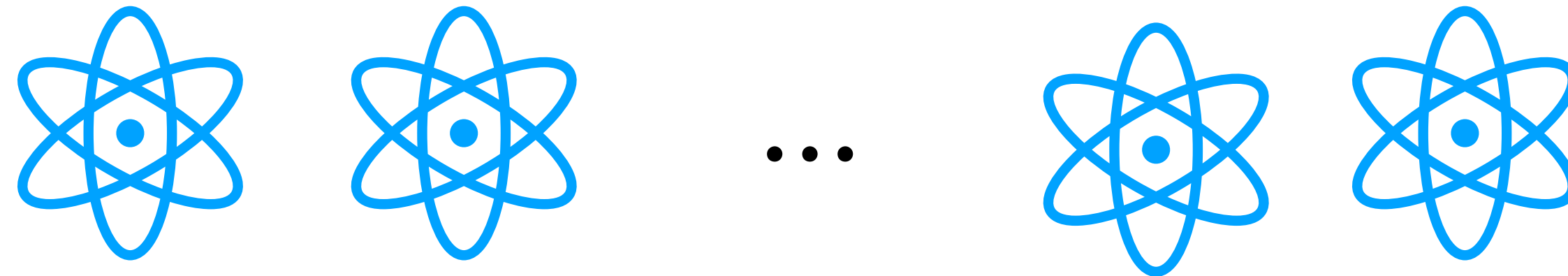
Using these copies, we estimate:

$$a_x = \langle X, \rho \rangle \approx \hat{a}_x$$

$$a_y = \langle Y, \rho \rangle \approx \hat{a}_y$$

$$a_z = \langle Z, \rho \rangle \approx \hat{a}_z$$

Back to our Learning Scenario



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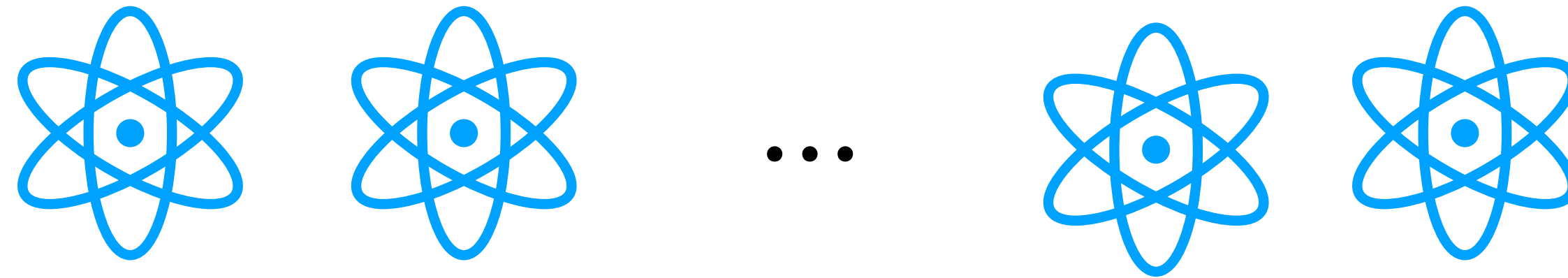
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$$\rho = \frac{I + a_x X + a_y Y + a_z Z}{2}$$

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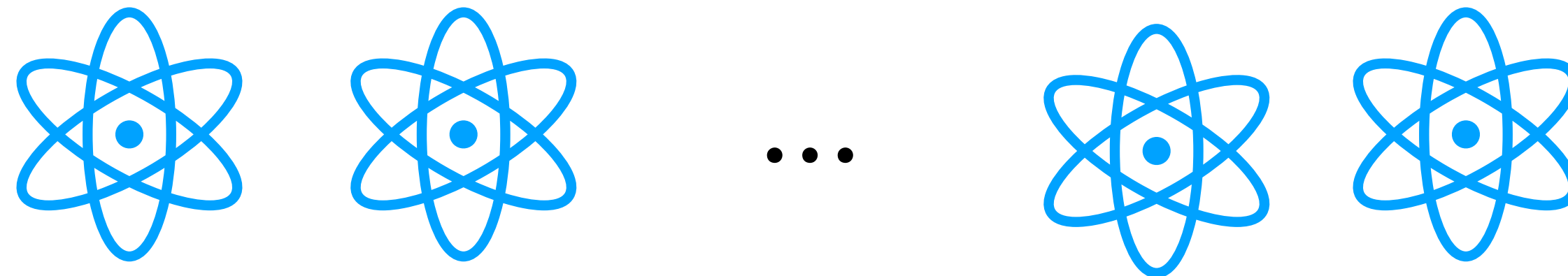
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Output the approximation

$$\rho = \frac{I + a_x X + a_y Y + a_z Z}{2} \approx \hat{\rho} = \frac{I + \hat{a}_x X + \hat{a}_y Y + \hat{a}_z Z}{2}$$

How can we learn an n-qubit state?

What is the quantum state $\rho \in \mathbb{C}^{2^n \times 2^n}$?



We will need maaaaa...any copies of the quantum state for n large!

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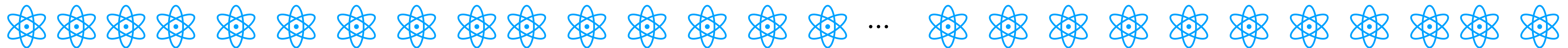
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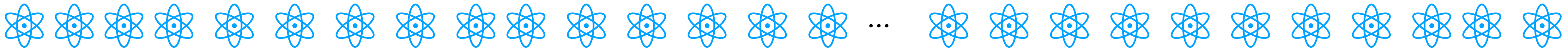
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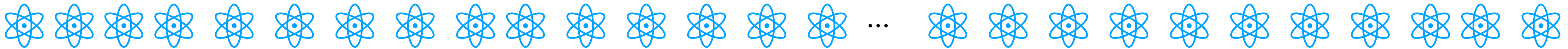


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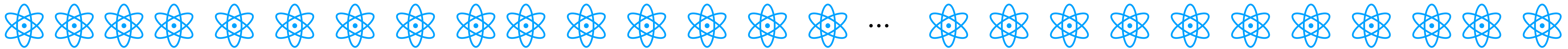
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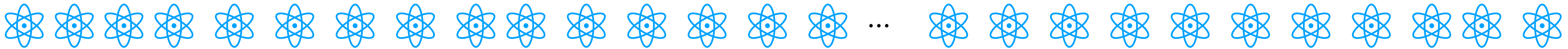
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No! Many notions of cheaper tomography known as Shadow Tomography!

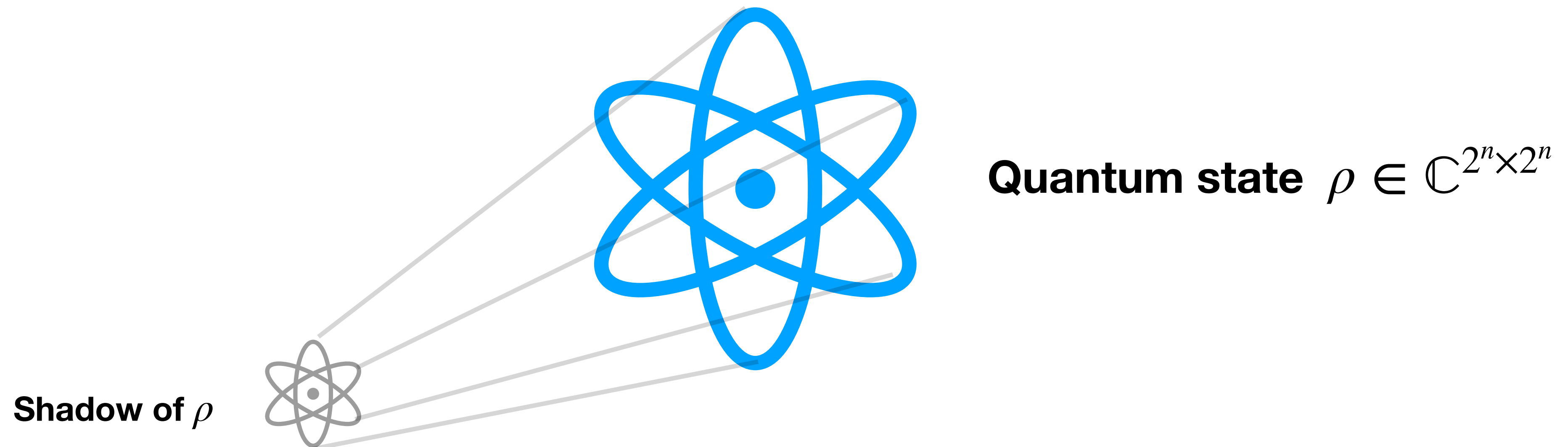
A Glimpse of an Active Research Topic:
The Shadow Tomography Case

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Goal: Given copies of $\rho \in \mathbb{C}^{d \times d}$ and precision parameter $\epsilon > 0$, estimate $\langle M_i, \rho \rangle \pm \epsilon, \forall i \in [m]$

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[Aaronson'16]

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Amazing: In particular, we can learn $4^{\Theta(n)}$ properties (as above) of n -qubit states
with only $n^{O(1)}$ copies!

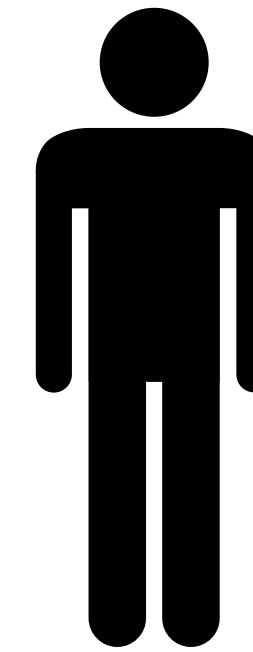
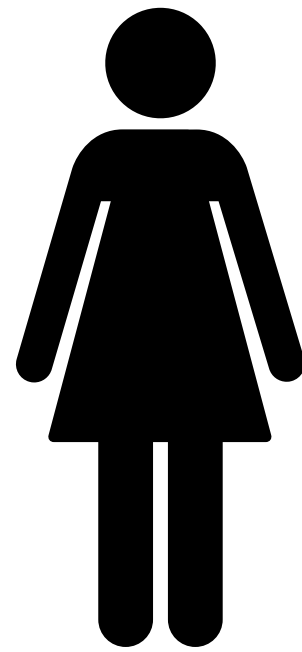
Shadow Tomography

There are many variants of Shadow tomography extending the seminal work of Aaronson with various guarantees

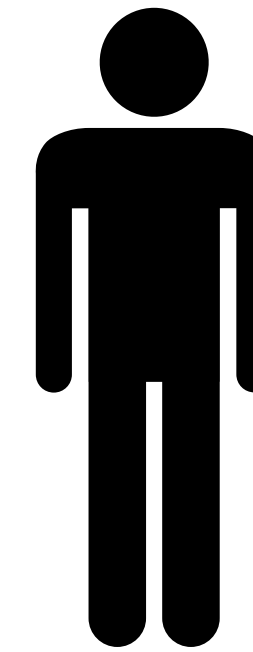
(this is an active research area)

Back to Alice in Scenario 1

Suppose Alice and Bob share an EPR pair

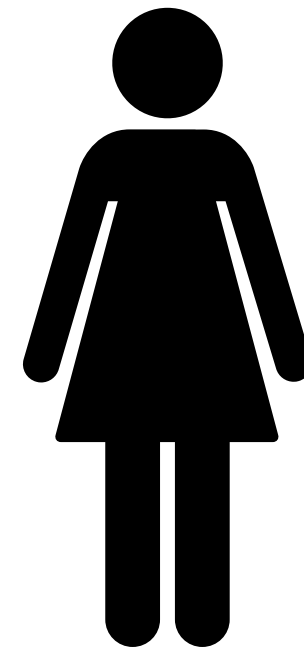


Suppose Alice and Bob share an EPR pair



$$|\mathbf{EPR}\rangle = \frac{1}{\sqrt{2}}|0\rangle^A|0\rangle^B + \frac{1}{\sqrt{2}}|1\rangle^A|1\rangle^B$$

Suppose Alice and Bob share an EPR pair



Goodbye, I have to go

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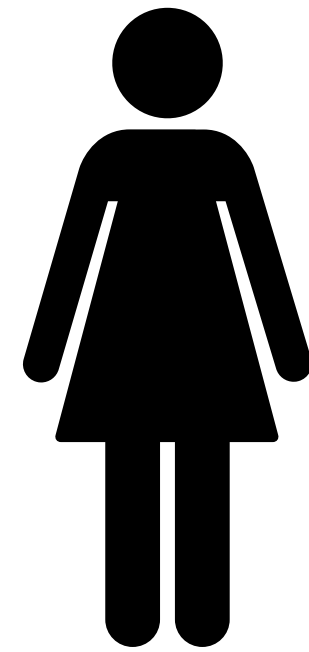


What is the state of Alice's quantum system?

$$|\mathbf{EPR}\rangle = \frac{1}{\sqrt{2}}|0\rangle^A|0\rangle^B + \frac{1}{\sqrt{2}}|1\rangle^A|1\rangle^B \neq |\psi_A\rangle \otimes |\psi_B\rangle$$

(no matter the choice of $|\psi_A\rangle, |\psi_B\rangle$)

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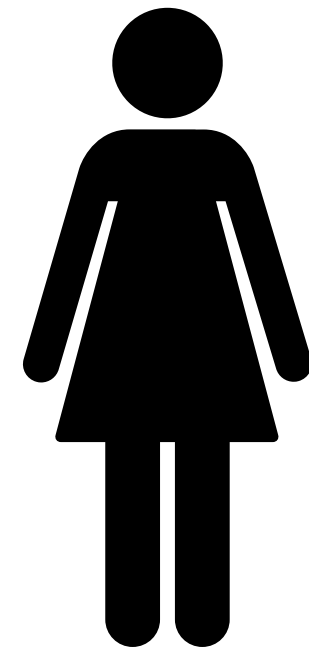
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Suppose Alice and Bob share an EPR pair



What is the state of Alice's quantum system?

It is not a pure state!

It is a mixed state given by a density operator!

$$|\mathbf{EPR}\rangle = \frac{1}{\sqrt{2}}|0\rangle^A|0\rangle^B + \frac{1}{\sqrt{2}}|1\rangle^A|1\rangle^B \neq |\psi_A\rangle \otimes |\psi_B\rangle$$

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**How can we compute the density operator on
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We will need the partial trace operation!

Sub-system State via Partial Trace

Let $\rho_{AB} \in \mathcal{D}(\mathcal{H}_A \otimes \mathcal{H}_B)$ be a density operator on a bipartite tensor space

$\mathcal{H}_A \otimes \mathcal{H}_B$ with orthonormal basis $\{|i\rangle_A\}_{i \in [d_A]}$ and $\{|i\rangle_B\}_{i \in [d_B]}$

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The reduced density operator ρ_A on subsystem A is given by the partial trace $\text{Tr}_B(\rho_{AB})$

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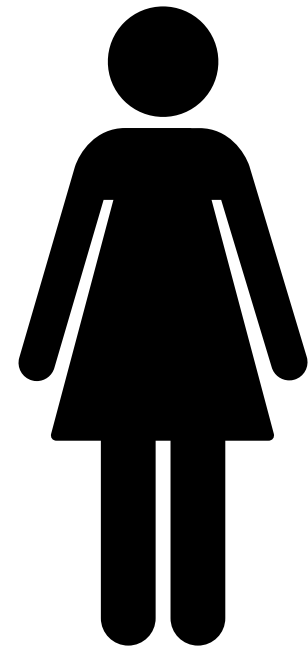
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where the partial trace $\text{Tr}_B(\cdot)$ is the linear operator acting on basis elements as

$$\text{Tr}_B(|i_A\rangle\langle j_A| \otimes |i_B\rangle\langle j_B|) = |i_A\rangle\langle j_A| \text{Tr}(|i_B\rangle\langle j_B|) = |i_A\rangle\langle j_A| \langle j_B|i_B\rangle$$

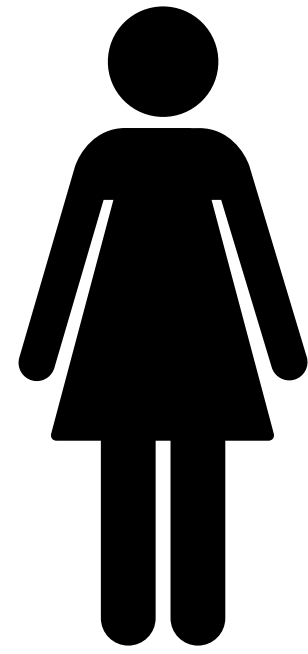
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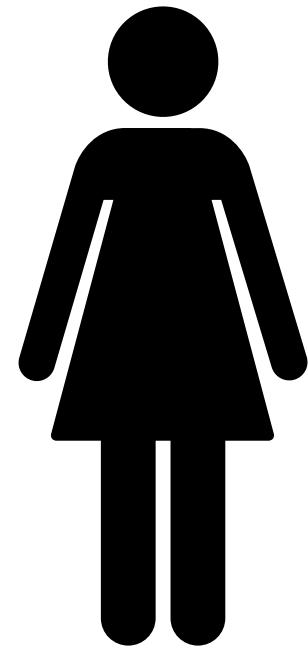
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$$\rho_{AB} = |\mathbf{EPR}\rangle\langle\mathbf{EPR}| = \frac{1}{2}|0\rangle^A|0\rangle^B\langle 0|^A\langle 0|^B + \frac{1}{2}|1\rangle^A|1\rangle^B\langle 0|^A\langle 0|^B + \frac{1}{2}|0\rangle^A|0\rangle^B\langle 1|^A\langle 1|^B + \frac{1}{2}|1\rangle^A|1\rangle^B\langle 1|^A\langle 1|^B$$

$$\rho_A = \mathbf{Tr}_B(\rho_{AB})$$

$$\begin{aligned} &= \frac{1}{2}\mathbf{Tr}_B(|0\rangle^A|0\rangle^B\langle 0|^A\langle 0|^B) + \frac{1}{2}\mathbf{Tr}_B(|1\rangle^A|1\rangle^B\langle 0|^A\langle 0|^B) \\ &\quad + \frac{1}{2}\mathbf{Tr}_B(|0\rangle^A|0\rangle^B\langle 1|^A\langle 1|^B) + \frac{1}{2}\mathbf{Tr}_B(|1\rangle^A|1\rangle^B\langle 1|^A\langle 1|^B) \\ &= \frac{1}{2}|0\rangle^A\langle 0|^A + \frac{1}{2}|1\rangle^A\langle 1|^A \end{aligned}$$

Alice's State via Partial Trace



$$|\mathbf{EPR}\rangle = \frac{1}{\sqrt{2}}|0\rangle^A|0\rangle^B + \frac{1}{\sqrt{2}}|1\rangle^A|1\rangle^B$$

$$\rho_{AB} = |\mathbf{EPR}\rangle\langle\mathbf{EPR}| = \frac{1}{2}|0\rangle^A|0\rangle^B\langle 0|^A\langle 0|^B + \frac{1}{2}|1\rangle^A|1\rangle^B\langle 0|^A\langle 0|^B + \frac{1}{2}|0\rangle^A|0\rangle^B\langle 1|^A\langle 1|^B + \frac{1}{2}|1\rangle^A|1\rangle^B\langle 1|^A\langle 1|^B$$

$$\rho_A = \mathbf{Tr}_B(\rho_{AB})$$

$$= \frac{1}{2}\mathbf{Tr}_B(|0\rangle^A|0\rangle^B\langle 0|^A\langle 0|^B) + \frac{1}{2}\mathbf{Tr}_B(|1\rangle^A|1\rangle^B\langle 0|^A\langle 0|^B)$$

$$+ \frac{1}{2}\mathbf{Tr}_B(|0\rangle^A|0\rangle^B\langle 1|^A\langle 1|^B) + \frac{1}{2}\mathbf{Tr}_B(|1\rangle^A|1\rangle^B\langle 1|^A\langle 1|^B)$$

$$= \frac{1}{2}|0\rangle^A\langle 0|^A + \frac{1}{2}|1\rangle^A\langle 1|^A$$

Maximally mixed state!

Due to *entanglement* between Alice and Bob,
Alice reduced state is mixed!

What is quantum *entanglement*?

What is quantum *entanglement*?

(We will focus on bipartite entanglement of pure quantum states)

Thank you!

Thank you!

More Questions?