

Lecture 1

Welcome to Phys 598 GTC

Goals for the course:

- Develop theory background for research in topological materials
- Understand modern developments in the study of noninteracting electrons
- Develop a foundation for understanding group theory in solid state physics

Guide to topics:

- ① Space group symmetries
- ② Wannier functions and band representations
- ③ Berry phases and band topology
- ④ Topological crystalline insulators

Course website: courses.physics.illinois.edu/phys598g/c

Course components: Lectures

HWs (5) \rightarrow graded on completeness
submitted via gradescope

Final presentations

Office hours: 4-5pm Mondays via Zoom link
on course website

I. Review/Intro to Group theory

Useful resources:

- Dresselhaus "Applications of Group Theory to the Physics of Solids"

• Bradley & Cracknell "Mathematical Theory of Symmetry in Solids"

• Serre "Linear Representations of Finite Groups"

Starting point: Hamiltonian $H = \frac{p^2}{2m} + V(x) + \dots$

Schrödinger Eqn: $H|\psi\rangle = E|\psi\rangle$

Find the set of transformations
 $\vec{x} \rightarrow \vec{x}'$

$$\vec{p} \rightarrow \vec{p}'$$

$$|\psi\rangle \rightarrow |\psi'\rangle$$

that leave the Schrödinger equation unchanged

$$H \rightarrow H' = H$$

$$\vec{x}' = R \vec{x} + \vec{d}$$

$$\vec{p}' = R \vec{p}$$

R - 3x3 rotation or reflection matrix $R^T = R^{-1}$

\vec{d} - translation vector

Intuitive facts: ① if I have two transformations, I can do one and then the other to get

a third transformation

(2) I can always undo a transformation

(3) $\begin{matrix} \vec{x} \rightarrow \vec{x} \\ \vec{p} \rightarrow \vec{p} \end{matrix}$ is a transformation
- identity transformation

Definition: a set G is called a group if

(1) there's some binary operation \cdot such that if $g_1 \in G, g_2 \in G$, then $g_1 \cdot g_2 \in G$, and
 $g \cdot (g_2 \cdot g_3) = (g_1 \cdot g_2) \cdot g_3$

③ $E \in G$ s.t. for all $g \in G$ $E \cdot g = g \cdot E = g$
 E is the identity

② If $g \in G$, then there exists g^{-1} such
that $g \cdot g^{-1} = g^{-1} \cdot g = E$

Examples of Groups: ① The set of unitary operators
on Hilbert space (d -dimensional)
 $U(d)$

- The binary operation is matrix multiplication
- $E = d \times d$ identity matrix
- $V \in U(d), V^T \in U(d)$
 $V^T V = E$
- $V_1 \in U(d), V_2 \in U(d)$
 $(V_1 V_2)^T = V_2^T V_1^T = (V_1 V_2)^{-1}$
 $\Rightarrow V_1 V_2 \in U(d)$

② The group of rotations in 3-dimensions

"special orthogonal group" $SO(3)$

3x3 matrices, determinant 1, transpose is their inverse

③ Translations in 3D space \mathbb{R}^3

- elements are vectors

- binary operation is addition of vectors

- $\vec{v} \in \mathbb{R}^3$ defines the transformation $\vec{x} \rightarrow \vec{x} + \vec{v}$

- identity element: $\vec{0}$

- inverses: $\vec{v}^{-1} = -\vec{v}$

Some important facts about groups

- Given a group G , we can consider subsets $H \subset G$ that are also groups $\leftarrow H$ is a subgroup of G

$H \subset G$ is a subgroup if:

1. $E \in H$

2. H is closed under multiplication

3. H is closed under taking inverses: $h \in H \Rightarrow h^{-1} \in H$

Examples: Consider $SO(3)$. Consider $SO(2) \subset SO(3)$

consisting of all rotations about a fixed axis
 \hat{n}

- translation group $\mathbb{R}^3 = \{(x, y, z), x, y, z \text{ real}\}$
pick 3 linearly independent vectors $\vec{t}_1, \vec{t}_2, \vec{t}_3$

$$T = \{n\vec{t}_1 + m\vec{t}_2 + l\vec{t}_3, n, m, l \in \mathbb{Z}\}$$

$T \subset \mathbb{R}^3$ is a subgroup known
as a Bravais lattice

We can use subgroups HCG to learn about the structure
of G

given a group G and a subgroup H we can define

right cosets $Hg = \{h \cdot g \mid h \in H\}$

$g \in G$

Important fact: every element $g' \in G$ is in exactly one right coset of H

proof: first: $E \in H$

$$Hg' = \{hg' \mid h \in H\} \ni Eg' = g'$$

so g' is in at least one right coset Hg'

to show it is only in one right coset; need to
show $g' \in Hg_1$ and $g' \in Hg_2 \Rightarrow Hg_1 = Hg_2$

$$\begin{array}{ccc} \downarrow & & \downarrow \\ h_1 g_1 = g' & & h_2 g_2 = g' \end{array}$$

$$h_1 g_1 = h_2 g_2$$

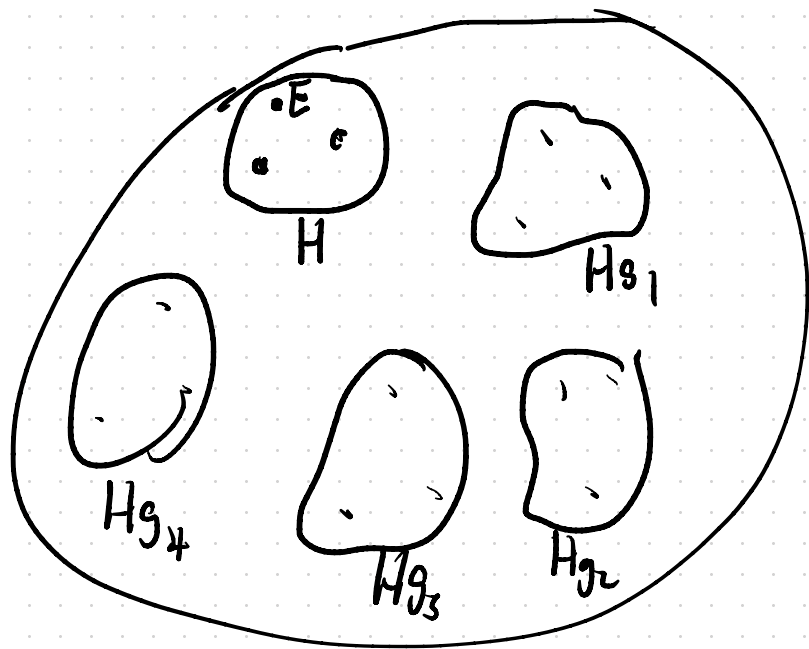
$$\Rightarrow h_2^{-1} h_1 g_1 = h_2^{-1} h_2 g_2$$

$$g_2 = h_2^{-1} h_1 g_1$$

$$g_2 g_1^{-1} = h_2^{-1} h_1 \in H$$

$$Hg_2g_1^{-1} = H$$

$$\Rightarrow Hg_2 = Hg_1$$



Right cosets of H partition G

G

$$G = \overset{E}{\downarrow} H \cup Hg_1 \cup Hg_2 \dots \cup Hg_{n-1} \leftarrow n \text{ right cosets of } H$$

n is known as the index of H in G $|G:H|$

$\{E, g_1, g_2, \dots, g_{n-1}\}$ are coset representatives of H