

Lecture 4

Recap: Representation ρ of a group G

- a vector space V
- a group homomorphism $\rho: G \rightarrow U(V)$
to unitary operators on V

Let G be a group, $\rho: G \rightarrow U(V)$ a representation on a vector space V

$$g \in G \quad \rho(g) \in U(V)$$

Given a vector $|v\rangle \in V$ $\rho(g)|v\rangle \in V$

We can look for $W \subseteq V$ such that

$$\rho(g)|w\rangle \in W \text{ for all } g \in G, |w\rangle \in W$$

such a subspace is called an invariant subspace

given $W \subseteq V$, we can consider W^\perp - orthogonal complement of W :

$$W^\perp = \left\{ |w_\perp\rangle \in V \mid \langle w_\perp | w \rangle = 0 \text{ for all } |w\rangle \in W \right\}$$

$$V = W \oplus W^\perp$$

every $|v\rangle \in V$ can be written uniquely as $|v\rangle = |w\rangle + |w_\perp\rangle$

Since $\rho(g)$ are all unitary, W^\perp is also an invariant subspace

p.f. let $|w\rangle \in W$, $|w^\perp\rangle \in W^\perp$

W invariant, $\rho(g)|w\rangle \in W$ for any $g \in G$

$$\Rightarrow \langle w^\perp | (\rho(g)|w\rangle) = 0$$

$$= (\langle w | (\rho(g)^\dagger |w^\perp\rangle)^\dagger = 0$$

$$\Rightarrow \langle w | \rho(g^{-1}) |w^\perp\rangle = 0$$

for any $g \in G$

for any $|w\rangle \in W$

for any $|w^\perp\rangle \in W^\perp$

$$\rightarrow \rho(g^{-1})|w^\perp\rangle \in W^\perp$$

$\Rightarrow W^\perp$ is an invariant subspace

orthonormal basis for V : $\{|v_1\rangle, |v_2\rangle, |v_3\rangle, \dots\}$

$\rho(g)$ as a matrix in this basis

$$\rho(g) = \begin{pmatrix} \langle v_1 | \rho(g) | v_1 \rangle & \langle v_1 | \rho(g) | v_2 \rangle & \dots \\ \langle v_2 | \rho(g) | v_1 \rangle & \langle v_2 | \rho(g) | v_2 \rangle & \dots \\ \vdots & \vdots & \ddots \end{pmatrix}$$

$\rho: G \rightarrow U(V)$ is a homomorphism

$$\rho(gg^{-1}) = \rho(g)\rho(g^{-1}) = E$$

$$= \rho(g)^\dagger \rho(g) \rho(g^{-1}) = \rho(g)^\dagger$$

$$\Rightarrow \rho(g^{-1}) = \rho(g)^\dagger$$

But instead, we can pick a different basis

$\{|w_1\rangle, |w_2\rangle, |w_3\rangle, \dots\}$ be a orthonormal basis for W

$\{|w_1^\perp\rangle, |w_2^\perp\rangle, |w_3^\perp\rangle, \dots\}$ be an orthonormal basis for W^\perp

a basis for V $\{|w_1\rangle, |w_2\rangle, \dots, |w_1^\perp\rangle, |w_2^\perp\rangle, \dots\}$

$$V = W \oplus W^\perp$$

$$|v\rangle = \underbrace{\sum_i |w_i\rangle \langle w_i | v \rangle}_{W} + \underbrace{\sum_j |w_j^\perp\rangle \langle w_j^\perp | v \rangle}_{W^\perp}$$

in this basis

$$\rho(\rho) = \begin{pmatrix} \langle w_i | \rho(s) | w_j \rangle & \langle w_i | \rho(s) | w_j^\perp \rangle \\ \langle w_i^\perp | \rho(s) | w_j \rangle & \langle w_i^\perp | \rho(s) | w_j^\perp \rangle \end{pmatrix} \begin{matrix} w \\ w^\perp \end{matrix} = 0 = \begin{pmatrix} \rho_w(s) & 0 \\ 0 & \rho_{w^\perp}(s) \end{pmatrix}$$

ρ_w and ρ_{w^\perp} are themselves representations

$$\rho \cong \rho_w \oplus \rho_{w^\perp}$$

ρ_w, ρ_{w^\perp} are subrepresentations of ρ

ρ is a reducible representation

up to a change of basis
 "unitarily equivalent to"

$$\rho_1 \cong \rho_2 \text{ if } \exists U \text{ s.t. } U \rho_1(s) U^\dagger = \rho_2(s)$$

$\forall g$

Any representation that is not reducible is called irreducible

Note if ρ is irreducible then the only invariant subspaces of ρ are $\{\vec{0}\}$ and V

Trivial example: let $[G]$ be any group, and $V = \mathbb{C}$
vector space of complex numbers

$$U(\mathbb{C}) = \{e^{i\phi} / \phi \in (0, 2\pi)\} = U(1)$$

$$\rho: g \rightarrow \rho(g) = \mathbb{1} = e^{i0}$$

$$\rho(g_1 g_2) = \mathbb{1} = \rho(g_1) \rho(g_2) = \mathbb{1} \times \mathbb{1}$$

trivial representation - irreducible

Nontrivial example: $SU(2)$ we know $V_{\frac{1}{2}} = \{|\uparrow\rangle, |\downarrow\rangle\}$
 $\rho_{\frac{1}{2}}: (\hat{n}, \theta) \rightarrow e^{-i\frac{\theta}{2}\vec{\sigma}\cdot\hat{n}}$

On the vector space of two spin- $\frac{1}{2}$ particles

$$V = V_{\frac{1}{2}} \otimes V_{\frac{1}{2}} = \{|\uparrow\uparrow\rangle, |\downarrow\uparrow\rangle, |\uparrow\downarrow\rangle, |\downarrow\downarrow\rangle\}$$

$$\rho_{\frac{1}{2} \times \frac{1}{2}}: (\hat{n}, \theta) \rightarrow e^{-i\frac{\theta}{2}\vec{\sigma}_1 \cdot \hat{n}} \otimes e^{-i\frac{\theta}{2}\vec{\sigma}_2 \cdot \hat{n}}$$

Are there (nontrivial) invariant subspaces for $\rho_{\frac{1}{2} \times \frac{1}{2}}$?

$$\left\{ \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle) \right\} = W \quad \text{-singlet state}$$

invariant subspace

$$W^\perp = \left\{ |\uparrow\uparrow\rangle, \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle), |\downarrow\downarrow\rangle \right\} \quad \text{-invariant}$$

subspace: triplet space

in this basis

$$\rho_{\frac{1}{2} \times \frac{1}{2}}(\hat{n}, \theta) = \left(\begin{array}{c|c} 1 & \vec{0} \\ \hline \vec{0} & e^{-i\hat{n} \cdot \vec{L}\theta} \end{array} \right)$$

spin-1 representation

$$\rho_{\frac{1}{2} \times \frac{1}{2}} \cong \rho_0 \oplus \rho_1 \quad \left(\frac{1}{2} \otimes \frac{1}{2} = 0 \oplus 1 \right)$$

Clebsch-Gordan coeffs: matrix elements of the unitary matrix that block-diagonalize

Schur's Lemma (2 $\frac{1}{2}$ parts)

Alternatively $\rho: G \rightarrow GL(V)$

"Weyl unitarity trick" $\langle v | w \rangle$ - Hermitian inner product

$$\langle \langle v | w \rangle \rangle = \sum_{g \in G} \langle \rho(g)v | \rho(g)w \rangle$$

Schur's Lemma part 1: let G be a group

$$\rho_1: G \rightarrow U(V_1)$$

$$\rho_2: G \rightarrow U(V_2)$$

two irreducible representations (irreps) of G

and a matrix $H: V_1 \rightarrow V_2$

$$\forall v_1 \in V_1$$

$$\forall v_2 \in V_2$$

$$H \forall v_1 \in V_1 \rightarrow v_2 \in V_2$$

IF $H \rho_1(g) = \rho_2(g) H$ for all $g \in G$

then either: ① $H = 0$

② H is invertible

PF: if $H = 0$ $\ker H = V_1$
 $\text{Im } H = \{ \vec{0} \}$

if H is invertible $\text{Ker } H = \{\vec{0}\}$

$$\text{Im } H = V_2$$

Lets look at $\text{Ker } H = \{ |v\rangle \in V_1 \mid H|v\rangle = 0 \}$

pick $|v\rangle \in \text{Ker } H$

$$0 = \rho(g) H |v\rangle = H(\rho_1(g) |v\rangle)$$

$\Rightarrow |v\rangle \in \text{Ker } H$ then so is $\rho_1(g) |v\rangle$ for all $g \in G$

$\Rightarrow \text{Ker } H$ is an invariant subspace of V_1

but ρ_1 is irreducible

$$\text{Ker } H = \begin{cases} \{0\} & - H \text{ is one-to-one} \\ V_1 & - H = 0 \end{cases}$$

Now let's look at $\text{Im } H = \{ |v_2\rangle \in V_2 \mid |v_2\rangle = H|v_1\rangle \text{ for some } |v_1\rangle \in V_1 \}$

$$|w\rangle \in \text{Im } H - |w\rangle = H|v\rangle \quad |v\rangle \in V_1$$

$$P_2(\rho)|w\rangle = P_2(\rho)H|v\rangle = H P_1(\rho)|v\rangle$$

$\text{Im } H \stackrel{\cap}{\rightarrow} \text{Im } H$ is an invariant subspace of P_2

$$\text{Im } H = \begin{cases} V_2 & : H \text{ is "onto"} \\ \vec{0} & : H = \vec{0} \end{cases}$$

Either $H = \vec{0}$ or

H is one-to-one and onto $\Rightarrow H$ is
invertible

Part 2: Suppose $V_1 = V_2 = V$ and $\rho_1 = \rho_2 = \rho$
are the same, ρ irreducible and finite-dimensional
and assume we have H that satisfies part 1 assumption

$$H\rho(s) = \rho(s)H \Rightarrow [H, \rho(s)] = 0$$

then either: $H = 0$ or

$$H = \lambda \text{Id}_V \leftarrow \text{identity matrix}$$

(if ρ is an irrep, the only matrix that commutes with every $\rho(g)$ is the identity)

Pf: Part 1 says $H=0$ or H invertible

Assume H invertible $\rightarrow H$ a finite dim square matrix

$\Rightarrow H$ has at least one eigenvector $|v\rangle$ w/ eigenvalue

$$\left\{ |n\rangle \mid n=0,1,2,\dots \right\} \quad \lambda$$

a^\dagger

$$B = H - \lambda \text{Id}_V \quad \rightarrow \quad [B, \rho(g)] = 0$$

$$\Rightarrow B = 0 \quad \text{or} \quad B \text{ is invertible}$$

\uparrow
 $B|v\rangle = 0$ so B not
invertible

$$B=0 \Rightarrow H = \lambda \text{Id}_V$$

If G is a symmetry group H

$$[H, \rho(g)] = 0$$

$$\{|\psi_i\rangle, i=1, \dots, N\}$$

$$\rho(g)|\psi_i\rangle = \sum_{j=1}^N |\psi_j\rangle \langle \psi_j | \rho(g) | \psi_i \rangle$$

$$\rho_{j,i}(g)$$

$$[H]_{ij} = \langle \Psi_i | H | \Psi_j \rangle$$

$$\sum_k [H]_{ik} \rho_{kj}(\mathcal{S}) - \rho_{ik}(\mathcal{S}) [H]_{kj} = 0$$

$$[H]_{ij} = E_i \delta_{ij}$$

\rightarrow

States transforming in an irrep of the symmetry group are degenerate

Ex: Nonrelativistic Hydrogen atom

fixed n , $\{ |nlm_z\rangle \mid l=0, \dots, n \quad m_z = -l, \dots, l \}$

$$V_1 = \{ |n l_1 m_z\rangle \} \leftarrow \text{spin } l_1 \text{ irrep of } SO(3)$$

$$V_2 = \{ |n l_2 m'_z\rangle \} \leftarrow \text{spin } l_2 \text{ irrep } SO(3)$$

$l_1 \neq l_2$ different dimensions

$$\text{Part I: } \langle n l_1 m_z | H | n l_2 m'_z \rangle = \delta_{l_1 l_2}$$

$$\text{Part II: } \langle n l_1 m_z | H | n l_2 m'_z \rangle = E \delta_{m_z m'_z}$$

Part 2.5.

Let G be a group

$$\rho_1: G \rightarrow U(V_1)$$

$$\rho_2: G \rightarrow U(V_2)$$

Finite dim irreps

$$\rho_2(s)H = H\rho_1(s)$$

H invertible $\rightarrow \rho_1 \sim \rho_2$

rep of
 $SU(2)$

$$\mathfrak{so}(3)$$

representation of
 $SO(3)$

2×2 unitary matrices
w/ determinant 1

$$\mathfrak{so}(3) + \frac{1}{2}$$

$$e^{-i\frac{\theta}{2}\vec{\sigma}\cdot\hat{n}}$$