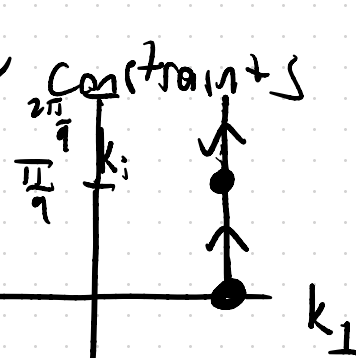


Lecture 22

Last time: Inversion symmetry
on Wilson loop

for $k_{\perp} = -k_{\perp} + \vec{b}_{\perp}$ $\vec{b}_{\perp} \in \Gamma$



$$W_{\frac{2\pi}{a} \sigma_0}(k_{\perp}) = \rho_{(0, k_{\perp})}(\mathbb{I}) W_{\frac{\pi}{a} \sigma_0}^{\dagger}(k_{\perp}) \rho_{(\frac{\pi}{a}, k_{\perp})}(\mathbb{I}) W_{\frac{\pi}{a} \sigma_0}(k_{\perp})$$

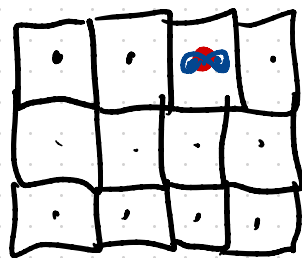
$$\rho_k^{nm}(\mathbb{I}) = \vec{u}_{nk}^{\dagger} \cdot V(\mathbb{I}k - k) B(\mathbb{I}) \vec{u}_{mk}$$

\mathbb{I} , the group representation matrix for \mathbb{I}

$$\det W = e^{i \sum_{\sigma=1}^{n_{occ}} \varphi_{\sigma}(k_{\perp})} = e^{i\pi \left(\Lambda_{-}^{(0, k_{\perp})} + \Lambda_{-}^{(\frac{\pi}{a}, k_{\perp})} \right)}$$

N_-^k - # of negative eigenvalues of $\rho_k(I)$

Lets extend this to 2D $\vec{k} \rightarrow (k_x, k_y)$



two band NN
tight binding model

1D

$$h(k_x) = (\Delta + t_1 \cos k_x a) \sigma_z + t_2 \sin k_x \sigma_y$$

↓ 2D

$$h(k_x, k_y) = (\Delta + t_1 \cos k_x a) \sigma_z + t_2 \sin k_x \sigma_y$$

This is boring

Two phases of the 1D chain

① $\Delta > t_1 > 0$ $\varphi = 0$

② $0 < \Delta < t_1$ $\varphi = \pi$

$h(k_x, k_y=0)$ — in phase ①

$h(k_x, k_y=\frac{\pi}{b})$ — in phase ②

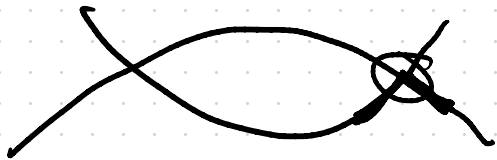
$\Delta > t_1 > 0$

$$h(k_x, k_y) = (\Delta + \Delta \cos k_y b + t_1 \cos k_x a) \sigma_z + t_2 \sin k_x a \sigma_y$$

$$h(k_x, k_y=0) = (2\Delta + t_1 \cos k_x a) \sigma_z + t_2 \sin k_x a \sigma_y \quad \text{①}$$

$$h(k_x, k_y=\frac{\pi}{b}) = t_1 \cos k_x a \sigma_z + t_2 \sin k_x a \sigma_y \quad \text{②}$$

$$E_{\pm} = \pm \sqrt{(\Delta + \Delta \cos k_y b + t_1 \cos k_x a)^2 + t_2^2 \sin^2 k_x a}$$



$$k_x = \frac{\pi}{a}$$

$$\cos k_y b = 1 - \frac{t_1}{\Delta} \rightarrow E_+ = E_- = 0$$

To fix this add an extra term to h

$$+ f(\vec{k}) \sigma_x$$

$$E_{\pm} = \pm \sqrt{(\Delta + \Delta \cos k_y b + t_1 \cos k_x a)^2 + t_2^2 \sin^2 k_x a + f(k)^2}$$

$$B(I) = \sigma_z$$

$$B(T) = \sigma_0 \mathcal{R}$$

Inversion:

$$\sigma_z f(\vec{k}) \sigma_x \sigma_z = f(-k) \sigma_z$$
$$\Rightarrow f(k) = -f(-k)$$

Need to
give up TRS
for $f \neq 0$

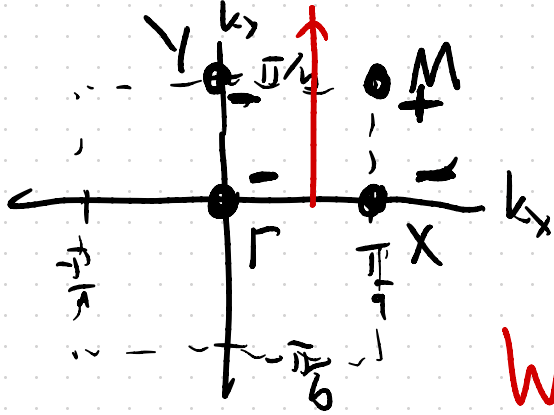
TRS

~~$$f^*(k) = f(k) = f(-k)$$~~

$$f(\vec{k}) = -t_2 \sin k_y b$$

$$h^{2D}(k_x, k_y) = \left[\Delta (1 + \cos k_y b) + t_1 \cos k_x a \right] \sigma_z + t_2 \sin k_x a \sigma_y - t_2 \sin k_y a \sigma_x$$

Gapped $\forall k$ when $\Delta > t_1 > 0$, $t_2 \neq 0$



Consider Wilson loops

I-eigenvalues of
- energy band

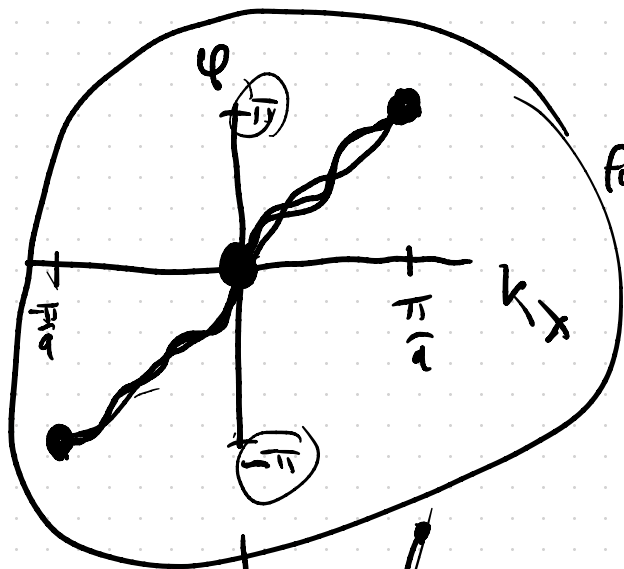
$$W_{\frac{2\pi}{b}c_0}(k_x) = \vec{u}_{-(k_x, 0)}^\dagger V\left(\frac{2\pi}{b}y\right) \prod_{k_y}^{\frac{2\pi}{b}c_0} P(k_x, k_y) \cdot \vec{u}_{-(k_x, 0)}$$

at $k_x=0$

$$\det W_{\frac{2\pi}{b}c_0}(k_x=0) = e^{i\pi 2} = 1 = e^{i\varphi_-(k_x=0)}$$

$$\varphi_-(k_x=0) = 0 \pmod{2\pi}$$

at $k_x = \frac{\pi}{a}$



for our model

$$\det W_{\frac{2\pi}{b} \leftarrow 0} \left(k_x = \frac{\pi}{a} \right) = -1$$

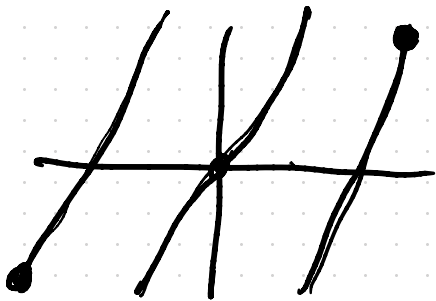
$$\varphi_{-} \left(k_x = \frac{\pi}{a} \right) = \pi \pmod{2\pi}$$

$\varphi_{-}(k_x)$ is related to $P_y P$ eigenvalues

$$P_y P |W_{k_x R_y}\rangle = \left(R_y + \frac{b}{2\pi} \varphi_{-}(k_x) \right) |W_{-k_x R_y}\rangle$$

Inversion: $\varphi_{-}(-k_x) = -\varphi_{-}(k_x)$

define an integer C - the # of times $\varphi(k_x)$



$$(-1)^C = (-1)^{\sum_{\text{TRANS}} n_k^-}$$

winds from $-\pi$ to π as
 k_x goes from $-\frac{\pi}{a} \rightarrow \frac{\pi}{a}$

C must be odd for our model

Formula for C

$$e^{i\varphi_-(k_x)} = \det W(k_x) = \det W(k_x + \frac{2\pi}{a}) = e^{i\varphi_-(k_x + \frac{2\pi}{a})}$$

$$\varphi_-(k_x + \frac{2\pi}{a}) = \varphi_-(k_x) + 2\pi C$$

$$\varphi_-(k_x) = C a k_x + f(k_x)$$

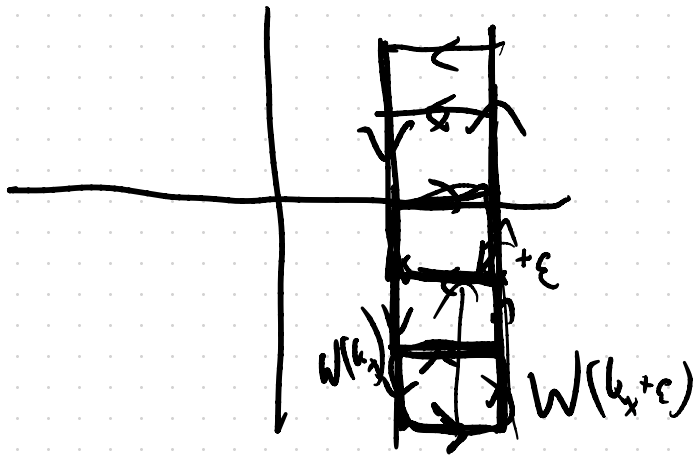
$$f(k_x + \frac{2\pi}{a}) = f(k_x)$$

$$\frac{1}{2\pi} \int_{-\frac{\pi}{a}}^{\frac{\pi}{a}} dk_x \frac{\partial \varphi}{\partial k_x} = \frac{1}{2\pi} \int_{-\frac{\pi}{a}}^{\frac{\pi}{a}} dk_x \left(C a + \cancel{\frac{\partial \varphi}{\partial k_x}} \right) \ominus$$

$$= C = \frac{1}{2\pi} \int_{-\frac{\pi}{a}}^{\frac{\pi}{a}} dk_x \frac{\partial}{\partial k_x} \text{Im} \log \det W(k_x)$$

$$\frac{\partial}{\partial k_x} \text{Im} \log \det W = \text{Im} \lim_{\varepsilon \rightarrow 0} \left[\frac{\log \det W(k_x + \varepsilon) - \log \det W(k_x)}{\varepsilon} \right]$$

$$= \text{Im} \lim_{\varepsilon \rightarrow 0} \frac{1}{\varepsilon} \log \det W(k_x + \varepsilon) W^\dagger(k_x)$$



Stokes thm

$$\begin{aligned}
 \frac{\partial}{\partial k_x} \text{Im} \log \det W &= \text{Im} \lim_{\epsilon \rightarrow 0} \frac{1}{\epsilon} \sum_{\text{Squares}} \log \det W_{\square} \\
 &= \lim_{\epsilon \rightarrow 0} \frac{1}{\epsilon} \sum_{\text{Squares}} \oint_{\text{Square}} dk \text{tr} A(\vec{k})
 \end{aligned}$$

$$= \int dk_y \operatorname{tr} \Omega_{xy}(k_x, k_y)$$

where $\Omega_{xy}(k_x, k_y) = \frac{\partial A_y}{\partial k_x} - \frac{\partial A_x}{\partial k_y} - i[A_x, A_y]$

$$C = \frac{1}{2\pi} \int dk_x \int dk_y \operatorname{tr} \Omega_{xy}(k_x, k_y)$$

$$= \frac{1}{2\pi} \oint d^2k \operatorname{tr} \Omega_{xy}(k_x, k_y)$$

\hat{C} Chern number $\in \mathbb{Z}$

Consider HWFs $|W_{-, k_x, k_y}\rangle$

$$P_y P |W_{-, k_x, R_y}\rangle = \left[R_y + \frac{b}{2\pi} \varphi(k_x) \right] |W_{-, k_x, R_y}\rangle$$

$$P_y P |W_{-, k_x + \frac{2\pi}{a}, R_y}\rangle = \left[R_y + \frac{b}{2\pi} \varphi(k_x + \frac{2\pi}{a}) \right] |W_{-, k_x + \frac{2\pi}{a}, R_y}\rangle$$

$$= \left[\underbrace{(R_y + Cb)}_{\text{smiley}} + \frac{b}{2\pi} \varphi(k_x) \right] |W_{-, k_x + \frac{2\pi}{a}, R_y}\rangle$$

$$\Rightarrow |W_{-, k_x + \frac{2\pi}{a}, R_y}\rangle = |W_{-, k_x, R_y + Cb}\rangle$$

Consequences

① Apply electric field $\vec{E} = E_0 \hat{x}$

$$\vec{A} = -E_0 t \hat{x} \quad \vec{E} = -\frac{\partial A}{\partial t}$$

$$\vec{p} \rightarrow \vec{p} - q\vec{A} = \vec{p} + qE_0 t \hat{x}$$

$$k \rightarrow k(t) = (k_x + qE_0 t, k_y)$$

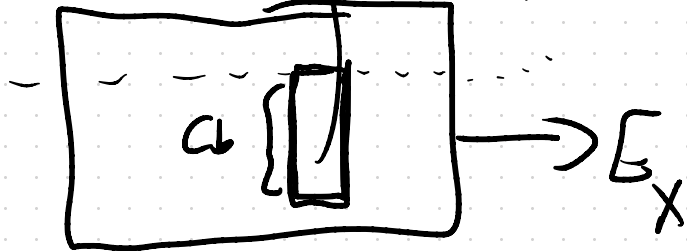
$$|W_{-k_x, R_y}\rangle \rightarrow |W_{-k_x + qE_0 t, R_y}\rangle$$

after one period

$$T = \frac{2\pi\hbar}{aqE_0}$$

$|W_{-k_x(t), R}\rangle$ has

moved C unit cells in the
 y direction



$$\Delta Q = C q N_x \quad V$$

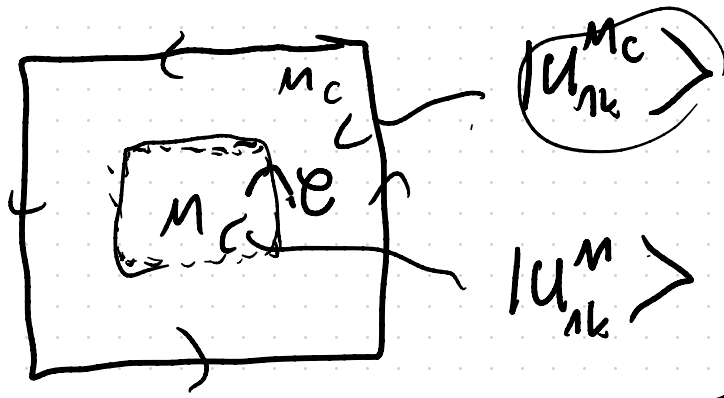
$$\Rightarrow I_y = \frac{\Delta Q}{T} = C q \left(\frac{q}{2\pi\hbar} \right) \overbrace{N_x a E_0}^V$$

$$= \underbrace{C \frac{q^2}{h}}_{G_H} V_x$$

G_H - Hall conductance

\rightarrow Integer $\times \frac{q^2}{h}$ for an insulator

Claim $C \neq 0$ the occupied bands do not have exponentially localized Wannier fns



$$C = \frac{1}{2\pi} \oint_{\mathcal{C}} d^2k \operatorname{tr} \Omega = \frac{1}{2\pi} \left[\iint_M d^2k \operatorname{tr} \Omega + \iint_{MC} d^2k \operatorname{tr} \Omega \right]$$

$$= \frac{1}{2\pi} \left[\oint_{\mathcal{C}} \operatorname{tr} A_m d^2k - \oint_{\mathcal{C}} \operatorname{tr} A_{mc} d^2k \right]$$

on \mathcal{C} $\operatorname{tr} A_{mc} = \operatorname{tr} A_m + \nabla \varphi$

$$C = \frac{1}{2\pi} \oint_C \nabla \varphi \cdot d\vec{k}$$

$\Rightarrow C \neq 0$ means I can't choose a single smooth gauge for my Bloch functions

\rightarrow No exponentially localized WFs.

$$|W_{n\vec{R}}\rangle = \frac{1}{(2\pi)^3} \int d^3k |\tilde{\psi}_k\rangle e^{-ik \cdot \vec{R}}$$

