

Gleyzes, S., Kuhr, S., Guerlin, C. et al. Quantum jumps of light recording the birth and death of a photon in a cavity. Nature 446, 297–300 (2007).

Requirements for a Quantum Measurement

- Quantum states, observables, operators, the environment
- Quantum measurement is statistical (repeatable measurements, large in number)
- Often need very well isolated systems (High Q; low energy loss)
- Decoherence rates must outlive measurement times



Repeatable Quantum Measurements

- What quantum behavior can we observe and how do we do that?
- Quantum non-demolition (QND) measurement of quantum state fluctuations
 - QND preserves observable uncertainty and evolution; repeatable



• Quantum jumps have been observed in trapped massive particles, but what about light quanta?

McIntyre, David H., et al. Quantum Mechanics. Pearson, 2016.

Measuring Quantum Jumps in Light

- Quantum jumps in photon number
- QND probes are Rb atoms in Rydberg states (dipole polarizability)
 - Detectable light shifts of single-photon resolution
- High Q cavity coupled to qubit (atom)
 - Enables measuring n(t)
 - Photon exchange (qubit and cavity)





A. Mortezapour, Quantum Inf Process 19, 136 (2020).

Experimental Set-up for Detecting Quantum Jumps

- B: box containing rubidium atoms, prepared in state |g>
- R_1, R_2 : Ramsey cavities
- S: classical microwave source
- *C*: photon box, a cavity that hosts the photons to be detected
- *D*: state selective field ionization detector
- *R*₁-*C*-*R*₂ interferometric arrangement cooled to 0.8K and shielded from thermal radiation



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1. Rubidium atoms in the circular Rydberg state $|g\rangle$ (n=50) are prepared and emitted from the box *B*. Travel along the blue axis at 250 ms⁻¹.



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- 1. Rubidium atoms in the circular Rydberg state $|g\rangle$ (n=50) are prepared and emitted from the box *B*. Travel along the blue axis at 250 ms⁻¹.
- 1. In R_1 , atoms undergo a change to a superposition state of $|e\rangle$ (n=51) and $|g\rangle$, $(|e\rangle + |g\rangle)/\sqrt{2}$



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3. Interaction of the atom with cavity *C*. Different photon state in cavity \rightarrow different interaction. Atomic state gains a phase shift $\Phi(n,\delta)$, become $(|e\rangle + \exp[i\Phi(n,\delta)]|g\rangle)/\sqrt{2}$



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- 3. Interaction of the atom with cavity *C*. Different photon state in cavity \rightarrow different interaction. Atomic state gains a phase shift $\Phi(n,\delta)$, become $(|e\rangle + \exp[i\Phi(n,\delta)]|g\rangle)/\sqrt{2}$
- 3. In R_2 , the Ramsey pulse brings the atom to state $|g\rangle$ if n=0, and to state $|e\rangle$ if n=1



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- 3. Interaction of the atom with cavity *C*. Different photon state in cavity \rightarrow different interaction. Atomic state gains a phase shift $\Phi(n,\delta)$, become $(|e\rangle + \exp[i\Phi(n,\delta)]|g\rangle)/\sqrt{2}$
- In R₂, the Ramsey pulse brings the atom to state |g⟩ if n=0, and to state |e⟩ if n=1

5. Detector counts the number of



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Experiment 1: How does it detect quantum jumps?

Detecting quantum jumps

- Measuring the states of atoms \rightarrow infer the photonic state
- If $|g\rangle \rightarrow |0\rangle$, if $|e\rangle \rightarrow |1\rangle$
- Obtain photon numbers as a function of time

Is it perfect?

- No! Conditional probability: P(g|1) = 13%, P(e|0) = 9%
- Majority vote involving 7 more atoms at any time
- Reduced probabilities: P(g|1) = 0.14%, P(e|0) = 0.025%

Experiment 2: Decay of the |1> state

- Just some additional steps compared to experiment 1
- Photonic state initialized to |0> by using ~10 atoms in |g>, tuned to resonance with cavity mode. No photon left in *C*.



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Experiment 2: Decay of the |1> state

- 2. 1 rubidium atom in state |e⟩ sent to C. Interaction time adjusted such that it exits C in |g⟩ and leaves C in |1⟩
- 2. Same as the steps in experiment 1



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Data and Results - Field Fluctuations

- 2.5 s experiment of 2241 counts
- Creation event at t = 1.054 s
 - o 0.476 s lifetime
- Background average occupation $n_0 = 0.063 \pm 0.005$ $n_t = 0.049 \pm 0.004$
- Characterize atom emission heating as 10⁻⁴ per atom per second



Data and Results - Single Photon State Decay



Critiques

- Multiple readings required to find critical details.
- Claiming to record the birth and death of photons is contrived when using a polling method.
- Setup measures n=1 and n!=1.
- Setup is not precise enough to distinguish expected stimulated decay contributions in experiment 2.
 - $T_{1,measured} = 0.097 \pm 0.005s; T_{1,expected} = 0.102 \pm 0.004s$
 - What do the error bars mean?
- Time between atomic detections is random.
 - \circ P_{measurement} = 0.063

"In this experiment, the detection does not distinguish between [state 2] and [state 0]."



Author's Conclusions

- Successful quantum non-demolition measurements on photons
- Observation of cavity photon's "birth" and "death"
- Demonstration of well-established ensemble behavior

Citation Evaluation

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Quantum jumps of light recording the birth and death of a photon in a cavity (Article) (Open Access)

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- "A spectacular recent achievement is the ability to perform quantum 'nondemolition' experiments, in which photons in a microwave cavity can be monitored without destroying them, revealing the progressive collapse of the wavefunction under successive measurements." (R. J. Schoelkopf & S. M. Girvin, 2008)
- "As their most spectacular sensing application, Rydberg atoms in vacuum have been employed as single-photon detectors for microwave photons in a cryogenic cavity in a series of experiments that was highlighted by the Nobel prize in physics in 2012." (Degen *et al.*, 2017)