Hund's rules state that, to determine the (a) (S, L, J) ground state of an atom, ne (DFirst, maximize S. (2) Then, maximize L. (3) Finally, for more than half-filled shells, maximize J; otherwise minimize J. When applying Hund's rules, we need to take care to only consider valid states. Applying Hundis rules to Li, B, N, Na, and K, we get Element S term symbol J 2 S 1/2 L: (2=3) 1/2 1/2 Ö B (2=5) 1/2 2P1/2 1/2 N (Z=7) 453/2 3/2 3/2 · 0 Na (7=11) 1/2 2542 1/2 D K (Z=19) 1/2 2S1/2 1/2 ·D Nitrogen is a little tricky. It has three valence electrons in the 2p shell. with $l_1 = l_2 = l_3 = 1$.

(acon) To maximize S to 3/2, we need that $m_{s_1} = m_{s_2} = m_{s_3} = \pm 1/2$ in other words that all three electrons' spins are aligned. So, in the 1 l, meinsi ; l2 meznsz ; l3 mizms 3 > basis, we know that the ground state most satisfy 11 me 1/2; 1 mez 1/2; 1 mez 1/2). Since these are electrons, they are termions and cannot occupy the same state. Therefore, me, mez, mez must all be different. Since mei 62-1,0,13 for li=1, the only possible state satisfying S= 3/2 and li=lz=lz=1 is the $m_{l_{1}} = -1$, $m_{l_{2}} = 0$, $m_{l_{3}} = +1$ state: 11-1 42; 10 1/2; 1 1/2>. This state has ML=D, which implies that L=O. This means that the only possible J is J=312.

(b) The lowest frequency line in the absorption spectrum of Na corresponds to an electron in the 3s level absorbing a photon and being excited to a 3p state. There are two such states, split by the spin-orbit effect. The splitting is given by Eso ≈ - ms·B where $\bar{\mu_s} = -g_{s}\mu_{i3}\,\bar{S}/\hbar$ is the magnetic moment of the electron and B is an effective magnetic field. Semiclassically, we can think of the B Aeld as avising from a charged nucleus rotating around the electron: B & Q X E. (related to Birt-Savart law) The E field is that of a point charge $\tilde{E} \approx \frac{e}{r^2} \hat{r}$. $(\vec{r} = r\hat{r})$ $\vec{B} \propto \frac{qe}{cr^2} \vec{y} \times \hat{r} = \frac{qe}{cmr^2} (m \vec{y} \times \hat{r})$ (Ľ=mvxř) $= \frac{qe}{cmr^3} \left(mv xr^2 \right) = \frac{qe}{cmr^3} L$

This then tells us that the (bcon) spin-orbit splitting can also be written as Eso = -Ms. B $= - \begin{bmatrix} g_{SMB} & \vec{S} \end{bmatrix} \cdot \begin{bmatrix} g_e & \vec{L} \end{bmatrix}$ - <u>gs Mage</u> <u>1</u> L.S mct r3 ~ L·S·+3 For the particular states considered, Land S are fixed and non-zero, so the splitting is proportional to (+3). n=-3. (c) For potassium K(Z=19), there is one 4s valence election in the ground state. The term symbol corresponding to this (L=0, S= 1/2, J=1/2) state is S1/2. The lowest single valence election states are OnexOnar Omin Ethend's rales 25.12 (S, L, J) = (12, 0, 1/2)[Av] 4s 2P1/2 (1/2, 1, 1/2) [Ar] 4p 2P312 (1/2, 1, 3/2)

The energy levels look like (Con) (4p) 2P312 $- (M_{J} = -\frac{3}{2}, -\frac{1}{2}, \frac{1}{2}, \frac{3}{2})$ $- (M_{J} = -\frac{1}{2}, \frac{1}{2}, \frac{1}{2})$ 2P1/2 (4s) 251/2 - (Mg=-1/2, 1/2) The energy of the electron's magnetic dipole in a magnetic field $\vec{B} = B\hat{z}$ is $\hat{H}' = -\hat{\mu}_{J}\cdot\hat{B} = \stackrel{e}{=} \stackrel{e}{=} \stackrel{f}{=} \stackrel{f}{=} \stackrel{f}{=} \stackrel{g}{=} \stackrel{f}{=} \stackrel{g}{=} \stackrel{f}{=} \stackrel{g}{=} \stackrel{f}{=} \stackrel{g}{=} \stackrel{f}{=} \stackrel{g}{=} \stackrel{f}{=} \stackrel{g}{=} \stackrel{g$ where $\vec{\mu}_J = -\frac{e}{2m}g^{\vec{J}}$ is the magnetiz dipole of the electron in an (5,2,5) atomic state (g is the landé-g factor, which depends on S, L, J). So, larger g and MJ correspond to larger changes in energies due to magnetic fields. The perturbed energy levels look like: $\begin{array}{c} (4_{p}) & ^{2}P_{312} & (g=4/3) \\ & ^{2}P_{1/2} & (g=2/3) \\ (4_{s}) & ^{2}S_{1/2} & (g=2) \end{array}$ My= 31/2 1 +1/2 (exaggerated slopes) Energy Din Bomagnetic field

(d) All of the allowed EI (dipole) transitions to the 25112 states are 2P312 2P1/2 -----25 1/2 which were determined by following the selection rales.