## Phys 486 Discussion 14 - Angular Momentum Addition : Clebsch-Gordan Coefficients

We are learning how to add two angular momenta, $j_{1}$ and $j_{2}$. We can add them in two ways: we can

- list all the eigenstates $\left|m_{1} m_{2}\right\rangle$ of the operators $j_{1 z} \& j_{2 z}$ of the INDIVIDUAL ang. momenta, OR
$\bullet$ list all the eigenstates $|J M\rangle$ of the TOTAL ang. mom. operators $J^{2} \equiv\left|\vec{j}_{1}+\vec{j}_{2}\right|^{2} \& J_{z} \equiv\left(j_{1 z}+j_{2 z}\right)$.
The eigenstates $\left|m_{1} m_{2}\right\rangle$ and $|J M\rangle$ provide different bases with which to describe the sum of one angular momentum $j_{1}$ with another one $j_{2}$. In class, we learned how to read tables of Clebsch-Gordan coefficients to express a $|J M\rangle$ as a linear combination of $\left|m_{1} m_{2}\right\rangle$ 's. These tables are now on our website, filename Formulae-CGtables.pdf. In this discussion, we will learn how to calculate the CG coefficients so that we understand exactly what they are, and where they come from, and that the scary-looking CG tables are not in the least bit mysterious.


## Problem 1 : Deuterium Atom

We'll continue to explore our example from class, which was to add these two angular momenta:

- $j_{1}=1$ is the $\operatorname{spin} s_{d}$ of a deuterium nucleus (a.k.a. a "deuteron")
- $j_{2}=1 / 2$ is the spin $s_{e}$ of an electron.

If we bring the deuteron and electron together, we get a deuterium atom with total spin $j_{1} \oplus j_{2}$. The " $\oplus$ " symbol means "an addition that's more complex than $2+3=5$ ". As we discussed, " $j_{1} \otimes j_{2}$ " is also used.
(a) Make a plot with the $m_{1}$ axis pointing upward and the $m_{2}$ axis pointing sideways, and mark with solid circles all the points $\left(m_{1}, m_{2}\right)$ where $\left|m_{1} m_{2}\right\rangle$ is a physically-possible state. How many states do you have?
(b) Make a plot with the $J$ axis pointing upward and the $M$ axis pointing sideways, and mark with solid circles all the points $(J, M)$ where $|J M\rangle$ is a physically possible state. Leave lots of space between your tick marks on the horizontal ( $M$ ) axis ... like an inch of space $\ldots$. we're going to write things under these tick marks. To make your plot, you need the first of the two angular momentum addition rules we learned today:
(1) The total $J$ quantum number runs from $\left|j_{1}-j_{2}\right|$ to $\left|j_{1}+j_{2}\right|$ in steps of 1 .
(2) The total $M$ quantum number is additive: $M=m_{1}+m_{2}$.

You also need the rule restricting the $m$ quantum number for any angular momentum:

- The $m$ (or $M$, or $m_{2}, \ldots$ ) quantum number runs from $-j$ to $+j$ (or $-J$ to $+J$, or $-j_{2}$ to $+j_{2}, \ldots$ ) in steps of 1 . How many states do you have? It should be the same as in part (a)!
(c) Underneath each of the $M$ tick marks that has at least one solid circle above it - i.e. under each $M$ value that has at least one valid $|J M\rangle$ eigenstate - write a list of all the $\left|m_{1} m_{2}\right\rangle$ states that might contribute to said $M$ value. HINT: You need the second of our two angular momentum addition rules.
(d) As we discussed, $m_{1}+m_{2}=M$, but $j_{1}+j_{2} \neq J$. The $m$ quantum number is additive, but the $j$ quantum number is not. Do you understand why this is? If not, check the hint ${ }^{1}$... if it's still unclear, ask!!!

[^0](e) Consult the Clebsch-Gordan tables on the web and write the eigenstate $|J, M\rangle=\left|\frac{1}{2},-\frac{1}{2}\right\rangle_{J, M}$ as a linear combination of $\left|m_{1}, m_{2}\right\rangle$ eigenstates.
(f) Consult nothing and write the $\left|\frac{3}{2},-\frac{3}{2}\right\rangle_{J, M}$ eigenstate as a linear combination of $\left|m_{1}, m_{2}\right\rangle$ eigenstates.
(g) You just realized that the $|J, M\rangle$ states $\left|\frac{3}{2}, \pm \frac{3}{2}\right\rangle_{J, M}$ each match a SINGLE $\left|m_{1}, m_{2}\right\rangle$ state: $\left| \pm 1, \pm \frac{1}{2}\right\rangle_{m_{1}, m_{2}}$.

These are called stretched states because all $j$ vectors are maximally aligned along or against the $z$ axis. The stretched states should be found at the top-right and top-left corners of the $(M, J)$ plot you've been making. Circle the stretched state with positive $M$. This will be the starting point for our Clebsch-Gordan calculations because it is always uniquely determined : $\left|J_{\max }, M_{\max }\right\rangle=\left|m_{1 \max }, m_{2 \max }\right\rangle=\left|+j_{1},+j_{2}\right\rangle$. Simple starting point!
(h) THE WHOLE THING: Here is the procedure for building all the $|J, M\rangle$ states from $\left|m_{1}, m_{2}\right\rangle$ states, and thereby calculating all the Clebsch-Gordan (CG) coefficients for the given $j_{1}$ and $j_{2}$ values:
A. Build the stretched state $\left|J_{\text {max }}, M_{\text {max }}\right\rangle=$ top-right point on your plot.
B. Build the states with the same $J$ but lower $M$ - step to the LEFT on your plot - by applying the step-down operator $J_{-}=j_{1-}+j_{2_{-}}$. This step has two parts:

1. use $\hat{J}_{ \pm}|j, m\rangle=\hbar \sqrt{j(j+1)-m(m \pm 1)}|j, m \pm 1\rangle \rightarrow$ apply $J_{-}$to the $|J, M\rangle$ form
2. use the concept $\hat{Q}_{1+2}=\hat{Q}_{\text {lonly }}+\hat{Q}_{\text {2only }} \quad \rightarrow$ apply $J_{-}=j_{1-}+j_{2-}$ to the $\left|m_{1}, m_{2}\right\rangle$ form

Step-down until you have filled in the whole line with the $J$ you are currently working on. Check the CG tables on the course website $\rightarrow$ are your answers correct?
C. Go back to the right-hand-side of the $J$ line you've just completed, move one step to the left, to $|J, J-1\rangle_{J, M}$, then move DOWN one step and build $|J-1, J-1\rangle_{J, M}$ using orthonormality.

How does that work exactly? $\rightarrow$ Look at the list of $\left|m_{1} m_{2}\right\rangle$ states you wrote below this $M$ column, and you'll see that there are exactly the same number of them as there are $|J M\rangle$ states in the column. It must be so. This whole exercise is a change of basis: we are building an orthonormal basis $|J M\rangle$ from a different orthonormal basis $\left|m_{1} m_{2}\right\rangle$. The states with a particular $M$ value - call it $M_{0}$ - are a subset of both bases, satisfying $M=M_{0}$ on the one hand and $m_{1}+m_{2}=M_{0}$ on the other hand. The number of $\left|m_{1} m_{2}\right\rangle$ states in this subspace must be the same as the number of $|J M\rangle$ states since both sets span the same subspace. Now realize that you have already built all but one of the $|J M\rangle$ states in this column. Since you only have one left, and you know all the $\left|m_{1} m_{2}\right\rangle$ states you can use to build it, you can build the final $|J M\rangle$ state in the column by making it orthonormal to the other $|J M\rangle$ states in the column. Get the idea? You will when you try it!

Repeat steps B and C until all the states $|J M\rangle$ are written as linear combinations of the states $\left|m_{1} m_{2}\right\rangle$. The figure at right may help you to visualize the procedure:
A. START at stretched state with maximum $M$
B. move LEFT via step-down operator $J_{-}=j_{1-}+j_{2-}$
C. move DOWN at bottom of a column via orthonormality


Add a spin-3/2 particle and a spin-1 particle to build the state $|J, M\rangle=\left|\frac{3}{2},+\frac{1}{2}\right\rangle_{J, M}$ from $\left|m_{1}, m_{2}\right\rangle$ states.
Calculate the result using the $A, B, C$ procedure above, doing only the steps you need (!!!) to reach the state you want. Check your answer against the CG table on the website, of course.


[^0]:    ${ }^{1} j$ and $m$ are quantum numbers. They label the eigenvalues of certain operators, but they are not eigenvalues. (For a more familiar example, recall the energy quantum number $n$ : it's an integer that labels energy eigenvalues but it is not an energy itself; we derived formulae like $E_{n}=-13.6 \mathrm{eV} / n^{2}$ to relate $n$ to the corresponding eigenvalue $E_{n}$. The relations between $j \& m$ and the e-values they label are: $\hat{j}_{z}|j m\rangle=\hbar m|j m\rangle$ and $\hat{j}|j m\rangle=\hbar \sqrt{j(j+1)}|j m\rangle$. The $m$ eigenvalue is linear in $m$, therefore $m$ 's are additive.

