

Q1

$$(a) \langle \Psi_2 | \Psi_2 \rangle = \left(\frac{4}{5}\right)^2 \langle \phi_1 | \phi_1 \rangle - \frac{12}{25} \langle \phi_1 | \phi_2 \rangle - \frac{12}{25} \langle \phi_2 | \phi_1 \rangle + \left(\frac{3}{5}\right)^2 \langle \phi_2 | \phi_2 \rangle = \left(\frac{4}{5}\right)^2 + \left(\frac{3}{5}\right)^2 = 1$$

One b/c ϕ_1, ϕ_2 are normalized
 zero b/c eigenstates are orthogonal

$$\langle \Psi_1 | \Psi_2 \rangle = \left(\frac{3}{5} \langle \phi_1 | + \frac{4}{5} \langle \phi_2 | \right) \left(\frac{4}{5} |\phi_1\rangle - \frac{3}{5} |\phi_2\rangle\right) = \frac{12}{25} \langle \phi_1 | \phi_1 \rangle + \frac{16}{25} \langle \phi_2 | \phi_1 \rangle - \frac{9}{25} \langle \phi_1 | \phi_2 \rangle - \frac{12}{25} \langle \phi_2 | \phi_2 \rangle$$

$$= \frac{12}{25} - \frac{12}{25} = 0$$

(b) From the 1st equation, we have $|\phi_1\rangle = \frac{5}{3} |\Psi_1\rangle - \frac{4}{3} |\phi_2\rangle$. Plugging this in to the 2nd equation gives $|\Psi_2\rangle = \frac{4}{5} \left(\frac{5}{3} |\Psi_1\rangle - \frac{4}{3} |\phi_2\rangle\right) - \frac{3}{5} |\phi_2\rangle$,
 or $|\phi_2\rangle = \frac{4|\Psi_1\rangle - 3|\Psi_2\rangle}{5}$. Plugging this in to our earlier eq,

$$|\phi_1\rangle = \frac{5}{3} |\Psi_1\rangle - \frac{4}{3} \left(\frac{4}{5} |\Psi_1\rangle - \frac{3}{5} |\Psi_2\rangle\right) = \frac{3|\Psi_1\rangle + 4|\Psi_2\rangle}{5}$$

(c) After measuring a_1 , we're in the eigenstate of \hat{A} w/ eigenvalue a_1 , so our new state is $|\Psi_1\rangle$

(d) We measure b_1 w/ probability $|\langle \phi_1 | \Psi_1 \rangle|^2 = \left|\frac{3}{5}\right|^2 = \frac{9}{25}$

We measure b_2 w/ probability $|\langle \phi_2 | \Psi_1 \rangle|^2 = \frac{16}{25}$

Eigenstate of our measurement operator \hat{B} .
 our current state

(e) If we measured b_1 : $P(a_1) = |\langle \Psi_1 | \phi_1 \rangle|^2 = \frac{9}{25}$
 If we measured b_2 : $P(a_1) = |\langle \Psi_1 | \phi_2 \rangle|^2 = \frac{16}{25}$
 So the total probability of getting a_1 is

$$P(b_1)P(a_1|b_1) + P(b_2)P(a_1|b_2)$$

$$= \frac{9}{25} \times \frac{9}{25} + \frac{16}{25} \times \frac{16}{25}$$

$$= \frac{337}{625} \approx 0.5392$$

So it is not 1!

(f) In this case, $P(a_1) = \frac{16}{25}$, as shown above.

(g) First, calculate $\hat{A}\hat{B}\psi_i$: $\hat{A}\hat{B}|\psi_i\rangle = \hat{A}\hat{B}\left[\frac{3}{5}|\phi_1\rangle + \frac{4}{5}|\phi_2\rangle\right]$

$$= \hat{A}\left(\frac{3b_1}{5}|\phi_1\rangle + \frac{4b_2}{5}|\phi_2\rangle\right)$$

$$= \hat{A}\left(\frac{3b_1}{5}\left(\frac{3}{5}|\psi_1\rangle + \frac{4}{5}|\psi_2\rangle\right) + \frac{4b_2}{5}\left(\frac{4}{5}|\psi_1\rangle - \frac{3}{5}|\psi_2\rangle\right)\right)$$

$$= \hat{A}\left(\left(\frac{9}{25}b_1 + \frac{16}{25}b_2\right)|\psi_1\rangle + \left(\frac{12}{25}b_1 - \frac{12}{25}b_2\right)|\psi_2\rangle\right)$$

$$= \left(\frac{9}{25}b_1 + \frac{16}{25}b_2\right)a_1|\psi_1\rangle + \left(\frac{12}{25}b_1 - \frac{12}{25}b_2\right)a_2|\psi_2\rangle$$

Now: $\hat{B}\hat{A}|\psi_i\rangle = \hat{B}a_i|\psi_i\rangle$

$$= \hat{B}\left(\frac{3}{5}a_1|\phi_1\rangle + \frac{4}{5}a_1|\phi_2\rangle\right)$$

$$= \frac{3}{5}a_1b_1|\phi_1\rangle + \frac{4}{5}a_1b_2|\phi_2\rangle$$

$$= \frac{3}{5}a_1b_1\left(\frac{3}{5}|\psi_1\rangle + \frac{4}{5}|\psi_2\rangle\right) + \frac{4}{5}a_1b_2\left(\frac{4}{5}|\psi_1\rangle - \frac{3}{5}|\psi_2\rangle\right)$$

$$= \left(\frac{9}{25}a_1b_1 + \frac{16}{25}a_1b_2\right)|\psi_1\rangle + \left(\frac{12}{25}a_1b_1 - \frac{12}{25}a_1b_2\right)|\psi_2\rangle$$

So in total,

$$[\hat{A}, \hat{B}]|\psi_i\rangle = \left(\frac{9}{25}a_1b_1 + \frac{16}{25}a_1b_2 - \frac{9}{25}a_1b_1 - \frac{16}{25}a_1b_2\right)|\psi_1\rangle$$

$$+ \left(\frac{12}{25}a_2b_1 - \frac{12}{25}a_2b_2 - \frac{12}{25}a_1b_1 + \frac{12}{25}a_1b_2\right)|\psi_2\rangle$$

$$= \frac{12}{25}(a_2b_1 - a_2b_2 - a_1b_1 + a_1b_2)|\psi_2\rangle$$

$$= \frac{12}{25}(a_1 - a_2)(b_1 + b_2)|\psi_2\rangle$$

We thus see $[\hat{A}, \hat{B}] \neq 0$ unless $a_1 = a_2$ or $b_1 = b_2$.

a) If this Hamiltonian is written in the basis $\{|\psi_1\rangle, |\psi_2\rangle, |\psi_3\rangle\}$, then in coords

Q2 we have $|\psi_1\rangle = \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix}$, $|\psi_2\rangle = \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix}$, $|\psi_3\rangle = \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$. Meanwhile,

$$H|\psi_1\rangle = \begin{pmatrix} a & 0 & b \\ 0 & c & 0 \\ b & 0 & a \end{pmatrix} \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix} = \begin{pmatrix} a \\ 0 \\ b \end{pmatrix} \neq K \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix}, \quad \text{thus } |\psi_1\rangle \text{ is not an eig of } H.$$

Similarly, you can show $|\psi_3\rangle$ is not an eig of H .

b) Note that $H|\psi_2\rangle = c|\psi_2\rangle$, so $|\psi_2\rangle$ is an eigenstate of H .

Thus, its time-dependence is given by multiplying by $e^{-iEt/\hbar}$. Here, $E=c$,

so

$$|\psi_2(t)\rangle = \begin{pmatrix} 0 \\ e^{-ict/\hbar} \\ 0 \end{pmatrix}$$

c) Note that $H|\psi_2\rangle = \begin{pmatrix} b \\ 0 \\ a \end{pmatrix} \neq K \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix}$, so $|\psi_2\rangle$ is not an eigenstate of

H . We thus need to expand $|\psi_2\rangle$ in eigenstates of H ; each eigenstate we will then be able to find the time-dependence of.

First: Find eigenstates of H .

$$\det(H - E\mathbb{1}) = 0, \text{ or } \begin{vmatrix} a-E & 0 & b \\ 0 & c-E & 0 \\ b & 0 & a-E \end{vmatrix} = 0, \text{ or}$$

$$0 = (c-E) \left[(a-E)^2 - b^2 \right] = (c-E) (a^2 - 2aE + E^2 - b^2) \\ = (c-E) (E-a+b)(E-a-b)$$

Thus, $E=c$ or $E=a-b$ or $E=a+b$

IF $E=c$, we have

$$\begin{pmatrix} a & 0 & b \\ 0 & c & 0 \\ b & 0 & a \end{pmatrix} \begin{pmatrix} e_1 \\ e_2 \\ e_3 \end{pmatrix} = c \begin{pmatrix} e_1 \\ e_2 \\ e_3 \end{pmatrix}, \text{ or } \begin{pmatrix} a e_1 + b e_3 \\ c e_2 \\ b e_1 + a e_3 \end{pmatrix} = \begin{pmatrix} c e_1 \\ c e_2 \\ c e_3 \end{pmatrix}$$

The ~~prop~~ equations are then

$$\begin{cases} a e_1 + b e_3 = c e_1 \\ b e_1 + a e_3 = c e_3 \end{cases} \quad c e_2 = c e_2 \quad \uparrow \text{always true}$$

→ solve for e_1 , get

$$e_1 = \frac{b}{c-a} e_3, \text{ plug in, get}$$

$$\left(\frac{b^2}{c-a} + a \right) e_3 = c e_3, \text{ or } 0 = \left(\frac{b^2}{c-a} + a - c \right) e_3, \text{ or}$$

① $0 = \frac{b^2 - (a-c)^2}{a-c} e_3$. Since $b^2 - (a-c)^2 \neq 0$, we have $e_3 = 0$.

Plugging $e_3 = 0$ to $e_1 = \frac{1}{c-a} e_3$, we find $e_1 = 0$. To normalize, we pick $e_2 = 1$, so our eigenstate is $\begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix}$.

If $E = a+b$

$$\begin{pmatrix} a & 0 & b \\ 0 & c & 0 \\ b & 0 & a \end{pmatrix} \begin{pmatrix} e_1 \\ e_2 \\ e_3 \end{pmatrix} = \begin{pmatrix} (a+b)e_1 \\ (a+b)e_2 \\ (a+b)e_3 \end{pmatrix}, \text{ or } \begin{pmatrix} ae_1 + be_3 \\ ce_2 \\ be_1 + ae_3 \end{pmatrix} = \begin{pmatrix} (a+b)e_1 \\ (a+b)e_2 \\ (a+b)e_3 \end{pmatrix}$$

Middle eqn gives $(a+b-c)e_2 = 0$, or $e_2 = 0$.

Top eqn gives $ae_1 + be_3 = ae_1 + be_1$, or $e_3 = e_1$.

For normalization, pick $e_1 = e_3 = \frac{1}{\sqrt{2}}$, so our eigenstate is $\begin{pmatrix} 1/\sqrt{2} \\ 0 \\ 1/\sqrt{2} \end{pmatrix}$.

If $E = a-b$

$$\begin{pmatrix} ae_1 + be_3 \\ ce_2 \\ ae_3 + be_1 \end{pmatrix} = \begin{pmatrix} ae_1 - be_1 \\ ae_2 - be_2 \\ ae_3 - be_3 \end{pmatrix}$$

Middle eqn gives $(a-c-b)e_2 = 0$, or $e_2 = 0$.

Top eqn gives $ae_1 + be_3 = ae_1 - be_1$, or

$$e_3 = -e_1.$$

For normalization, pick $e_1 = \frac{1}{\sqrt{2}}$, so our eigenstate is $\begin{pmatrix} 1/\sqrt{2} \\ 0 \\ -1/\sqrt{2} \end{pmatrix}$.

Second: Expand $\begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$ in eigenstates. We want to write

$$\begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} = A \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix} + B \begin{pmatrix} 1/\sqrt{2} \\ 0 \\ 1/\sqrt{2} \end{pmatrix} + C \begin{pmatrix} 1/\sqrt{2} \\ 0 \\ -1/\sqrt{2} \end{pmatrix}. \text{ We note that}$$

$$A = (0 \ 1 \ 0) \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} = 0, \quad B = (1/\sqrt{2} \ 0 \ 1/\sqrt{2}) \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} = \frac{1}{\sqrt{2}}, \quad C = (1/\sqrt{2} \ 0 \ -1/\sqrt{2}) \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} = -\frac{1}{\sqrt{2}}.$$

Thus,
$$\begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1/\sqrt{2} \\ 0 \\ 1/\sqrt{2} \end{pmatrix} - \frac{1}{\sqrt{2}} \begin{pmatrix} 1/\sqrt{2} \\ 0 \\ -1/\sqrt{2} \end{pmatrix}$$

We know the first eigenstate has time-dependence $e^{-i(a+b)t/\hbar}$, and the second has $e^{-i(a-b)t/\hbar}$, so

$$|\mathcal{B}(t)\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1/\sqrt{2} \\ 0 \\ 1/\sqrt{2} \end{pmatrix} e^{-i(a+b)t/\hbar} - \frac{1}{\sqrt{2}} \begin{pmatrix} 1/\sqrt{2} \\ 0 \\ -1/\sqrt{2} \end{pmatrix} e^{-i(a-b)t/\hbar} = \begin{pmatrix} \frac{1}{2} e^{-i(a+b)t/\hbar} - \frac{1}{2} e^{-i(a-b)t/\hbar} \\ 0 \\ \frac{1}{2} e^{-i(a+b)t/\hbar} + \frac{1}{2} e^{-i(a-b)t/\hbar} \end{pmatrix}$$