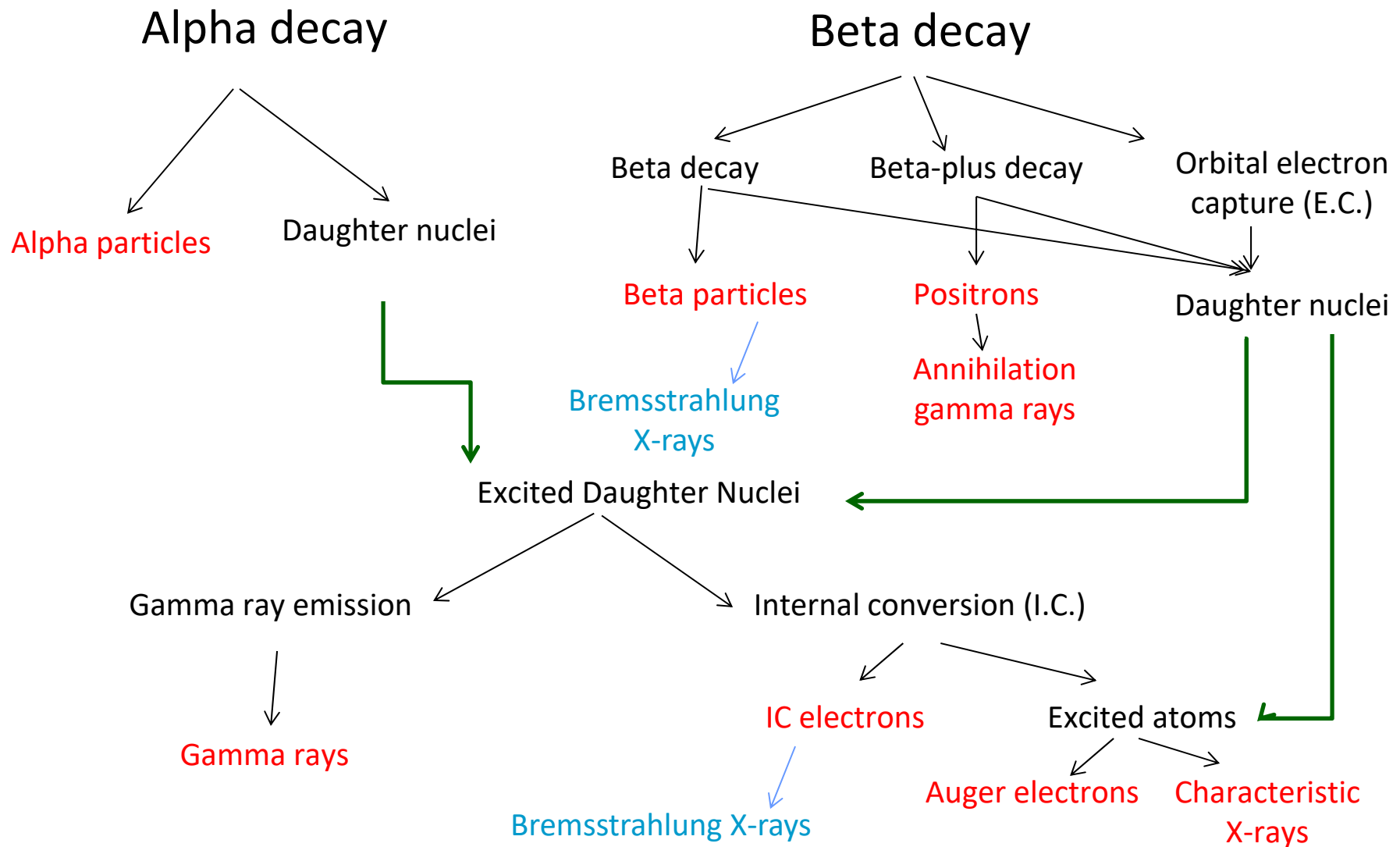


# Typical Decay Products from Unstable Radioisotopes





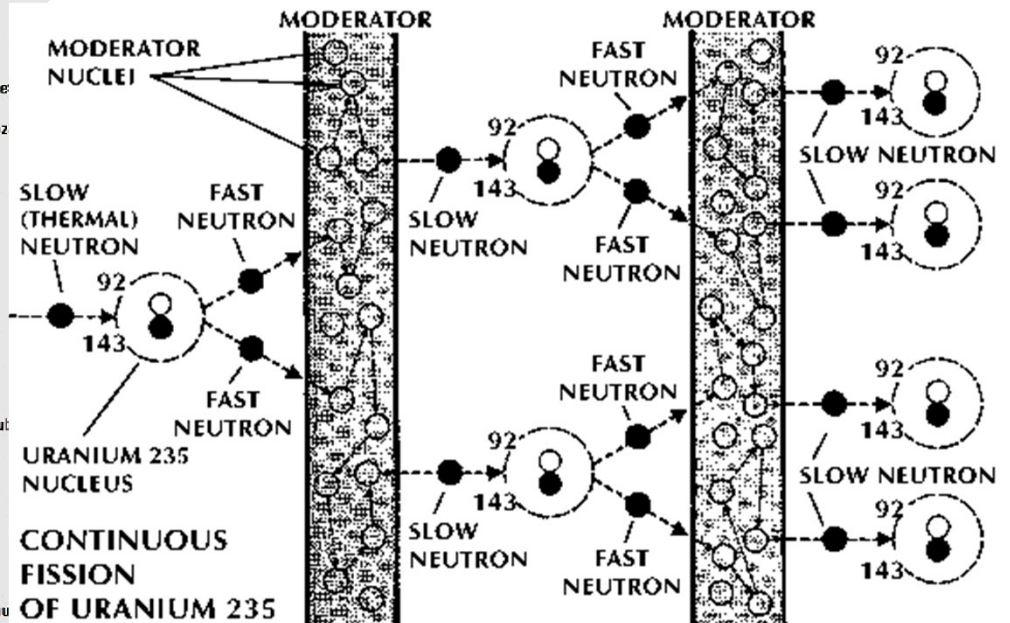
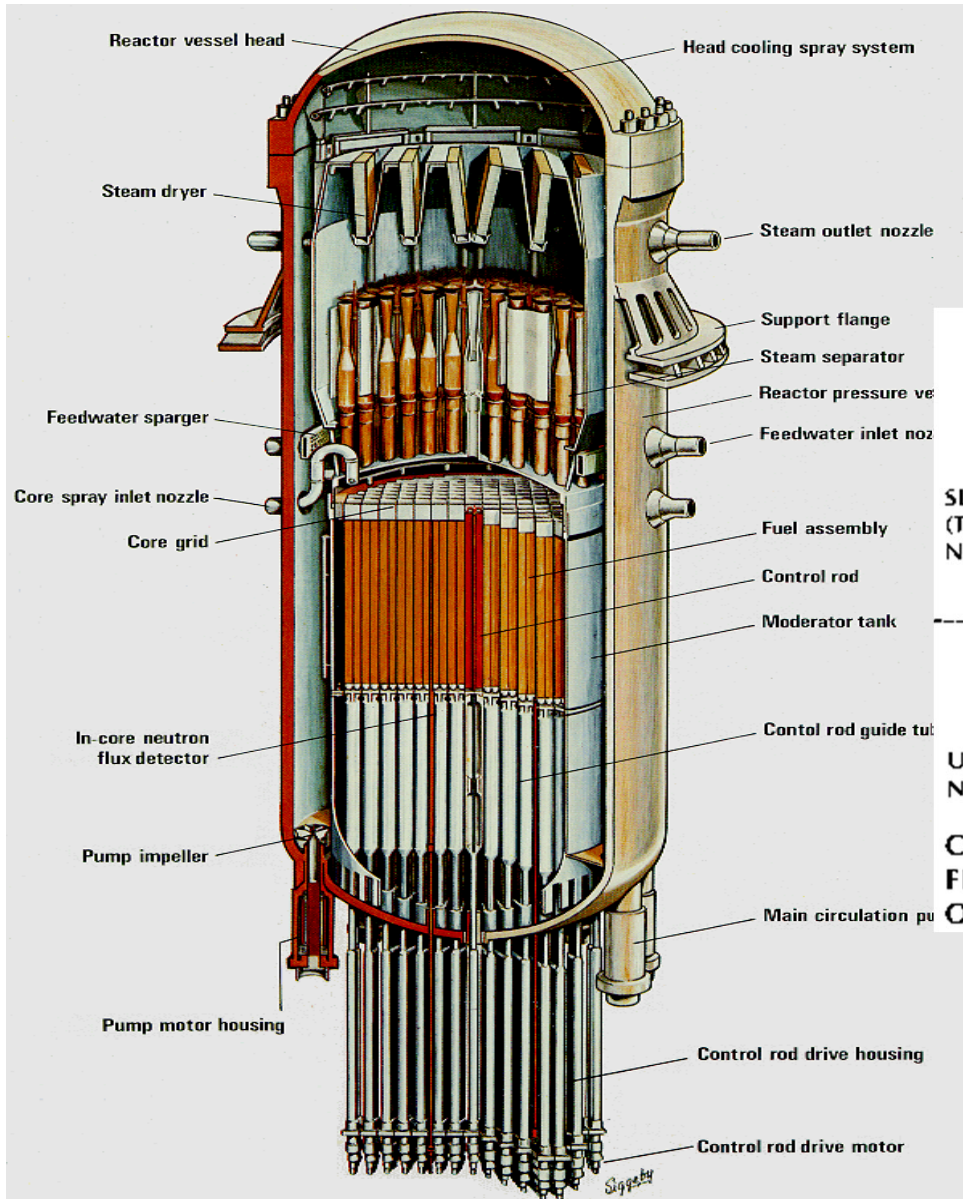
# Neutron Sources



## Typical Scenarios Involving Neutron Radiation

- Portable neutron sources
- Nuclear reactors
- Medical applications

# Inside a Nuclear Reactor



## Neutron Sources – Spontaneous Fission



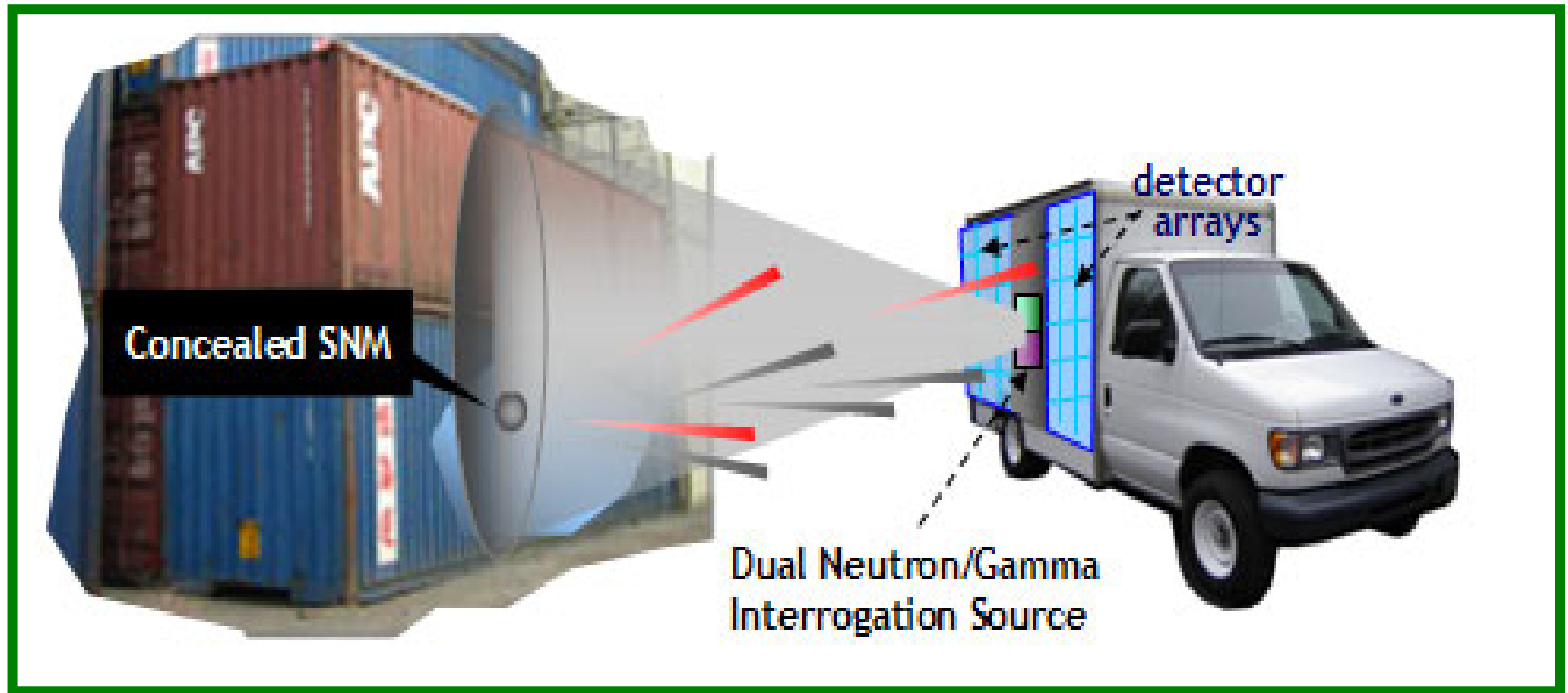
Cf-252 neutron source can be made extremely compact



An engineer tests the prototype Timed Neutron Detector, a device that detects landmines. The neutron source of the landmine detector holds a tiny amount of californium-252. (Photo credit: Pacific Northwest National Lab)

# Active Interrogation for Finding Special Nuclear Materials

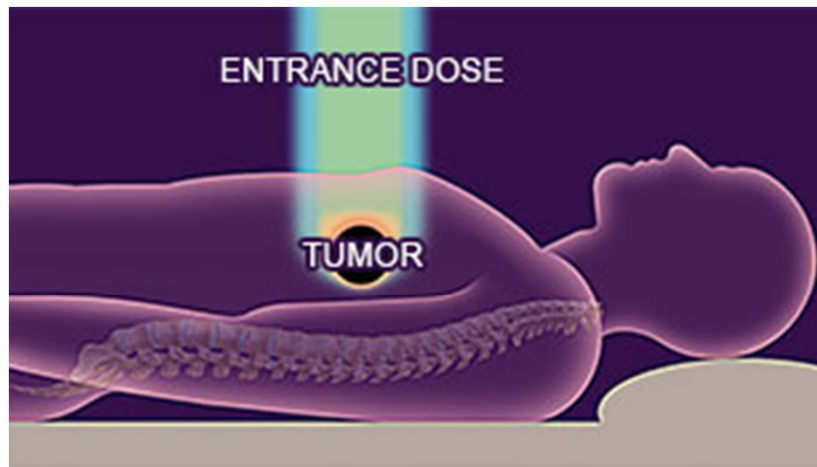
## Mobile Dual Neutron/Gamma Interrogation System



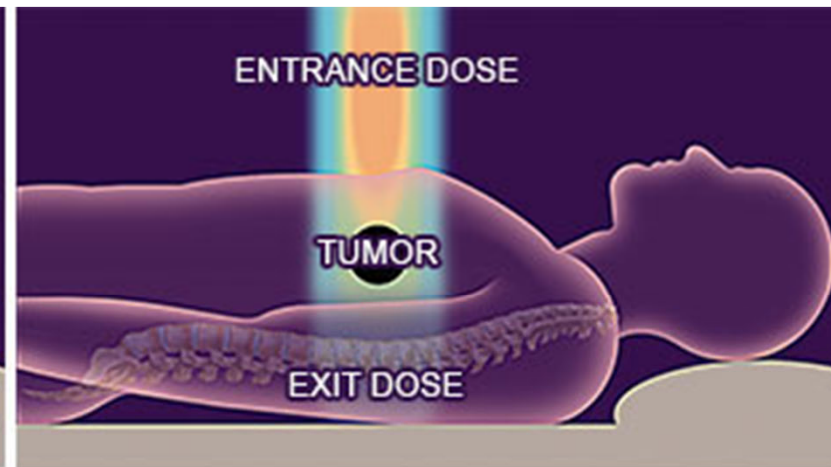


# Proton Therapy

<http://www.floridaproton.org/what-is-proton-therapy/benefits>



**TARGETED PROTON THERAPY:**  
Deposits most energy on target

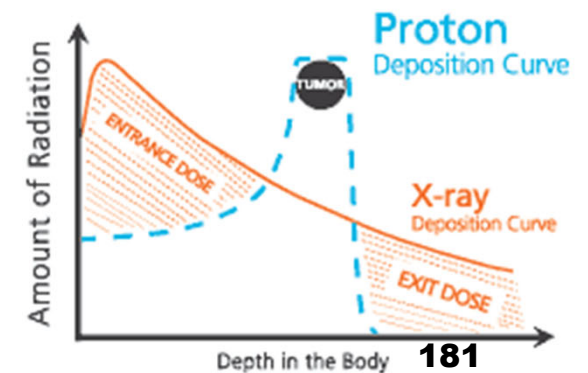
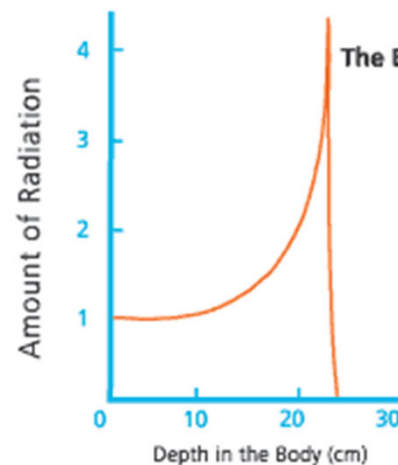


**CONVENTIONAL RADIATION THERAPY:**  
Deposits most energy before target

FIGURE 1

FIGURE 2

The downside of this strategy is that proton interactions with materials in the beamline will create high-energy secondary neutrons. The high linear energy transfer (LET) of neutrons makes them extremely efficient at ionization, and far more likely to cause cell death than low-LET particles, such as X-rays or protons.



# Thermal Neutron Capture Radiation Therapy

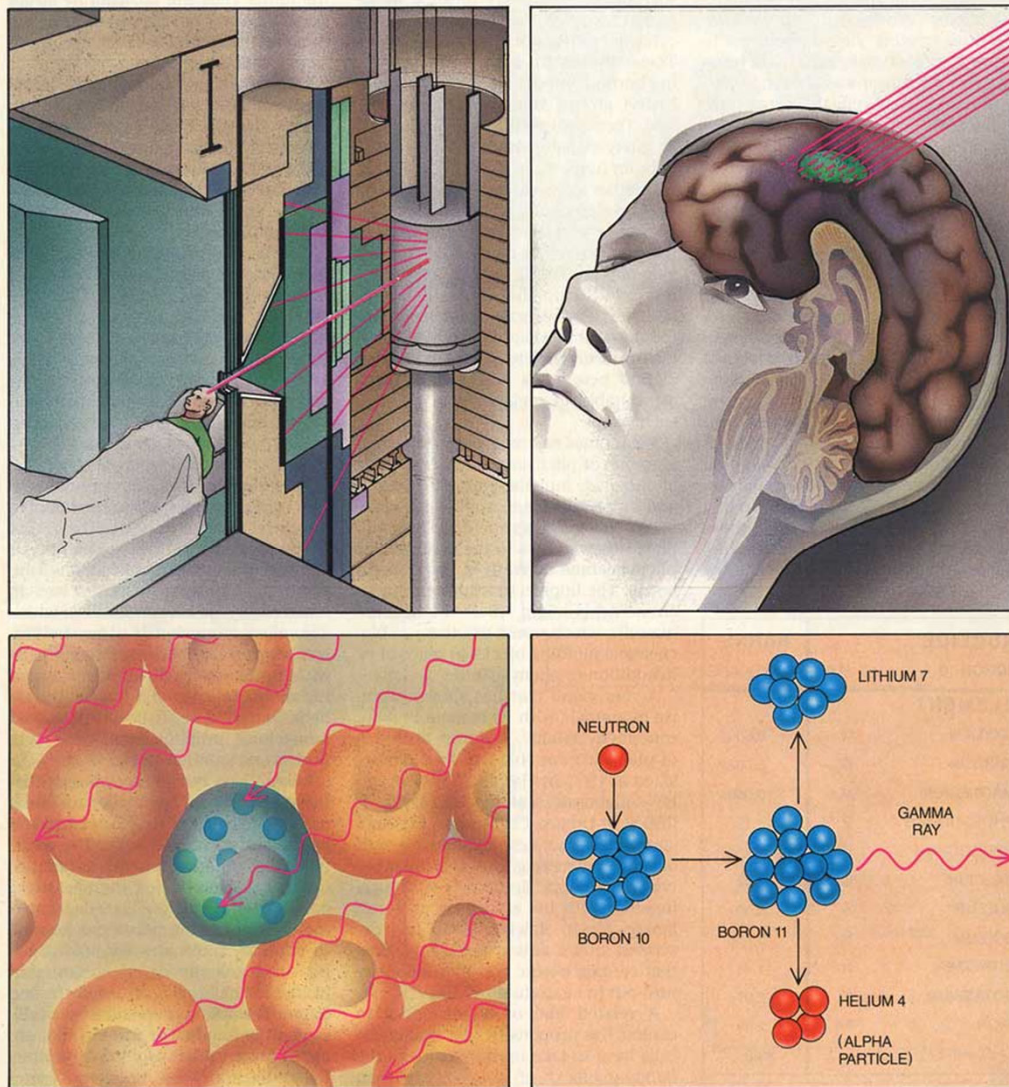


Fig.1 boron neutron capture therapy (BNCT) can be performed at a facility with a nuclear reactor or at hospitals that have developed alternative neutron sources. A beam of epithermal neutrons penetrates the brain tissue, reaching the malignancy. Once there the epithermal neutrons slow down and these low-energy neutrons combine with boron-10 (delivered beforehand to the cancer cells by drugs or antibodies) to form boron-11, releasing lethal radiation (alpha particles and lithium ions) that can kill the tumor.[1]

Figure from: <http://en.wikipedia.org/wiki/File:NeutronCaptureTherapyImage.jpg>

Barth, Rolf F.; Soloway, Albert H.; Fairchild, Ralph G. (1990). "Boron Neutron Capture Therapy for Cancer". *Scientific American* 263 (4): 100–3, 106–7. Bibcode:1990SciAm.263d.100B. doi:10.1038/scientificamerican1090-100. PMID 2173134. **182**





# Neutrons Sources



# Neutron Sources – Spontaneous Fission

## Neutron Sources – Spontaneous Fission



Cf-252 neutron source can be made extremely compact



An engineer tests the prototype Timed Neutron Detector, a device that detects landmines. The neutron source of the landmine detector holds a tiny amount of californium-252. (Photo credit: Pacific Northwest National Lab)

**TABLE 7.1 NEUTRON TERMINOLOGY**

Term	Energy Range	Velocity
Ultracold	$<2 \times 10^{-7}$ eV	6 m/s
Very cold	$2 \times 10^{-7}$ eV to $5 \times 10^{-5}$ eV	100 m/s
Cold neutrons	$5 \times 10^{-5}$ eV to 0.025 eV	—
Thermal <sup>c</sup>	0.025 eV	2200 m/s
Epithermal	1 eV–1 keV	$4.4 \times 10^5$ m/s
Cadmium	$<0.4$ eV	8800 m/s
Epicadmium	$>0.6$ eV	$1.1 \times 10^4$ m/s
Slow	$<1$ to 10 eV	$1.4 \times 10^4$ m/s
Resonance <sup>a</sup>	1 to 300 eV	$2.4 \times 10^5$ m/s
Intermediate	1 keV to 0.1 MeV	$4.4 \times 10^6$ m/s
Fast	$>0.1$ MeV	$1.4 \times 10^7$ m/s
Ultra fast (relativistic)	$>20$ MeV	—
Fission <sup>b</sup>	100 keV to 15 MeV	—

<sup>a</sup>In pile neutron physics usually refers to neutrons which are strongly captured in the resonance of U-238, and of a few commonly used detectors, e.g., In, Au.

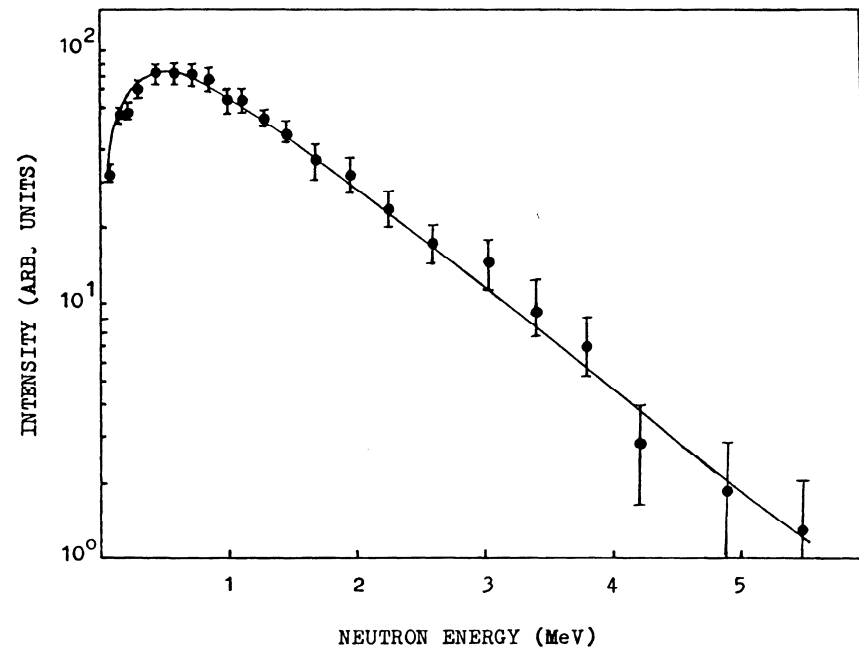
<sup>b</sup>Most probable energy 0.8 MeV, Average energy 2.0 MeV.

<sup>c</sup>Maxwellian distribution of 20°C extends to about 0.1 eV.

# Neutron Sources – Spontaneous Fission

Spontaneous fission of transuranic heavy nuclides, such as  $^{252}\text{Cf}$ , produces several fast neutrons, in addition to heavy fission products, prompt fission gamma rays and beta and gamma ray activities.

- Half-life: 2.65 years
- Neutron yield: 0.116n/s per Bq, or  $2.3 \times 10^6$  n/s per mg
- Neutron energy peaking at 0.5MeV and extends beyond 10MeV.



Measured neutron energy spectrum from spontaneous fission of  $^{252}\text{Cf}$





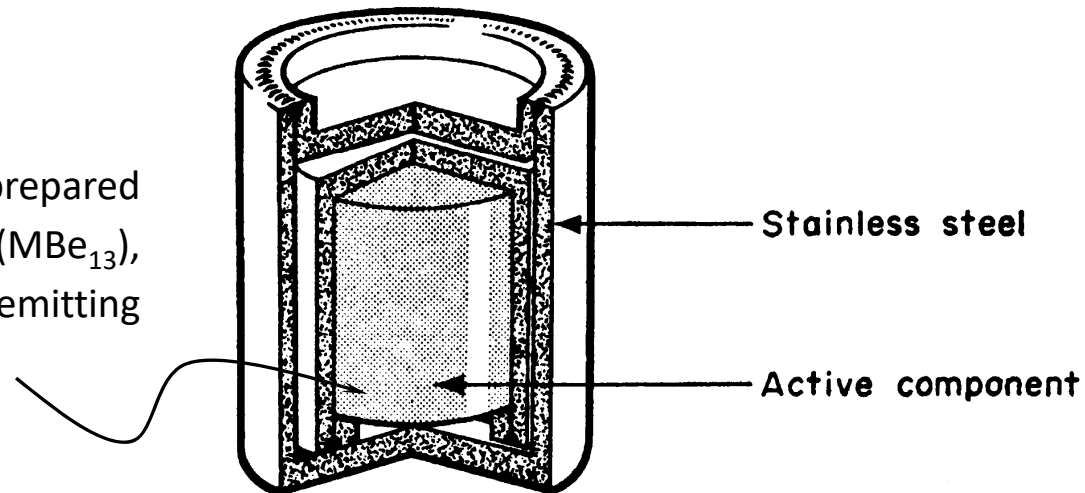
## Radioisotope ( $\alpha,n$ ) Sources

# Neutron Sources – Radioisotope ( $\alpha, n$ ) Sources

Energetic alpha particles can induce ( $\alpha, n$ ) reaction in certain target materials.

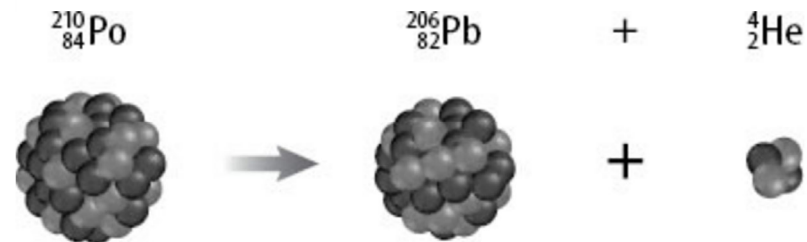


The source is normally prepared in the form of alloy ( $\text{MBe}_{13}$ ), where M is alpha-emitting radioisotopes



A practical neutron source

# Neutron Sources – Radioisotope ( $\alpha,n$ ) Sources

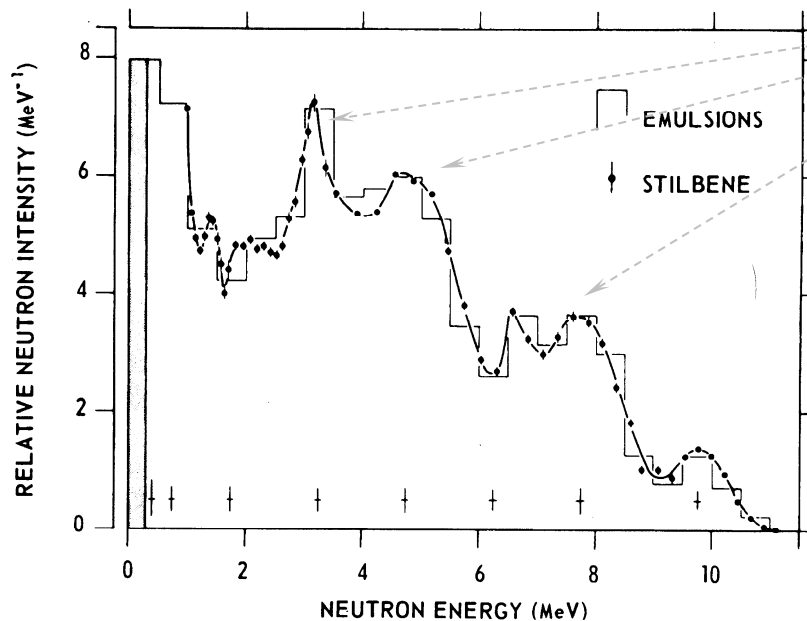


**TABLE 9.2. ( $\alpha,n$ ) Neutron Sources**

Source	Average Neutron Energy (MeV)	Half-life
$^{210}\text{PoBe}$	4.2	138 d
$^{210}\text{PoB}$	2.5	138 d
$^{226}\text{RaBe}$	3.9	1600 y
$^{226}\text{RaB}$	3.0	1600 y
$^{239}\text{PuBe}$	4.5	24100 y

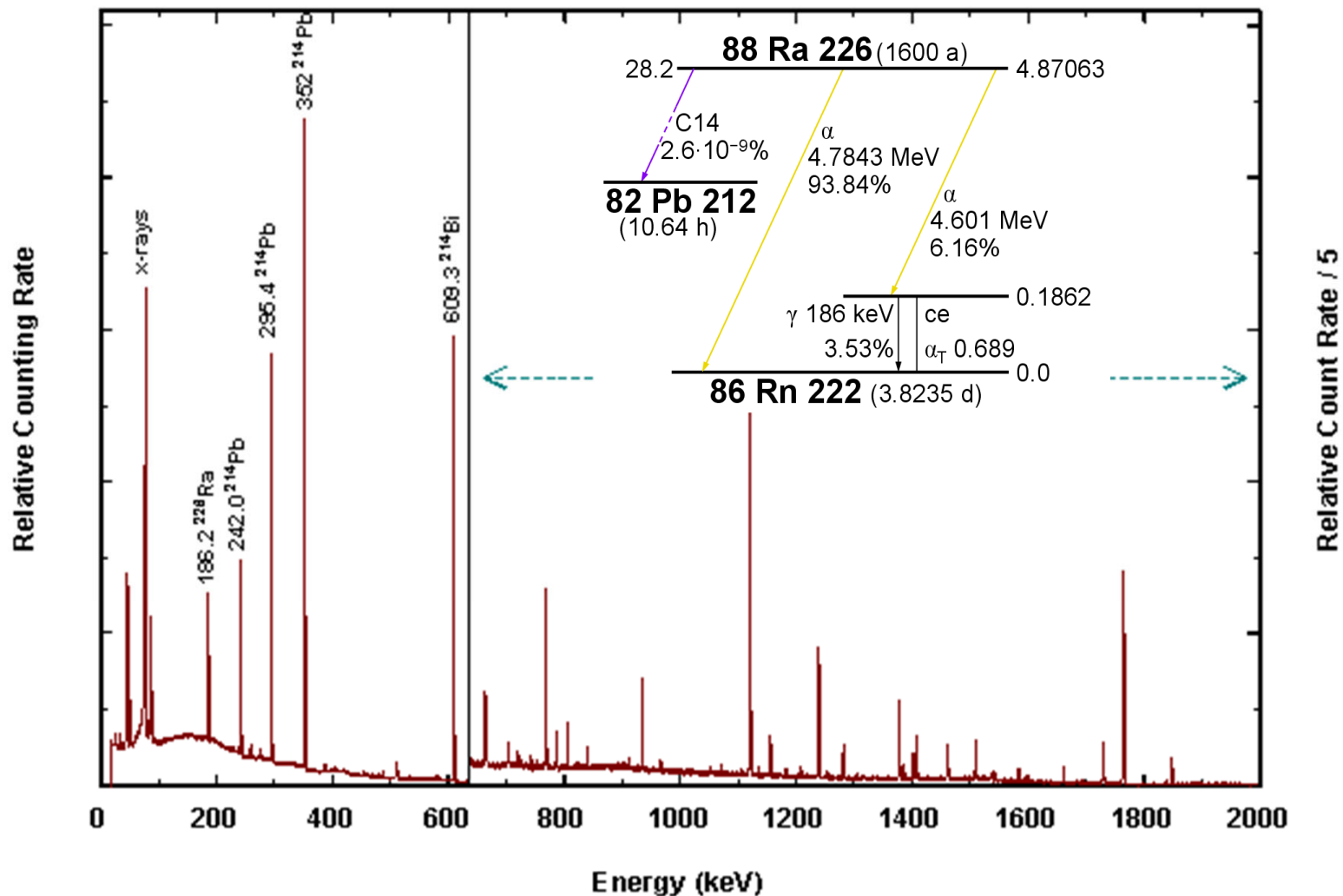
# Neutron Sources – Radioisotope ( $\alpha, n$ ) Sources

A typical neutron energy spectrum from an  $^{239}\text{Pu}/\text{Be}$  source.



- The various peak and valley are due to the distinct excited states of the  $^{12}\text{C}$  product nucleus.
- The continuum is the result of variable energy possessed by the alpha particles before reaction.

# Neutron Sources – Radioisotope ( $\alpha, n$ ) Sources



Radium-226 gamma ray spectrum from high purity germanium (HPGe) detector





## Neutron Sources – Radioisotope ( $\alpha$ ,n) Sources

Practical considerations for choosing appropriate  $\alpha$  emitter.

- Radioisotope ( $\alpha$ ,n) sources are normally associated with other significant background radiations, especially when  $^{226}\text{Ra}$  and  $^{227}\text{Ac}$  are used.
- Choice has to be made between specific activity of the alpha emitter (and therefore neutron yield), source life-time and the availability of the isotope.

## Neutron Sources – Photon-Neutron Sources

- Some radioisotope gamma ray emitters can also be used to produce neutrons when combined with an appropriate target material.



- A gamma ray photon with an energy greater than the negative of the Q-value is required.
- Some practical gamma ray emitter include:  ${}^{226}\text{Ra}$ ,  ${}^{124}\text{Sb}$ ,  ${}^{72}\text{Ga}$ ,  ${}^{140}\text{La}$  and  ${}^{24}\text{Na}$ .

## Neutron Sources – Photo-neutron Sources

If the gamma rays are monoenergetic, the neutrons are also nearly monoenergetic!

$$E_n(\theta) \cong \frac{M(E_\gamma + Q)}{m + M} + \frac{E_\gamma [(2mM)(m + M)(E_\gamma + Q)]^{1/2}}{(m + M)^2} \cos(\theta)$$

where

$\theta$  = angle between gamma photon and neutron direction

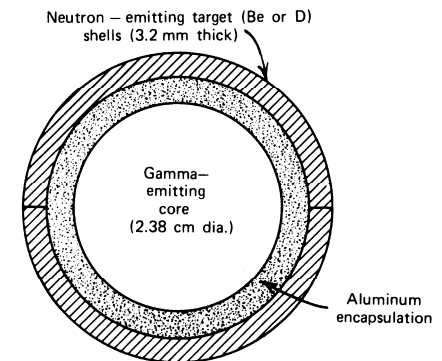
$E_\gamma$  = gamma energy

$M$  = mass of recoil nucleus  $\times c^2$

$m$  = mass of neutron  $\times c^2$

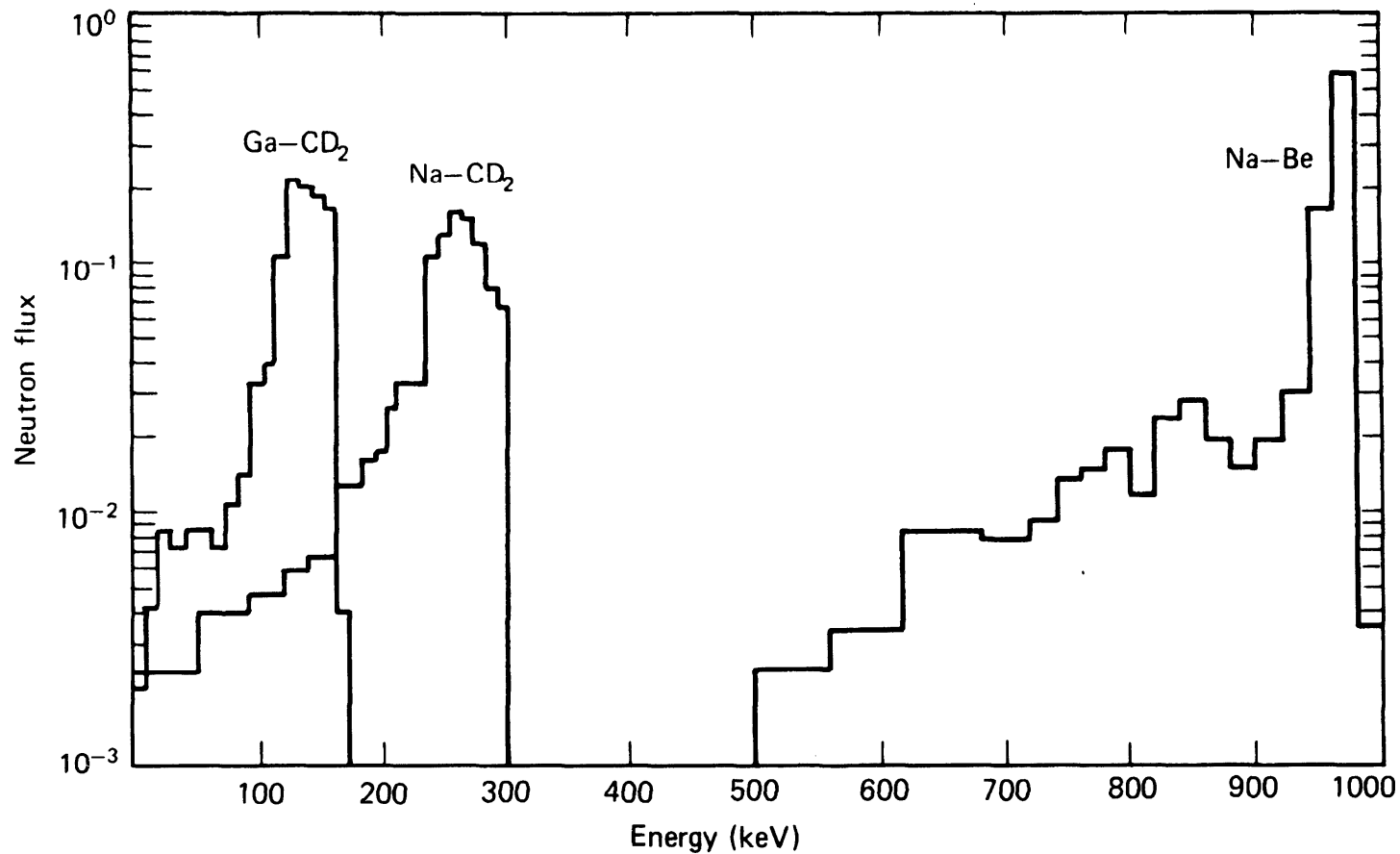
The neutron energy is blurred by

- The slight angular dependency.
- Neutron scattering inside the source.



Typical structure of photon neutron sources

# Neutron Sources – Photo-neutron Sources



Calculated neutron energy spectra

**TABLE 7.2 SOURCES OF NEUTRONS** (From ICRU 26, 1977)

Source	Reaction	Half Life	Average Neutron Energy (MeV)	Yield n/s/Ci	Character Problems
Mock Fission (Po+Be+B+Li+F)	$\alpha, n$	134.4 d	Fission spectrum	$4 \times 10^5$	$\alpha$
$^{24}\text{Na} + \text{Be}$	$\gamma, n$	15 h	0.83	$1.3 \times 10^5$	$\gamma$
$^{24}\text{Na} + \text{D}_2\text{O}$	$\gamma, n$	15 h	0.22	$2.7 \times 10^5$	$\gamma$
$^{56}\text{Mn} + \text{Be}$	$\gamma, n$	2.58 h	0.1 (90%) 0.3 (10%)	$2.9 \times 10^4$	$\gamma$
$^{56}\text{Mn} + \text{D}_2\text{O}$	$\gamma, n$	2.58 h	0.22	$3.1 \times 10^3$	$\gamma$
$^{72}\text{Ga} + \text{Be}$	$\gamma, n$	14.1 h	0.78	$5 \times 10^4$	$\gamma$
$^{72}\text{Ga} + \text{D}_2\text{O}$	$\gamma, n$	14.1 h	0.13	$6 \times 10^4$	$\gamma$
$^{88}\text{Y} + \text{Be}$	$\gamma, n$	107 d	0.16	$1 \times 10^5$	$\gamma$
$^{88}\text{Y} + \text{D}$	$\gamma, n$	107 d	0.31	$3 \times 10^3$	$\gamma$
$^{116}\text{In} + \text{Be}$	$\gamma, n$	14 s	0.30	$8.2 \times 10^3$	$\gamma$
$^{124}\text{Sb} + \text{Be}^b$	$\gamma, n$	60.2 d	0.024	$1.9 \times 10^5$	$\gamma$
$^{140}\text{La} + \text{Be}$	$\gamma, n$	40.3 h	0.62	$3 \times 10^3$	$\gamma$
$^{140}\text{La} + \text{D}_2\text{O}$	$\gamma, n$	40.3 h	0.15	$8 \times 10^3$	$\gamma$
$^{228}\text{Ra} + \text{Be}$	$\gamma, n$	5.75 y	0.83	$3.5 \times 10^4$	$\gamma$
$^{228}\text{Ra} + \text{D}_2\text{O}$	$\gamma, n$	5.75 y	0.20	$9.5 \times 10^4$	$\gamma$
$^{226}\text{Ra} + \text{Be}$	$\alpha, n$	1600 y	Spectrum	$3.0 \times 10^4$	$\alpha, \gamma, \text{Rn}$
$^{226}\text{Ra} + \text{Be}$	$\alpha, n$	1600 y	5.0	$1.7 \times 10^7$	$\alpha, \gamma, \text{Rn}$
$^{226}\text{Ra} + \text{B}$	$\gamma, n$	1600 y	3.0	$6.8 \times 10^6$	$\alpha, \gamma, \text{Rn}$
$^{226}\text{Ra} + \text{D}_2\text{O}$	$\alpha, n$	1600 y	0.12	$1 \times 10^3$	$\alpha, \gamma, \text{Rn}$
$^{222}\text{Rn} + \text{Be}$	$\alpha, n$	3.82 d	5	$1.5 \times 10^7$	$\alpha, \gamma, \text{Rn}$
$^{210}\text{Po} + \text{Be}$	$\alpha, n$	134.4 d	4	$3 \times 10^6$	$\alpha$
$^{210}\text{Po} + \text{B}$	$\alpha, n$	134.4 d	2.5	$9 \times 10^5$	$\alpha$
$^{210}\text{Po} + \text{F}$	$\alpha, n$	134.4 d	1.4	$4 \times 10^5$	$\alpha$
$^{210}\text{Po} + \text{Li}$	$\alpha, n$	134.4 d	0.42	$9 \times 10^4$	$\alpha$
$^{227}\text{Ac} + \text{Be}$	$\alpha, n$	21.8 y	—	—	$\alpha$
$^{238}\text{Pu} + \text{Be}$	$\alpha, n$	87.7 y	4.5	$2.3 \times 10^6$	$\alpha$
$^{239}\text{Pu} + \text{Be}$	$\alpha, n$	$2.41 \times 10^4$ y	4 (3.2)	$1.7 \times 10^6$	$\alpha$
$^{241}\text{Am} + \text{Be}$	$\alpha, n$	432 y	4.5	$2.2 \times 10^6$	$\alpha$
$^{241}\text{Am} + \text{Li}$	$\alpha, n$	432 y	0.54	$6.0 \times 10^4$	$\alpha$
$^{239}\text{Pu} (\text{WG})^c$	Spon. Fission	$2.41 \times 10^4$ y	1.94	63.6	$\alpha$
$^{252}\text{Cf}$	Spon. Fission	2.64 y	Fission spectrum <sup>a</sup> (2.35)	$10^6$	$\alpha$

<sup>a</sup> $3.80 \pm 0.035$  neutrons per fission.

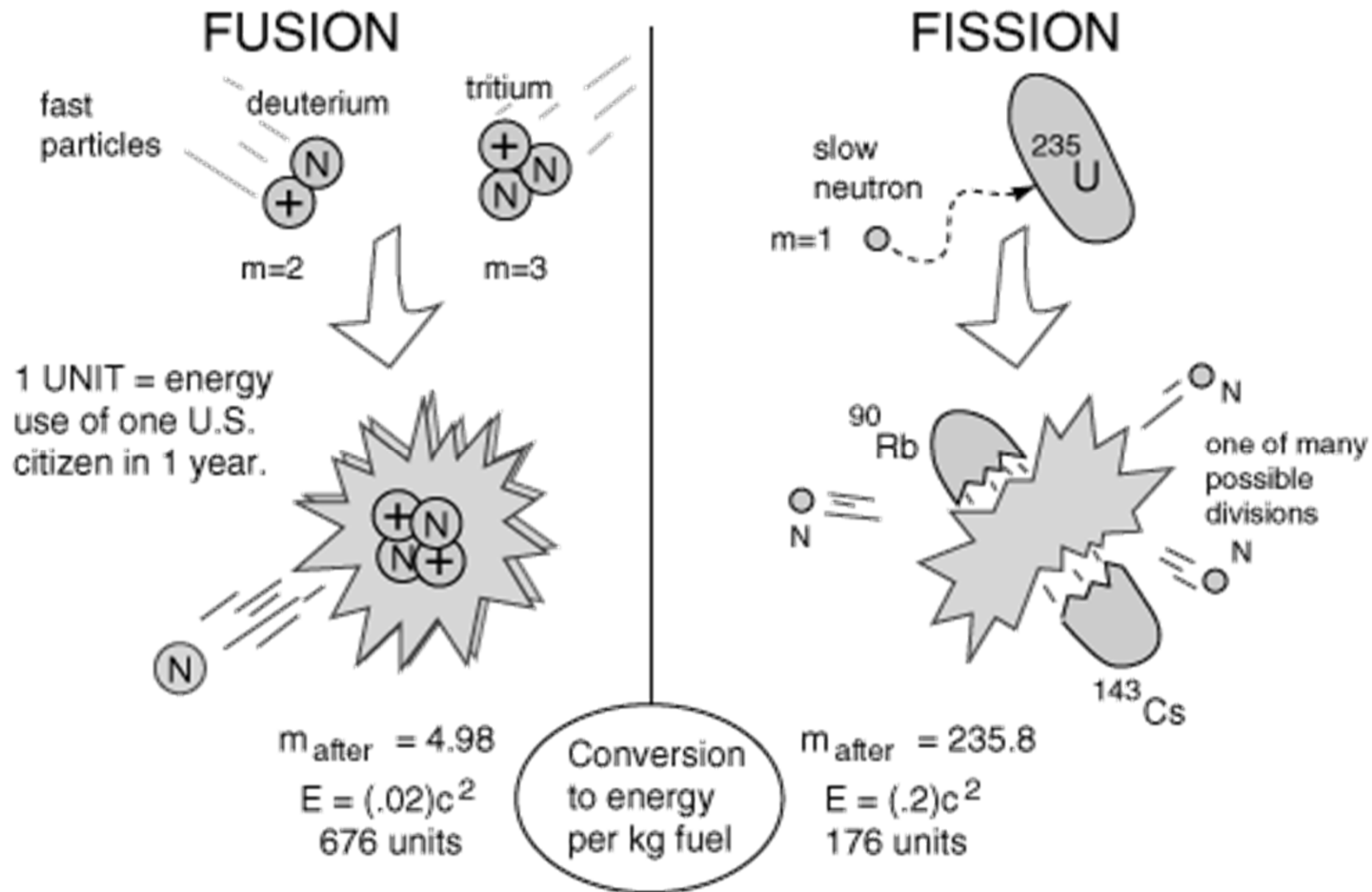
<sup>b</sup>Typically used in reactors – inserted as  $^{123}\text{Sb}$  and resultant activation to  $^{124}\text{Sb}$  occurs.

<sup>c</sup>WG = Weapons grade, >93% Pu-239.



# Average Binding Energy Per Nucleon

## Comparing Fusion and Fission Reactions



<http://230nsc1.phy-astr.gsu.edu/hbase/hframe.html>



# **Neutrons Generated by Accelerated Charged Particles**

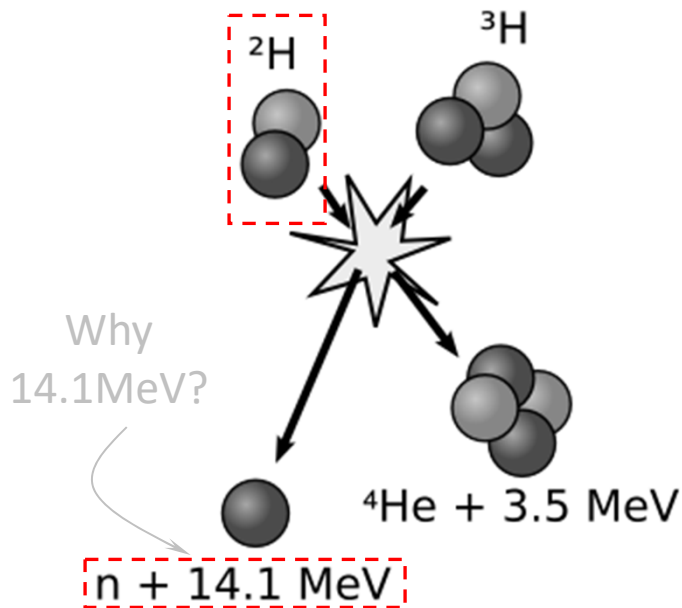
# Neutrons Generated by Accelerated Charged Particles

- Neutrons can be produced by nuclear reaction between accelerated charged particles.

The D-D reaction:  ${}^2_1\text{H} + {}^2_1\text{H} \Rightarrow {}^3_2\text{He} + {}^1_0\text{n}$ , Q-value: 3.26MeV,  $E_n=2.5\text{MeV}$

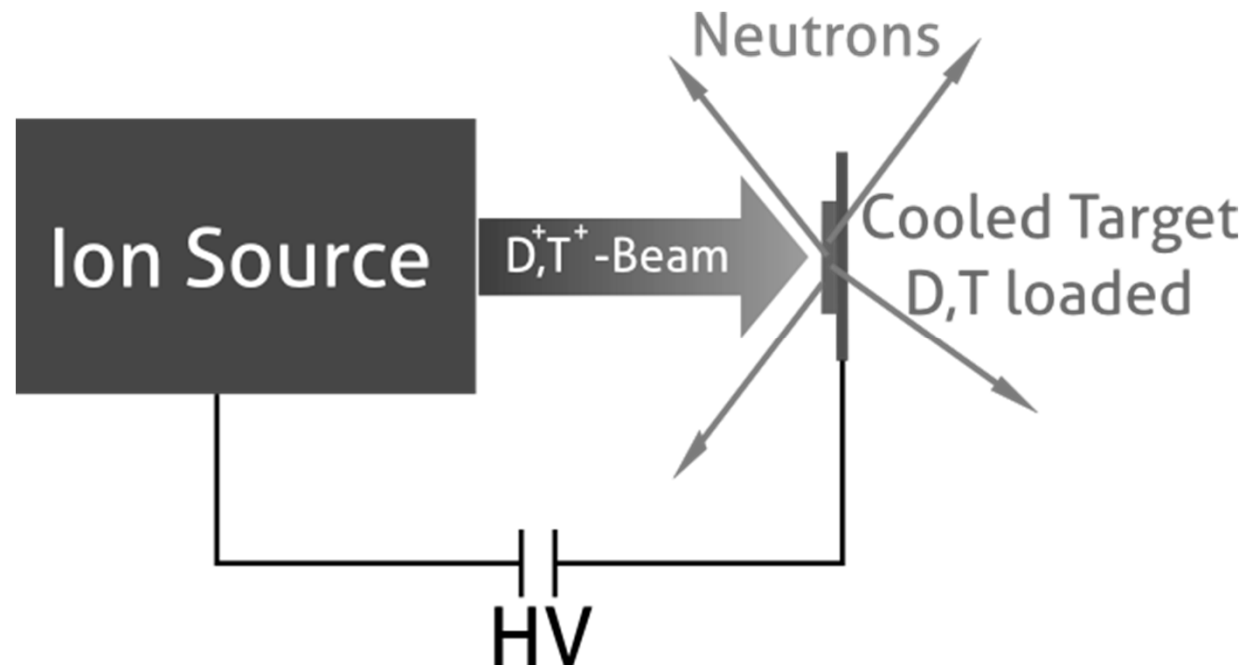
The D-T reaction:  ${}^2_1\text{H} + {}^3_1\text{H} \Rightarrow {}^4_2\text{He} + {}^1_0\text{n}$ , Q-value: 17.6MeV,  $E_n=14.1\text{MeV}$

Why accelerated?



- Due to the Coulomb barrier between the incident deuteron and the light target nucleus, a relatively small accelerating potential is required (about 100 to 300kV) to induce the reaction.
- The neutrons produced by a given nuclear reaction (D-D or D-T) have roughly the same energies.

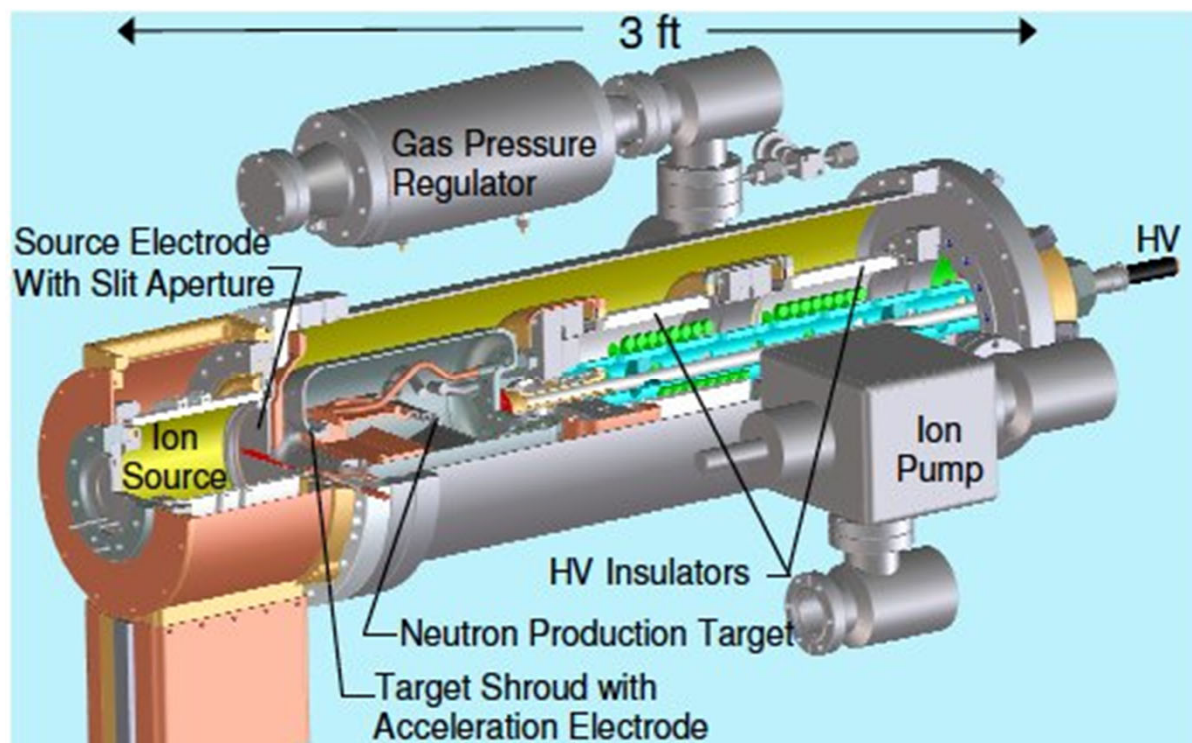
# Neutrons Generated by Accelerated Charged Particles



A schematic of a neutron generator. On the left hand side is an ion source from which a Hydrogen isotope ion beam is extracted. This beam is accelerated with a high voltage towards a target on the right hand side, loaded with Deuterium and Tritium, to produce neutrons in a nuclear fusion process.

<http://ibt.lbl.gov/neutrongamma.html>

## A Typical D-T Neutron Generator



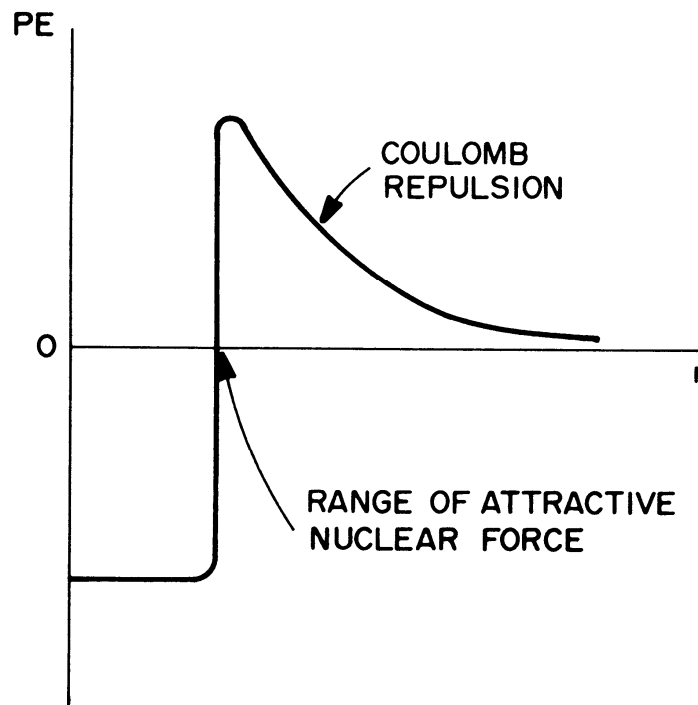
This compact and simple device can generate  $5 \times 10^{11}$  neutrons per second by accelerating deuterium or tritium (depending on the desired neutron spectrum) into a deuterated or tritiated neutron production target.

<http://ibt.lbl.gov/neutrongamma.html>

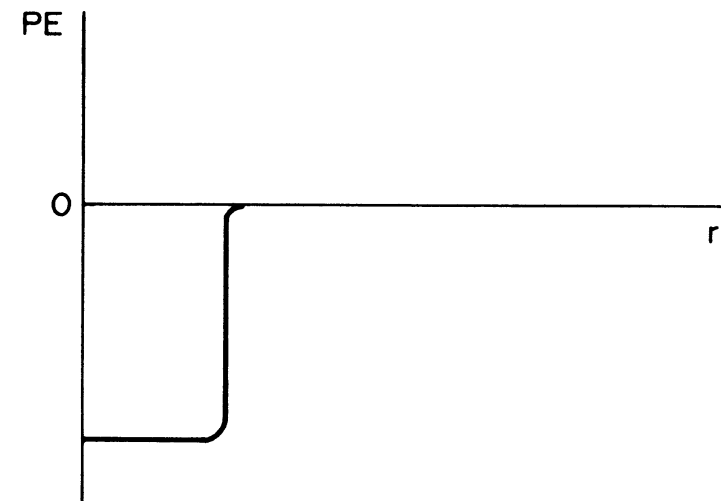


# Coulomb Barrier

- Nucleons are bounded together in nucleus by the strong force, which has a short range of  $\sim 10^{-15}\text{m}$ .
- The strong force is powerful enough to overcome the Coulomb repulsion between the positively charged protons.



(a) PROTON - NUCLEUS



(b) NEUTRON - NUCLEUS



## Neutrons Generated by Accelerated Charged Particles

**TABLE 9.1. Reactions Used to Produce Monoenergetic Neutrons with Accelerated Protons (p) and Deuterons (d)**

Reaction	Q Value (MeV)
${}^3\text{H}(\text{d},\text{n}){}^4\text{He}$	17.6
${}^2\text{H}(\text{d},\text{n}){}^3\text{He}$	3.27
${}^{12}\text{C}(\text{d},\text{n}){}^{13}\text{N}$	−0.281
${}^3\text{H}(\text{p},\text{n}){}^3\text{He}$	−0.764
${}^7\text{Li}(\text{p},\text{n}){}^7\text{Be}$	−1.65

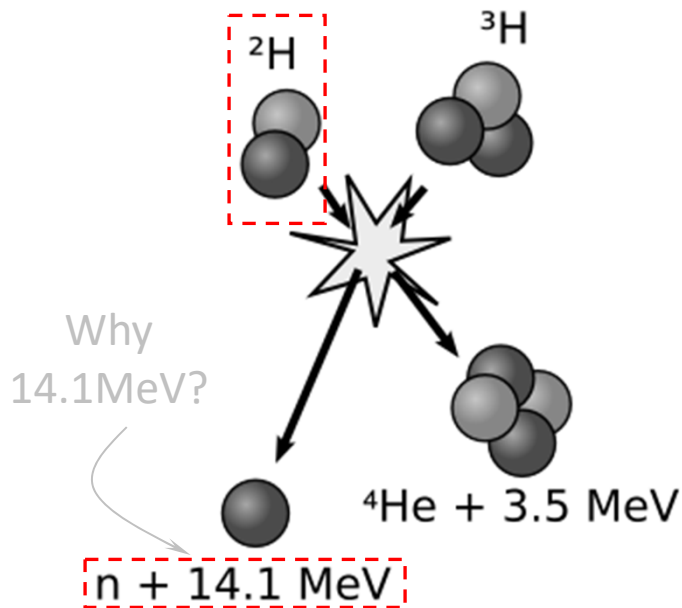
# Neutrons Generated by Accelerated Charged Particles

- Neutrons can be produced by nuclear reaction between accelerated charged particles.

The D-D reaction:  ${}^2_1\text{H} + {}^2_1\text{H} \Rightarrow {}^3_2\text{He} + {}^1_0\text{n}$ , Q-value: 3.26MeV,  $E_n=2.5\text{MeV}$

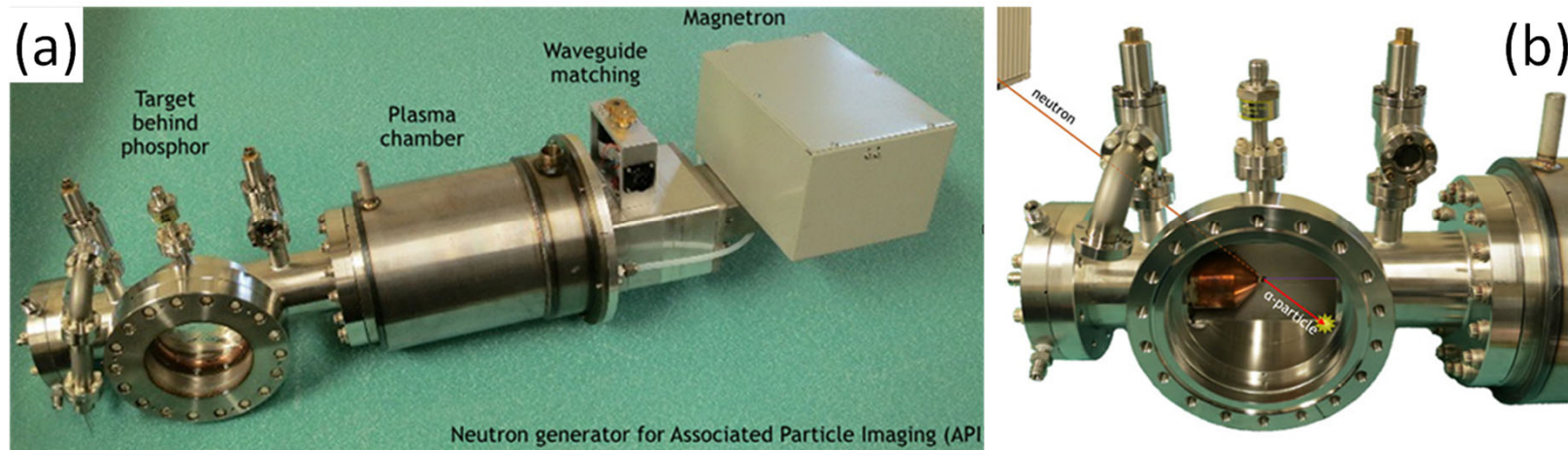
The D-T reaction:  ${}^2_1\text{H} + {}^3_1\text{H} \Rightarrow {}^4_2\text{He} + {}^1_0\text{n}$ , Q-value: 17.6MeV,  $E_n=14.1\text{MeV}$

Why accelerated?

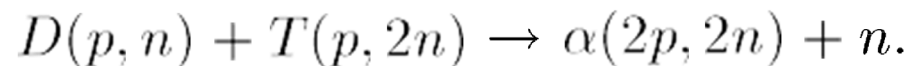


- Due to the coulomb barrier between the incident deuteron and the light target nucleus, only a relatively small accelerating potential is required (about 100 to 300kV) to induce the reaction.
- The neutrons produced by a given nuclear reaction (D-D or D-T) have roughly the same energies.

# NDE of Potential Misloading with Time-Tagged Neutron Interrogation Techniques (Technical Approach)



(a) The commercial associate particle neutron generator (DT108API) from Adelphi Technology. (b) The internal assembly of the anode and the alpha-detector inside the APNG. Both images from [23].





# Large Sized Neutron Sources

## Nuclear fission reactors

Nuclear fission which takes place within in a reactor produces very large quantities of neutrons and can be used for a variety of purposes including power generation and experiments.

## Nuclear fusion systems

Nuclear fusion, the combining of the heavy isotopes of hydrogen, also has the potential to produces large quantities of neutrons. Small scale fusion systems exist for research purposes at many universities and laboratories around the world. A small number of large scale nuclear fusion systems also exist including the National Ignition Facility in the USA, JET in the UK, and soon the recently started ITER experiment in France.

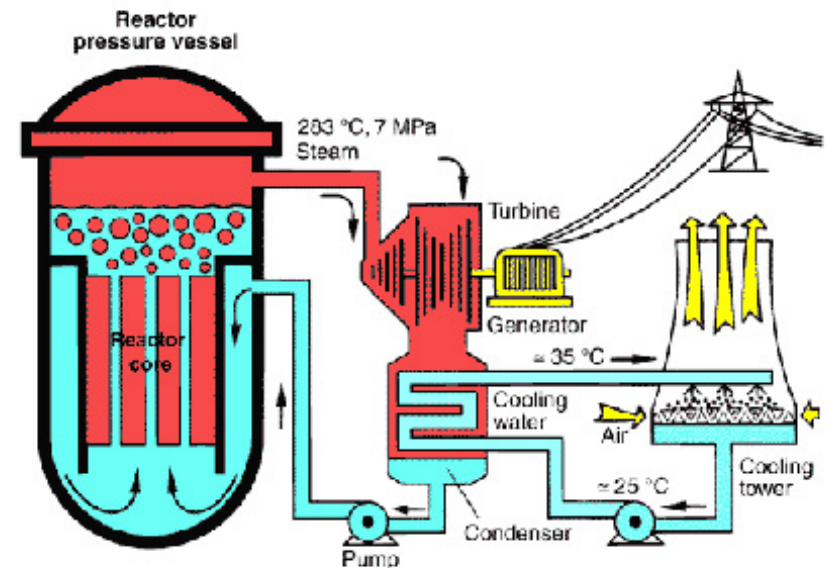
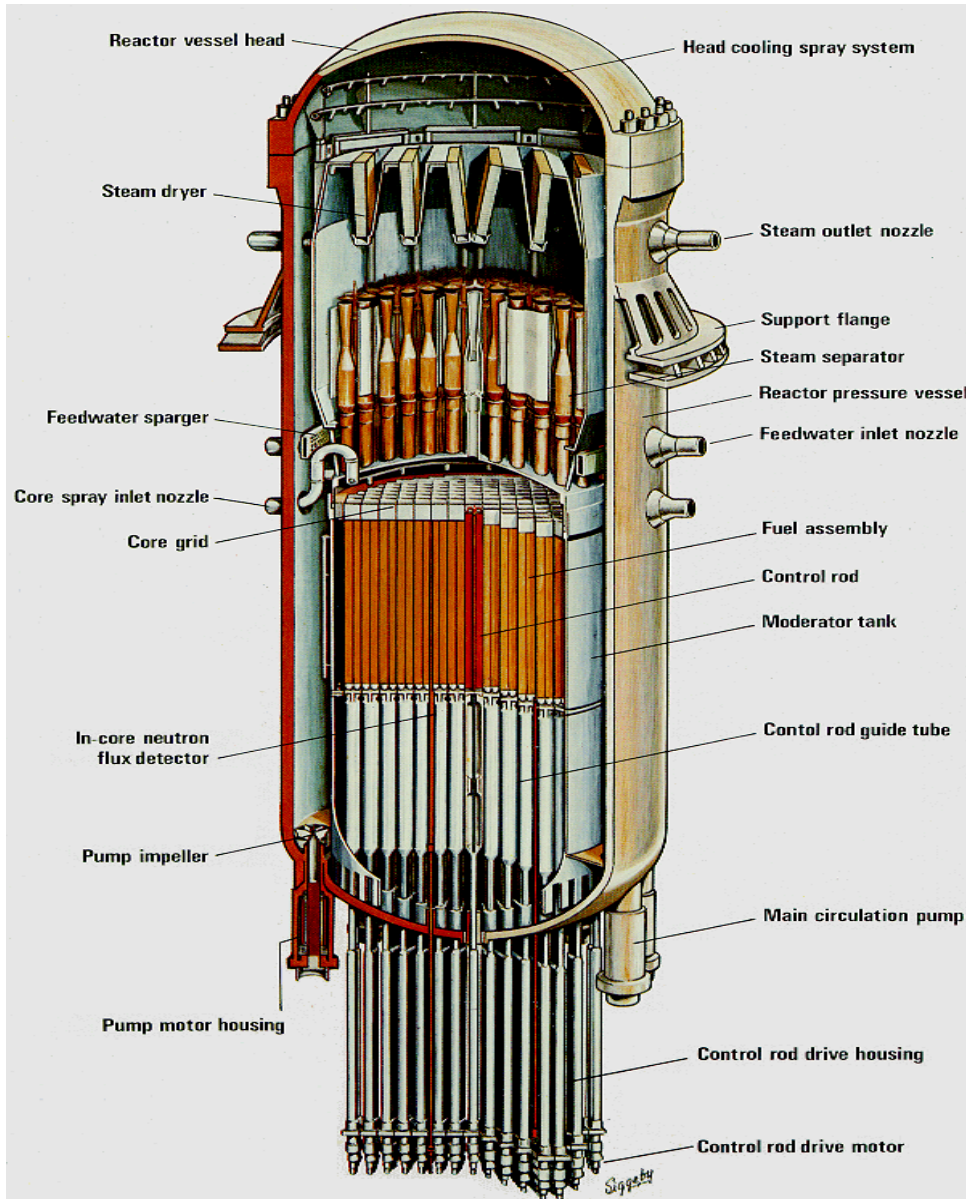
## High energy particle accelerators

A spallation source is a high-flux source in which protons that have been accelerated to high energies hit a target material, prompting the emission of neutrons.



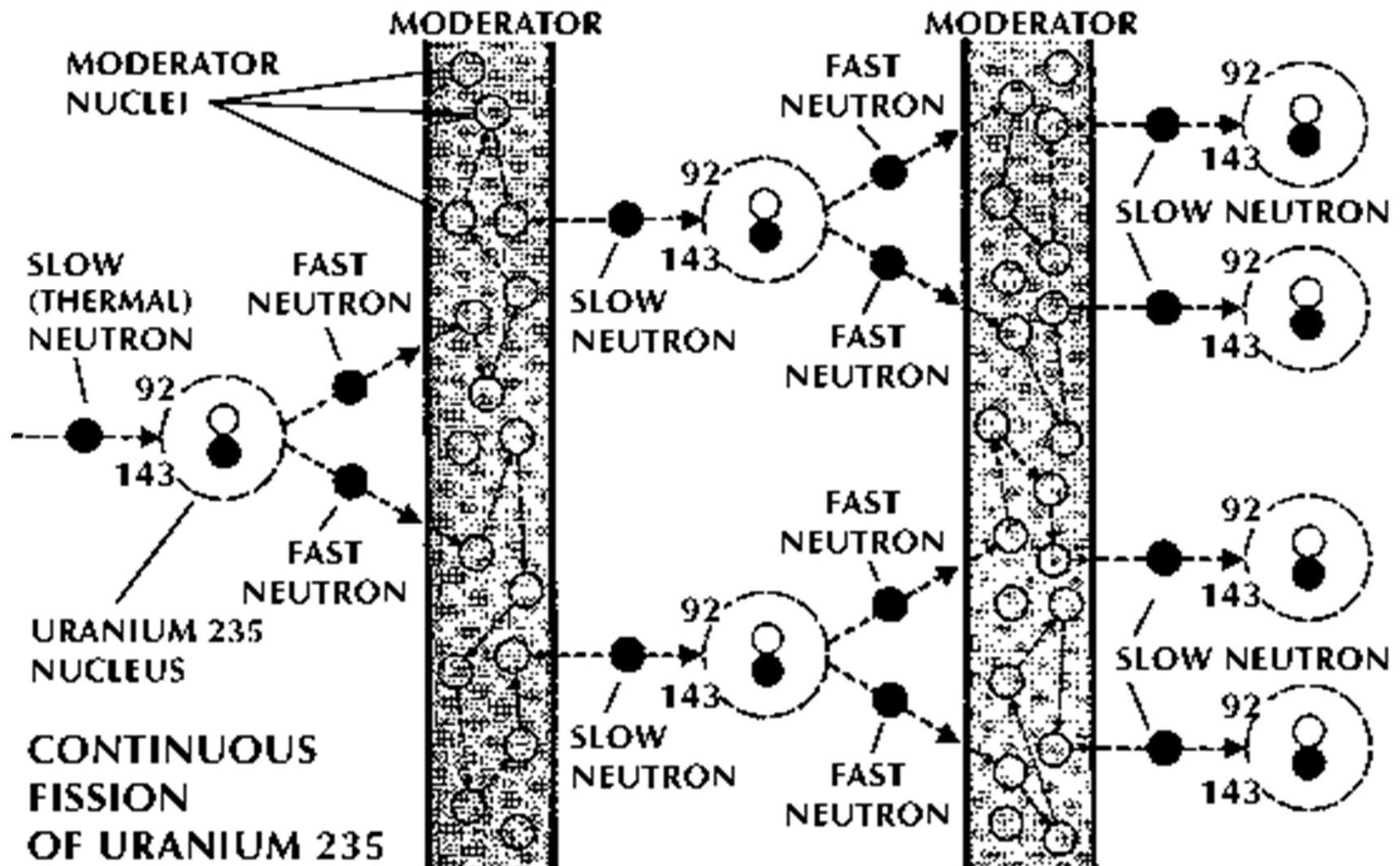
# Nuclear Reactor

# Inside a Nuclear Reactor

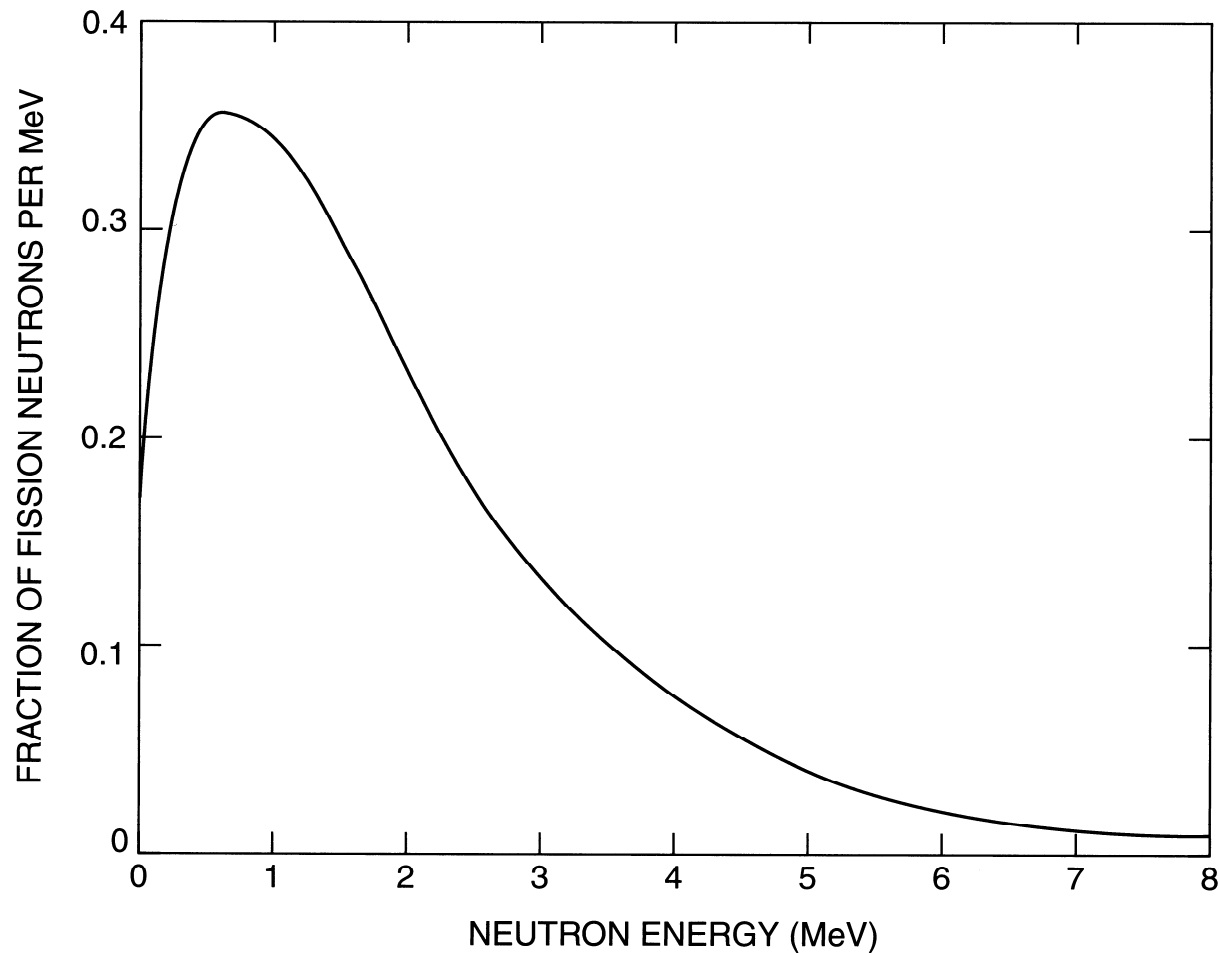




# Nuclear Fission from Slow Neutrons and Water Moderator



# Nuclear Fission from Slow Neutrons and Water Moderator



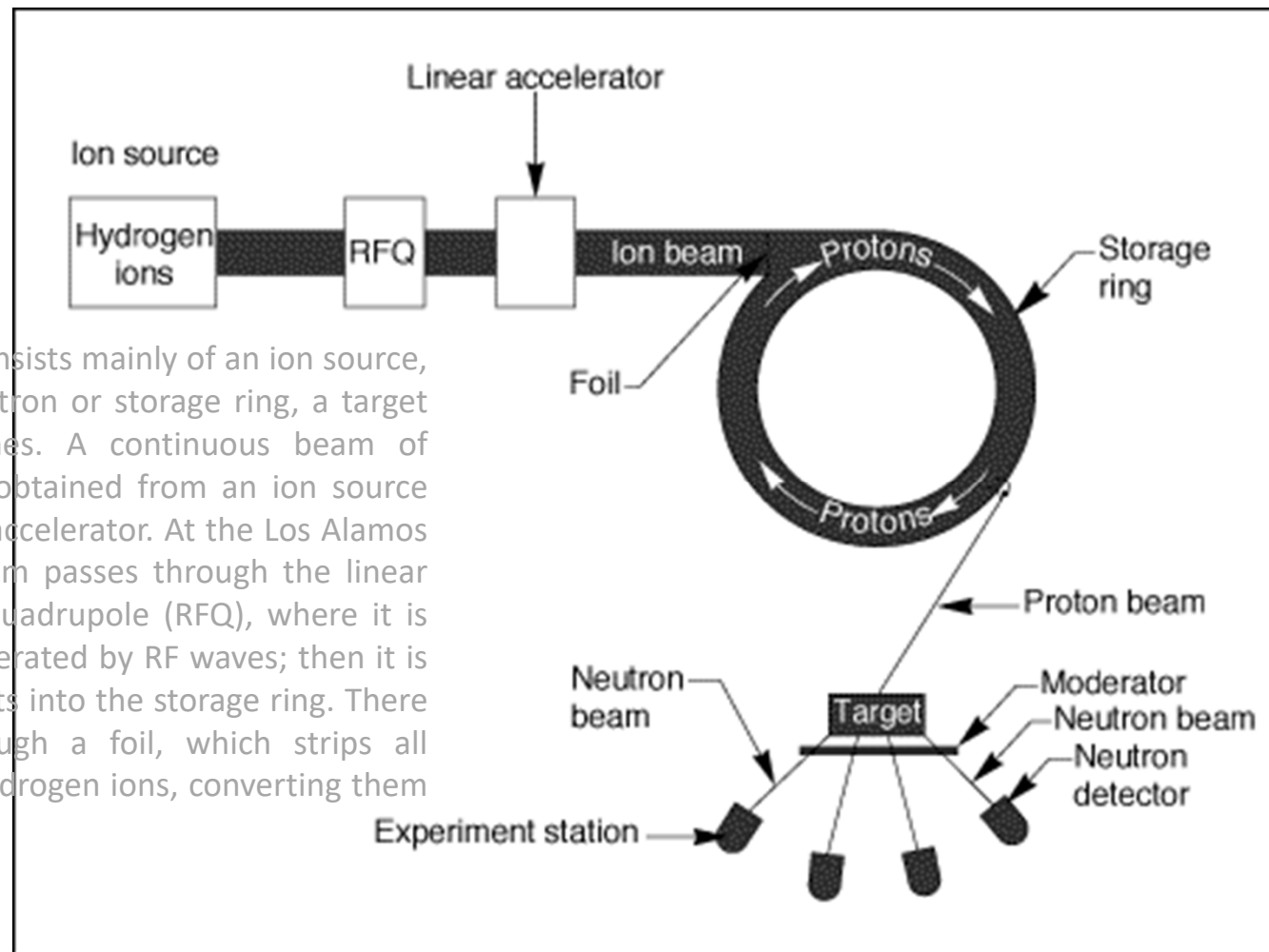
**FIGURE 7.3** ♦ Prompt Neutron Energy Spectrum for Thermal Fission of U-235.  
(from DOE-HDBK-1019/1-93)



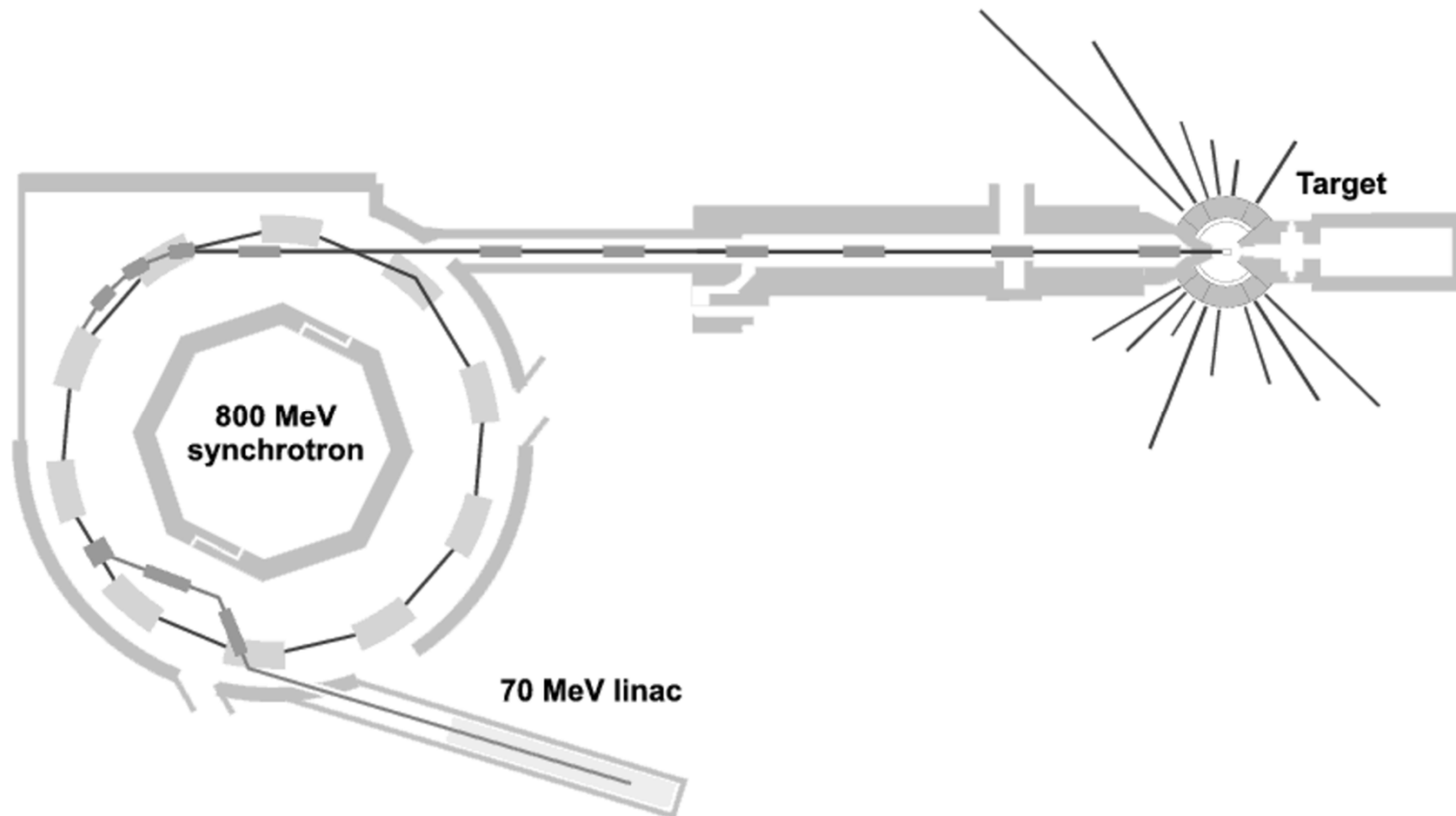
# Spallation Neutron Sources

# Spallation Neutron Sources

A spallation neutron source consists mainly of an ion source, a linear accelerator, a synchrotron or storage ring, a target area, and neutron beam lines. A continuous beam of negative ions of hydrogen is obtained from an ion source and accelerated by the linear accelerator. At the Los Alamos spallation source, the ion beam passes through the linear accelerator's radiofrequency quadrupole (RFQ), where it is bunched into pulses and accelerated by RF waves; then it is focused by quadrupole magnets into the storage ring. There the negative ions pass through a foil, which strips all electrons from the negative hydrogen ions, converting them into positive ions, or protons.



# Spallation Neutron Sources



A schematic diagram of a pulsed source (based on ISIS, UK)