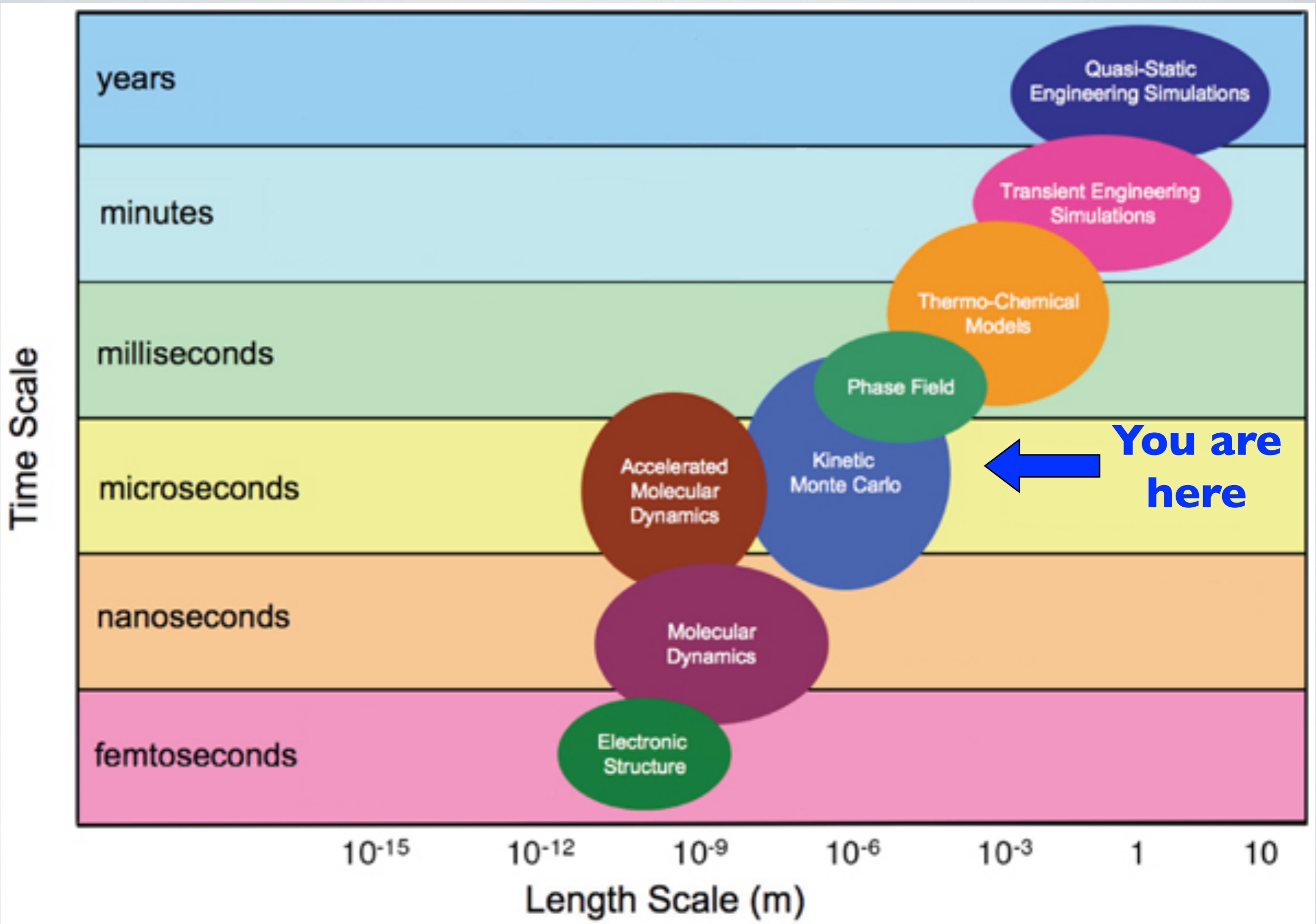


MODULE 3: DISLOCATION DYNAMICS

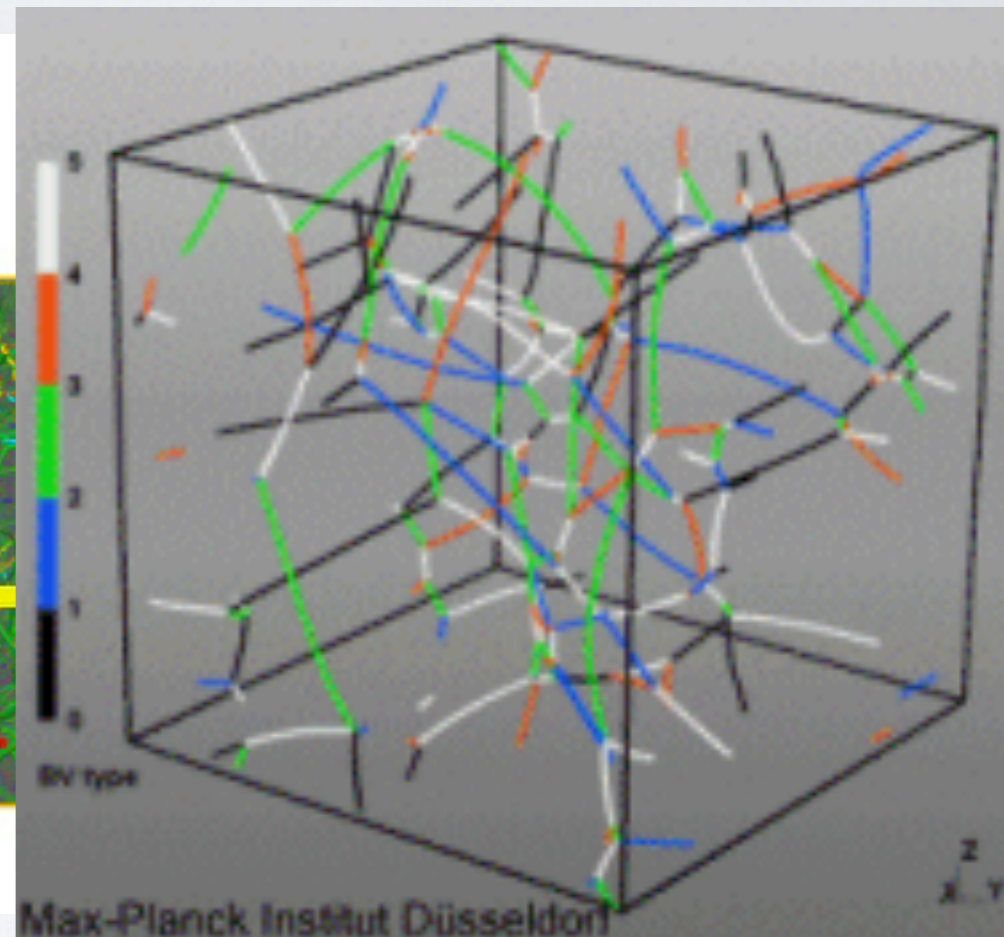
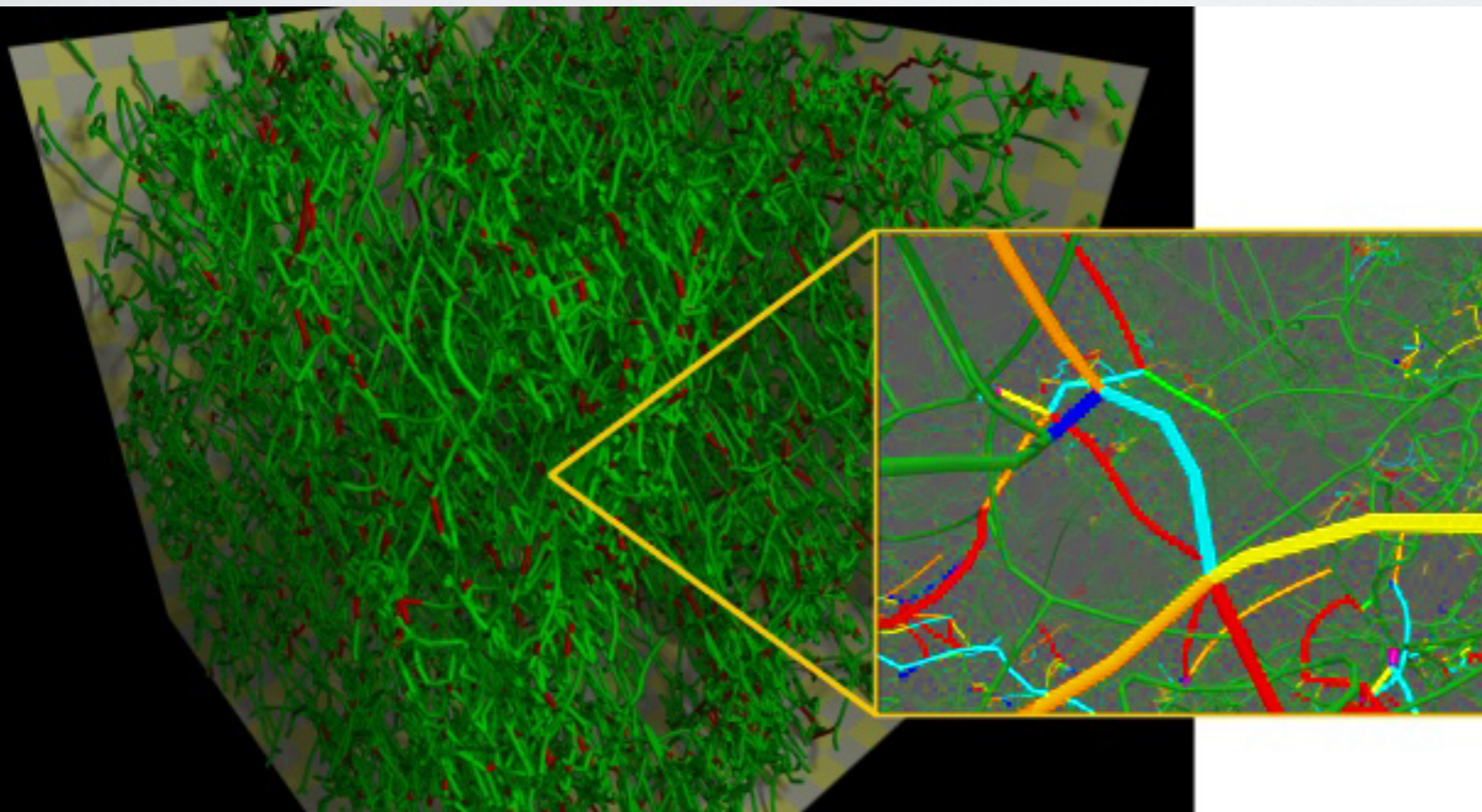
Principles and Theory

I. Introduction



What is dislocation dynamics?

- A (predictive?) guide to the evolution of materials microstructure
- An experiment on a computer
- A simulation of the “classical” dynamics of dislocations

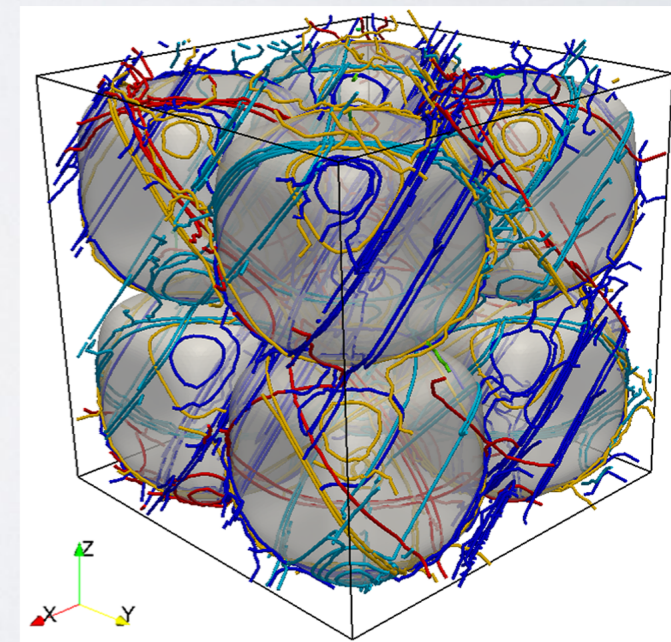
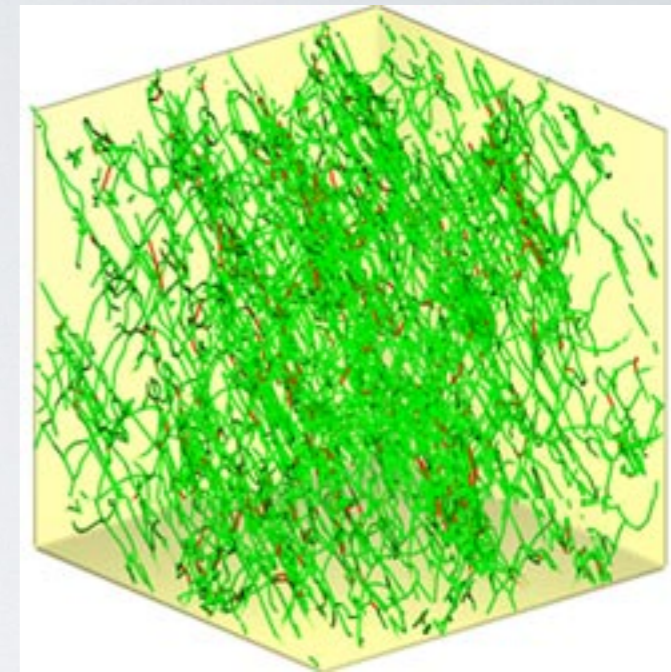


Why is it useful?

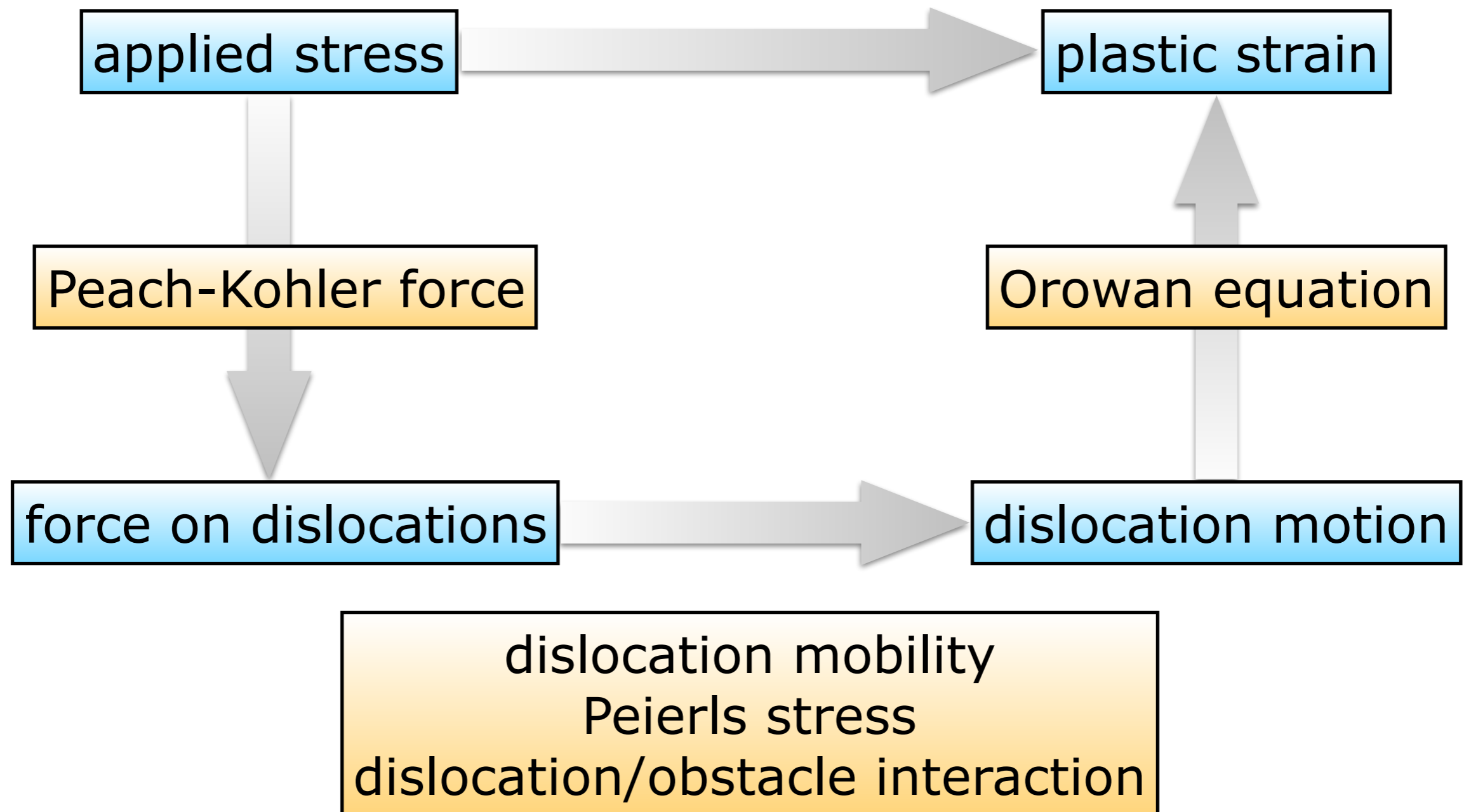
- By simulating dislocation motion, interaction, and growth, we can gain **mesoscale insight** into dislocation structure
- Advanced experimental techniques (TEM, high-energy X-ray diffraction) can resolve dislocation structure, but not dynamics
- We can **predict and understand** dislocation behavior and compare / interpret experimental observations
- Total control of dislocation interactions and initial conditions
- Nanomechanical measurement approaches scale of DD

What is it used for?

- Bridge from atomistic to mesoscale
 - dislocation behavior and interaction scale up
- Understanding strain hardening
 - dislocation density growth with strain & importance of interaction mechanisms
- Examining small-scale plasticity
 - micro- and nano pillars show unusual plastic behavior due to dislocations
- Parameterizing larger-scale models
 - dislocation-density-based crystal plasticity



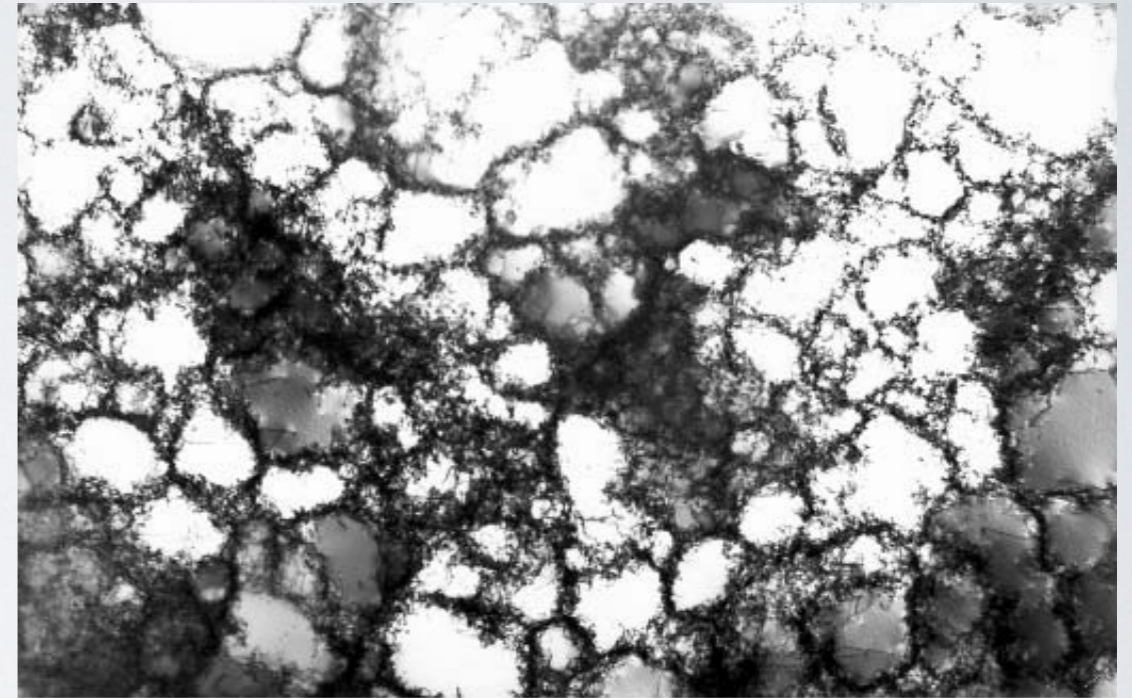
Dislocations and plastic deformation



II. History

First DDD simulations

- Robert Amodeo and Nasr Ghoniem (1997-1998) start direct numerical simulations of interacting dislocations to understand **cell formation**
- Range of efforts worldwide:
 - Cai/Bulatov/Arsenlis
 - Devincere
 - Ghoniem/Al Ezab
 - ...
- In addition, there are other approaches that attempt to model dislocations as entities:
 - Phase-field methods
 - Level-set methods
- All of the **discrete dislocation dynamics** methods benefit heavily from **parallelization**



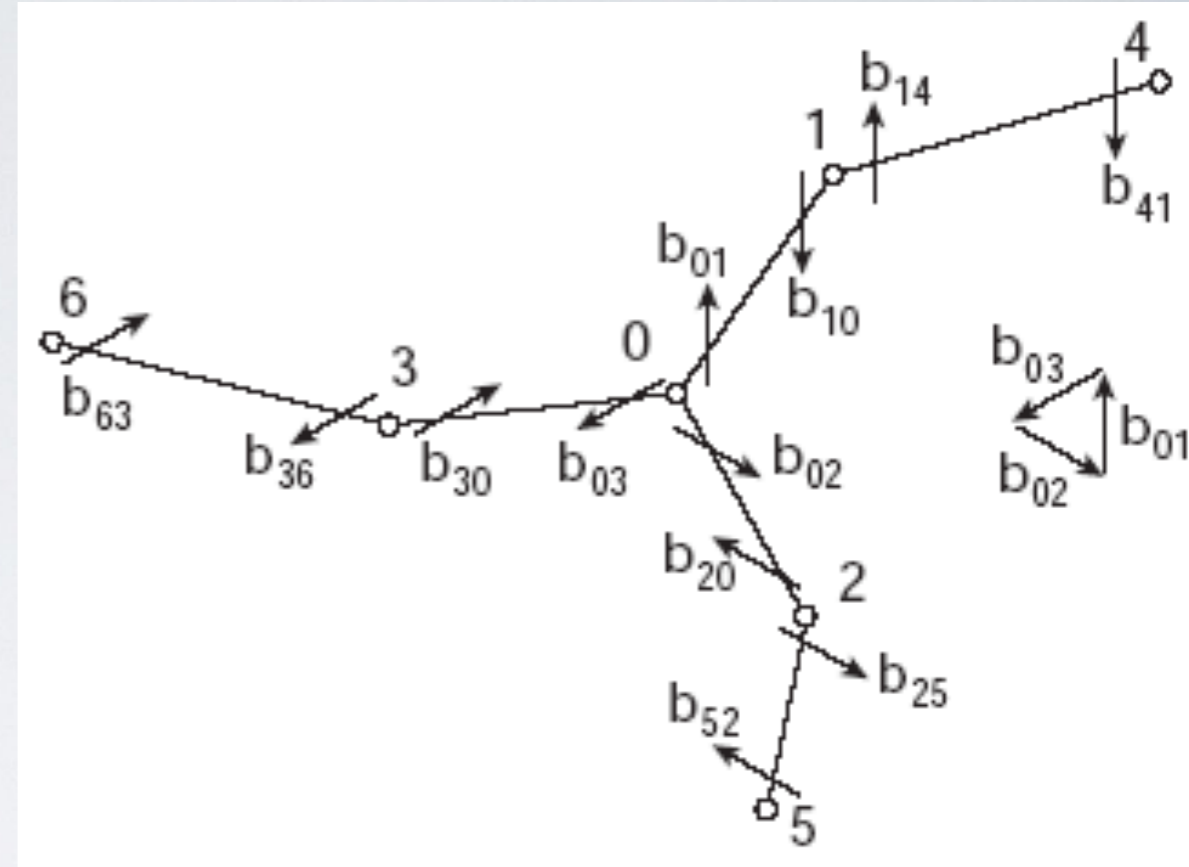
III. Basic Principles

The fundamental idea

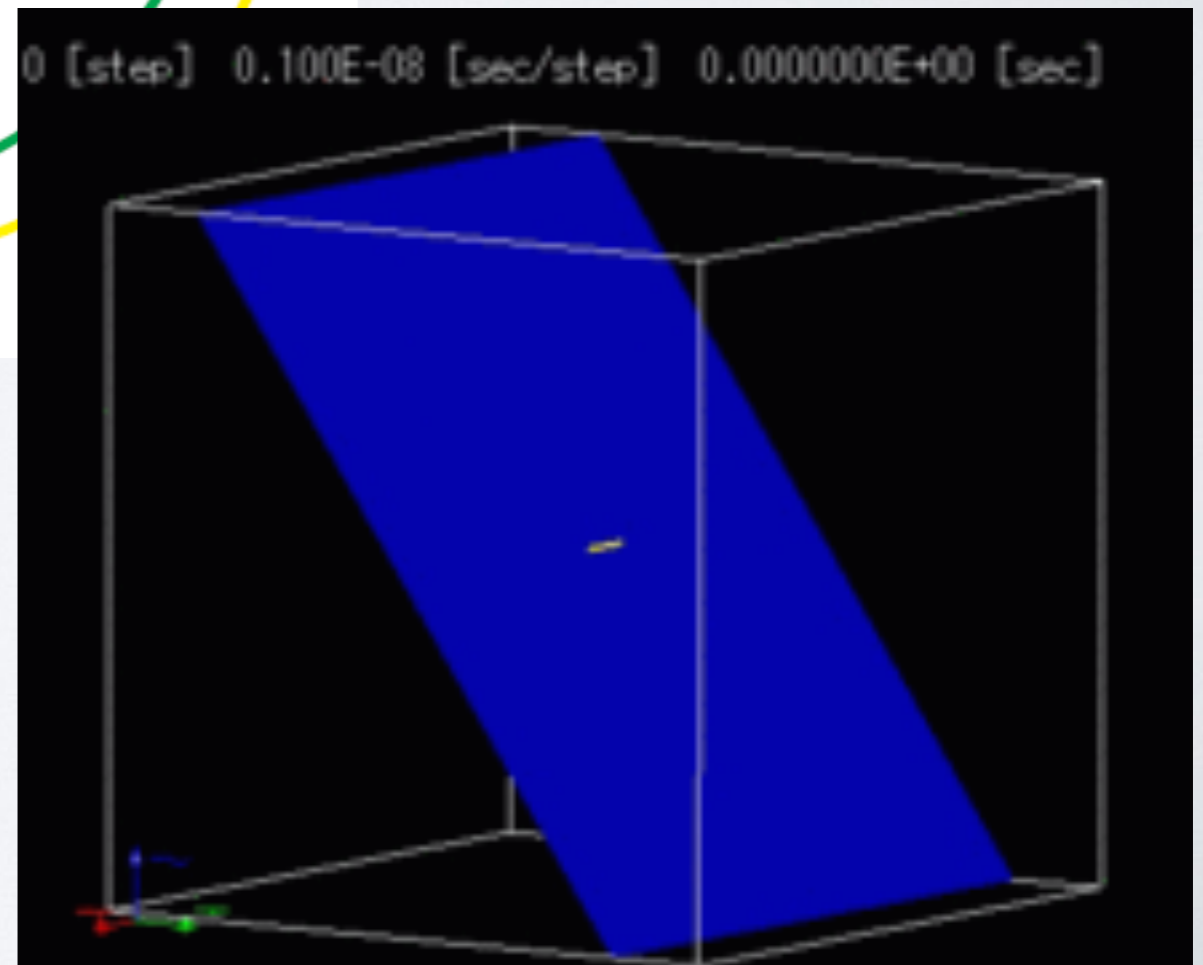
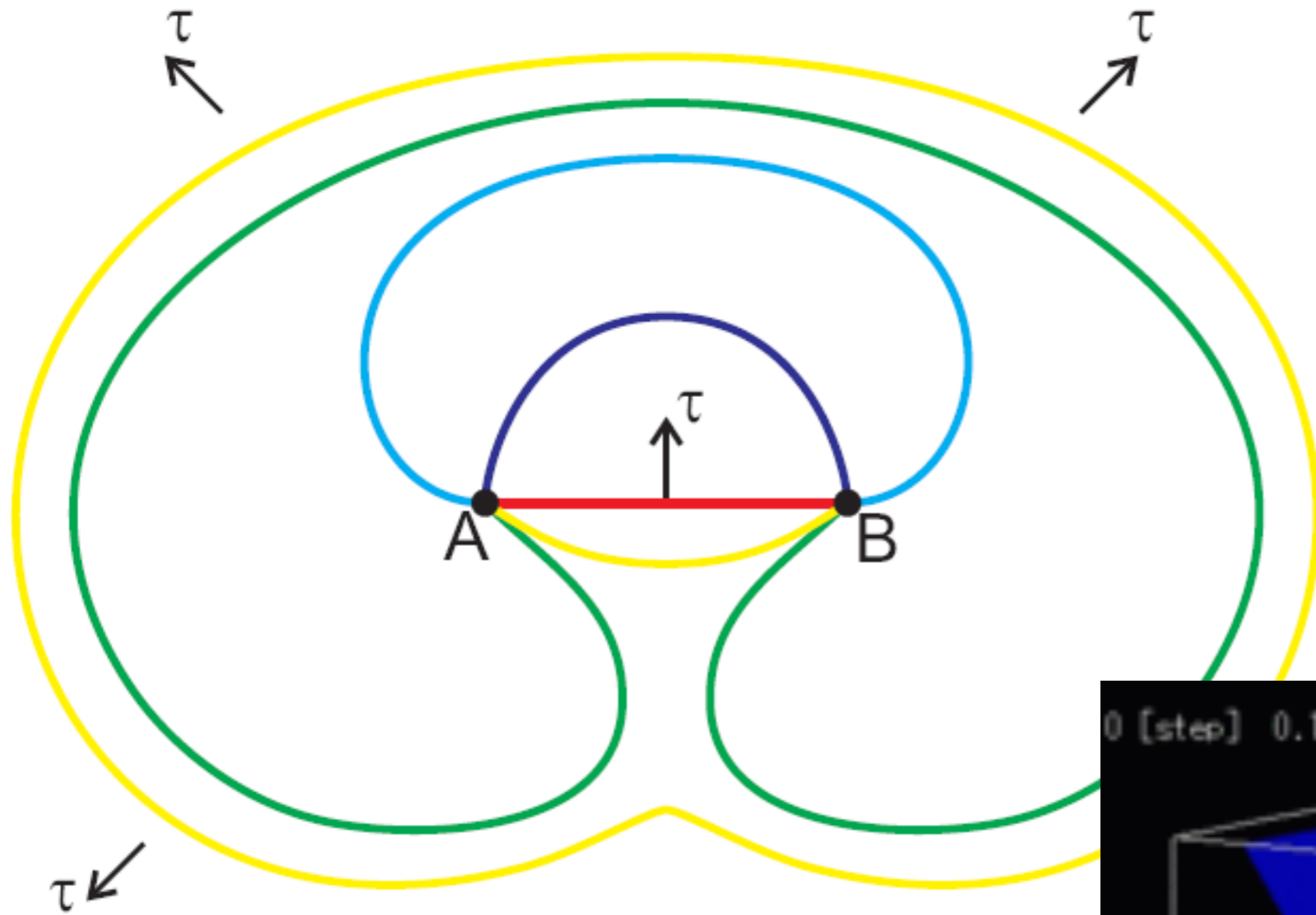
- DDD simulates dislocation motion using overdamped dynamics: similar to MD, but with a *mobility law* not mass.
- Running a simulation is like cooking - just follow the recipe!
- Three ingredients:
 1. An initial system configuration: dislocations
 2. Interaction between dislocations: Peach-Koehler
 3. Evolving dislocation geometry: mobility + remeshing

Ingredient 1: Initial configuration

- Represent dislocation curves
- Unlike MD, dislocations are one-dimensional objects: hence, *discretization*
- Different approaches:
 - Dislocation as line segments
 - Dislocation as cubic spline
 - Dislocation as connected arcs
 - Dislocation as “rastered” pure edge and screw segments
- Dislocation intersections have to be considered as well
- Typically use **Frank-Read** or **one-armed sources**



Ingredient I: Frank-Read source

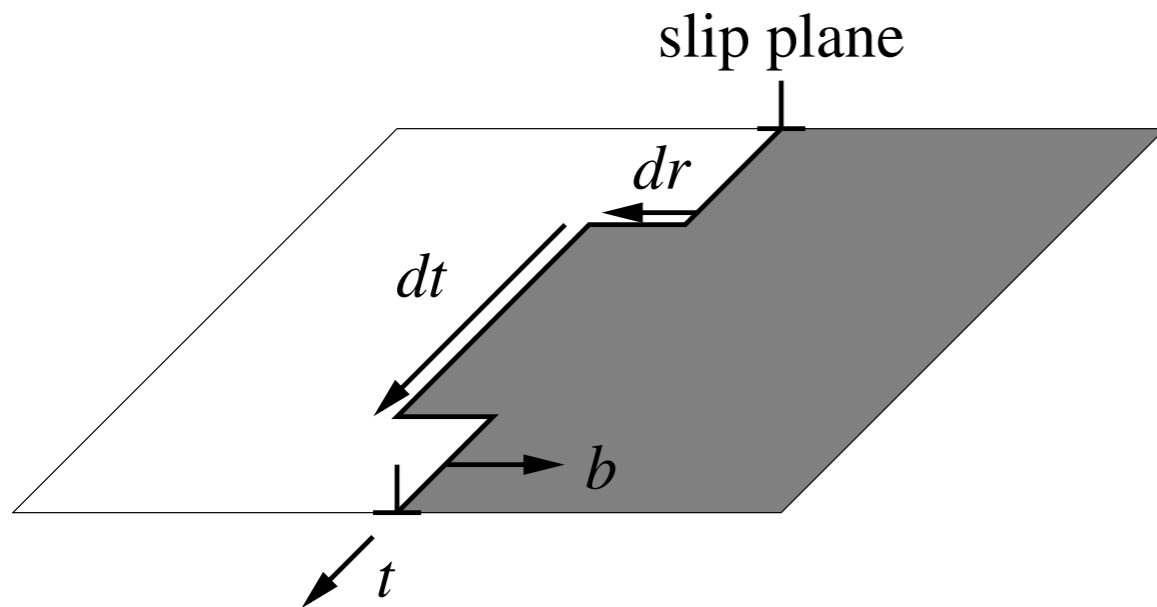


Ingredient 2: Interaction forces

- The net force acting on each dislocation in the system is a result of its interactions with all other dislocations + applied stress
- The underlying physical form is well-established: Peach-Koehler force

Dislocation motion under stress

- Dislocation line: separates “slipped” from “unslipped” parts of crystal
 - Sweeping out area displaces top part of crystal
 - Force on top area times displacement = –work done on dislocation



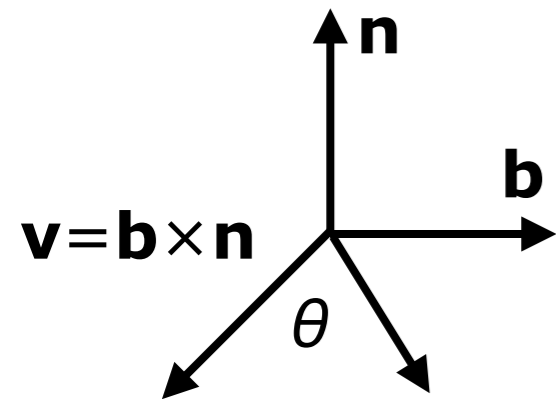
$$\mathbf{dF} = (\underline{\sigma} \cdot \mathbf{b}) \times \mathbf{dt}$$

Force/length on a dislocation

$$\begin{aligned} dW &= -(\text{force}) \cdot (\text{displacement}) \\ &= -(\text{stress} \cdot \text{area}) \cdot (\text{displacement}) \\ &= (\underline{\sigma} \cdot \mathbf{da}) \cdot \mathbf{b} = \mathbf{da} \cdot \underline{\sigma} \cdot \mathbf{b} \\ &= (\mathbf{dr} \times \mathbf{dt}) \cdot \underline{\sigma} \cdot \mathbf{b} \\ &= -(\mathbf{dt} \times \mathbf{dr}) \cdot (\underline{\sigma} \cdot \mathbf{b}) \\ &= -\left[(\underline{\sigma} \cdot \mathbf{b}) \times \mathbf{dt} \right] \cdot \mathbf{dr} \\ &= -\mathbf{dF} \cdot \mathbf{dr} \end{aligned}$$

Dislocation motion under stress

- Force per length (“Peach-Kohler force”)
 - Always perpendicular to dislocation line
 - Force **in slip plane**: *glide force*
 - Force **normal to slip plane**: *climb force* (edge) or *cross-slip* (screw)



$$\mathbf{t} = \mathbf{v} \cos \theta + \mathbf{b} \sin \theta$$

glide force

cross-slip force

climb force

$$d\mathbf{F} = (\underline{\sigma} \cdot \mathbf{b}) \times d\mathbf{t}$$

$$= \begin{pmatrix} \sigma_{vv} & \sigma_{vb} & \sigma_{vn} \\ \sigma_{vb} & \sigma_{bb} & \sigma_{bn} \\ \sigma_{vn} & \sigma_{bn} & \sigma_{nn} \end{pmatrix} \cdot \begin{pmatrix} 0 \\ b \\ 0 \end{pmatrix} \times \begin{pmatrix} \cos \theta \\ \sin \theta \\ 0 \end{pmatrix} dt$$

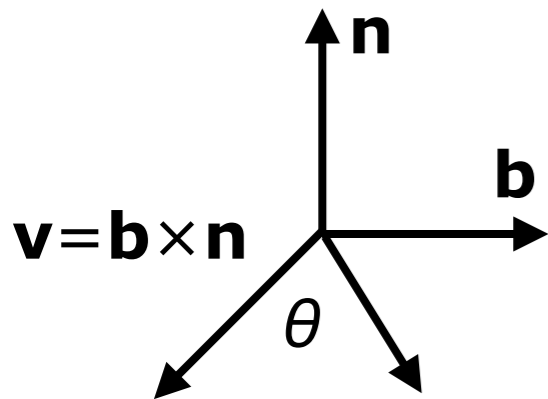
$$= \begin{pmatrix} \sigma_{bv}b \\ \sigma_{bb}b \\ \sigma_{bn}b \end{pmatrix} \times \begin{pmatrix} \cos \theta \\ \sin \theta \\ 0 \end{pmatrix} dt$$

$$= \begin{vmatrix} \mathbf{v} & \mathbf{b} & \mathbf{n} \\ \sigma_{bv}b & \sigma_{bb}b & \sigma_{bn}b \\ \cos \theta & \sin \theta & 0 \end{vmatrix} dt$$

$$= \sigma_{bn}b(-\mathbf{v} \sin \theta + \mathbf{b} \cos \theta) + \sigma_{bv}b \sin \theta \mathbf{n} - \sigma_{bb}b \cos \theta \mathbf{n}$$

Force on an edge and screw dislocation

- Force per length (“Peach-Kohler force”)
 - Always perpendicular to dislocation line
 - Force **in slip plane**: *glide force*
 - Force **normal to slip plane**: *climb force* (edge) or *cross-slip* (screw)



$$\mathbf{t} = \mathbf{v} \cos \theta + \mathbf{b} \sin \theta$$

$$\mathbf{dF} = \sigma_{bn} b (-\mathbf{v} \sin \theta + \mathbf{b} \cos \theta) + \sigma_{bv} b \sin \theta \mathbf{n} - \sigma_{bb} b \cos \theta \mathbf{n}$$

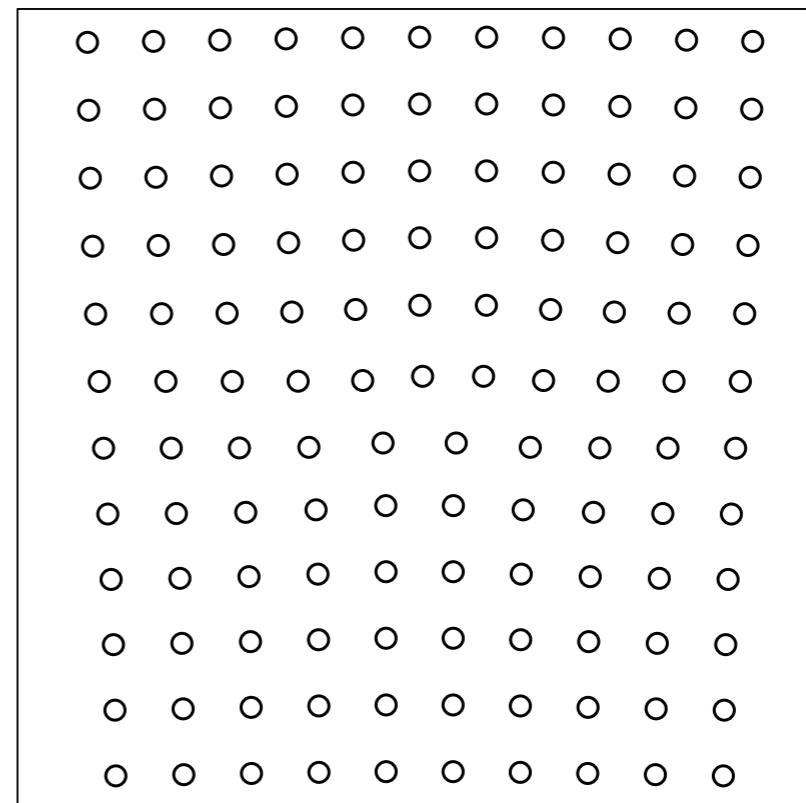
$$\mathbf{dF}_{\text{edge}} = \sigma_{bn} b \mathbf{b} - \sigma_{bb} b \mathbf{n}$$

$$\mathbf{dF}_{\text{screw}} = -\sigma_{bn} b \mathbf{v} + \sigma_{bv} b \mathbf{n}$$

glide force

cross-slip force

climb force



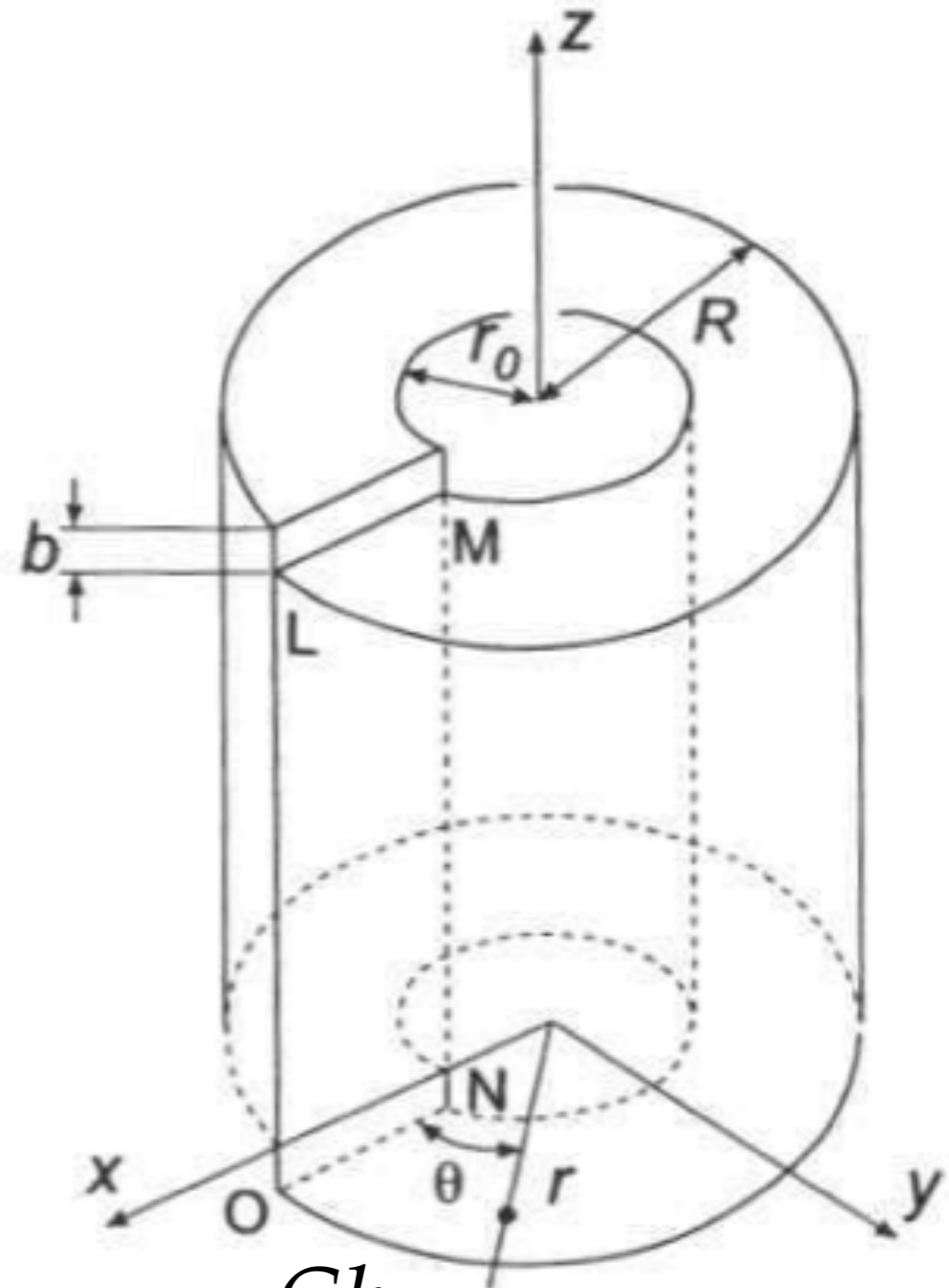
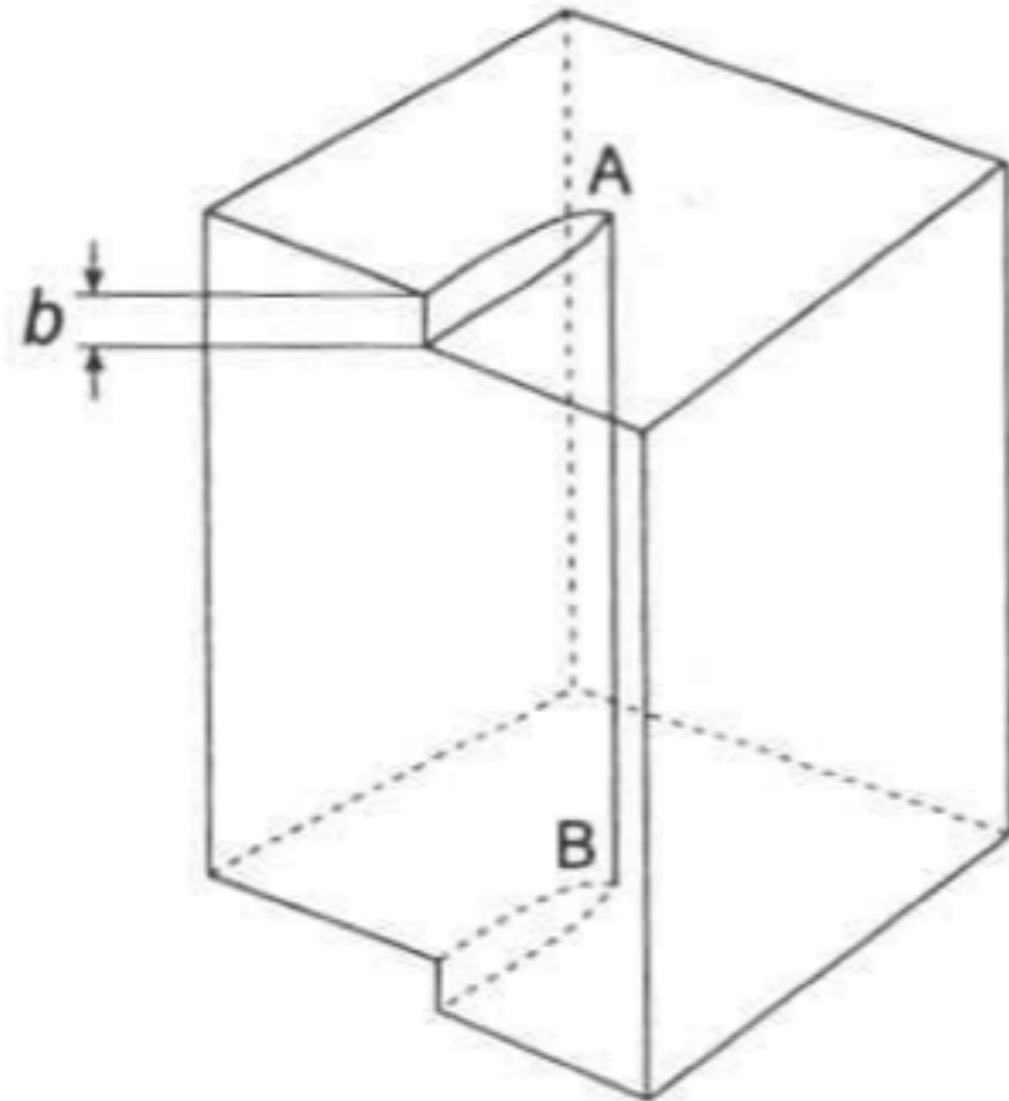
Stress and force

- To use, we need to need to be able to calculate stress at each point, *and* translate force on a dislocation segment onto the dislocation “degrees of freedom” (nodes)
- For elastically isotropic materials (with a core “cutoff”)

$$\sigma_{\alpha\beta}(\mathbf{x}) = \frac{\mu}{8\pi} \oint_C \partial_i \partial_p \partial_p R_a \left[b_m \epsilon_{im\alpha} dx'_\beta + b_m \epsilon_{im\beta} dx'_\alpha \right] \quad R = \|\mathbf{x} - \mathbf{x}'\|$$

$$+ \frac{\mu}{4\pi(1-\nu)} \oint_C b_m \epsilon_{ink} \left[\partial_i \partial_\alpha \partial_\beta R_a - \delta_{\alpha\beta} \partial_i \partial_p \partial_p R_a \right] dx'_k \quad R_a = \sqrt{\|\mathbf{x} - \mathbf{x}'\|^2 + a^2}$$

Screw dislocation: stress field



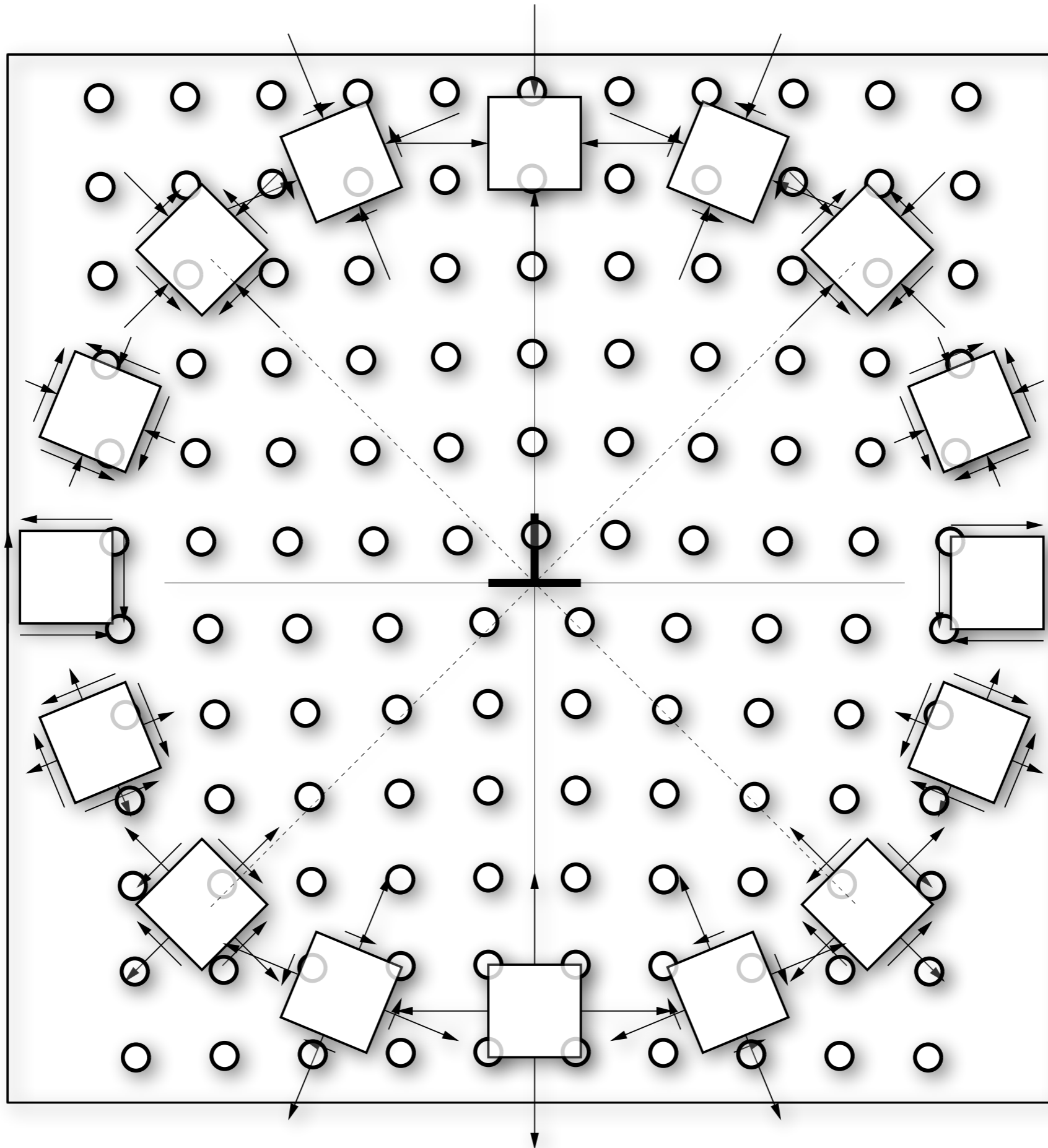
$$\gamma = \frac{b}{2\pi r}$$

$$\sigma_{z\theta} = \sigma_{\theta z} = \frac{Gb}{2\pi r}$$

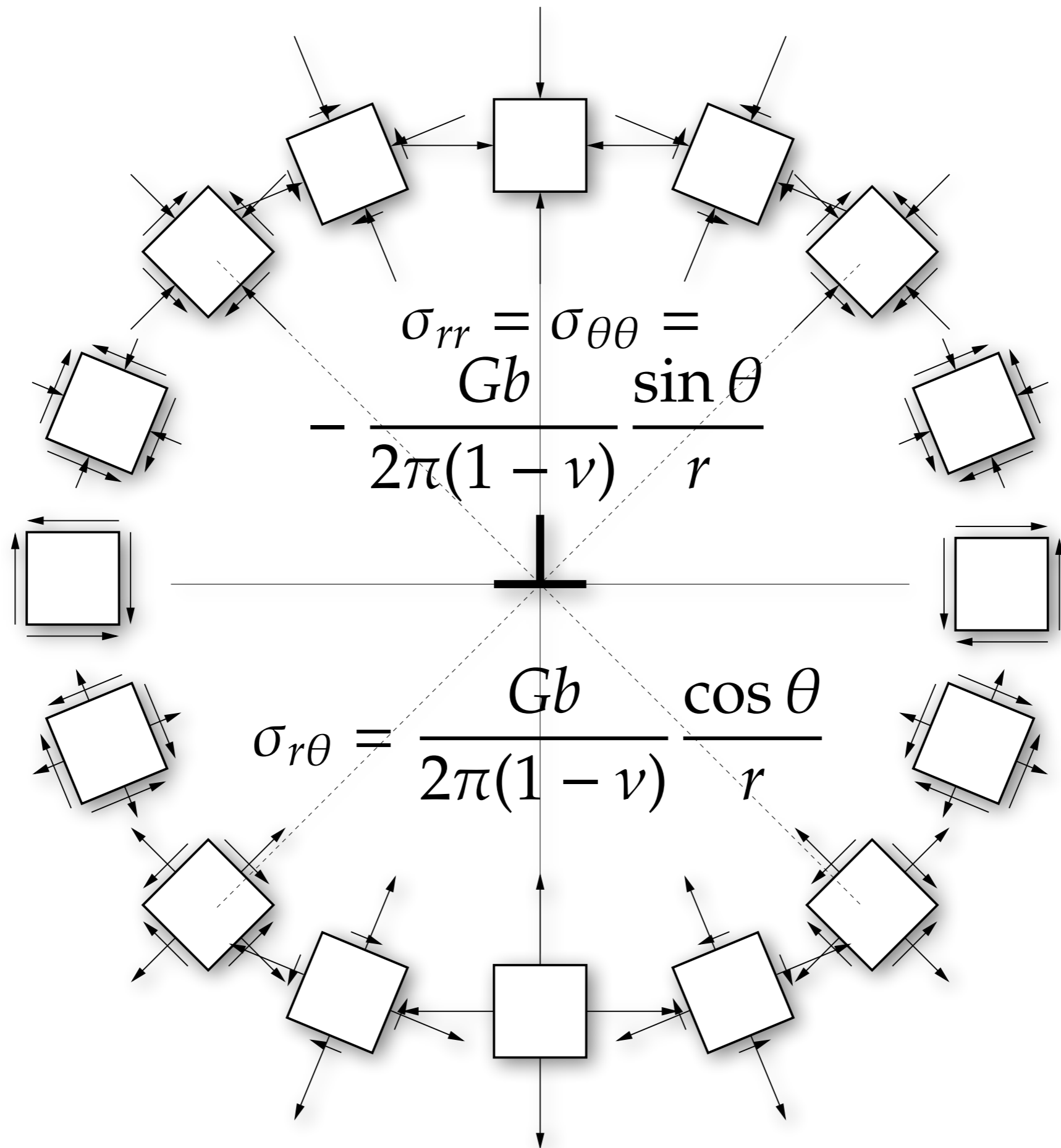
$$\sigma_{xz} = \sigma_{zx} = -\frac{Gb}{2\pi r} \sin \theta$$

$$\sigma_{yz} = \sigma_{zy} = \frac{Gb}{2\pi r} \cos \theta$$

Edge dislocation: stress field



Edge dislocation: stress field



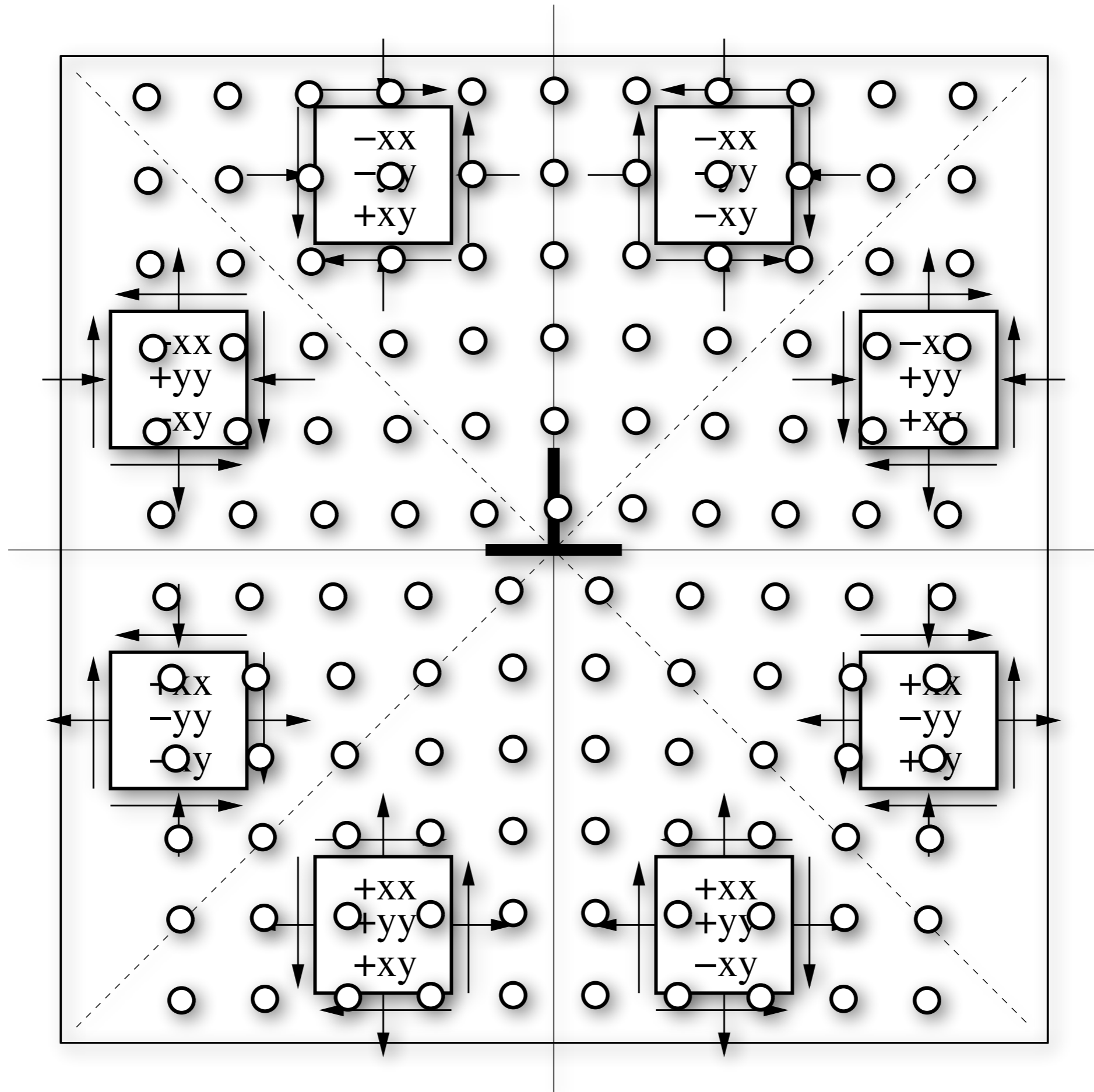
Edge dislocation: stress field

$$\sigma_{rr} = \sigma_{\theta\theta} = \frac{Gb}{2\pi(1-\nu)} \frac{\sin \theta}{r} \quad \sigma_{r\theta} = \frac{Gb}{2\pi(1-\nu)} \frac{\cos \theta}{r}$$

$$\theta_{\text{cart,polar}} = \begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix}$$

$$\begin{aligned} \begin{pmatrix} \sigma_{xx} & \sigma_{xy} \\ \sigma_{xy} & \sigma_{yy} \end{pmatrix} &= \begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} \sigma_{rr} & \sigma_{r\theta} \\ \sigma_{r\theta} & \sigma_{\theta\theta} \end{pmatrix} \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix} \\ &= \frac{Gb}{2\pi(1-\nu)r} \begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} -\sin \theta & \cos \theta \\ \cos \theta & -\sin \theta \end{pmatrix} \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix} \\ &= \frac{Gb}{2\pi(1-\nu)r} \begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} -2\sin \theta \cos \theta & -\sin^2 \theta + \cos^2 \theta \\ 1 & 0 \end{pmatrix} \\ &= \frac{Gb}{2\pi(1-\nu)r} \begin{pmatrix} -\sin \theta(1 + \cos^2 \theta) & \cos \theta \cos 2\theta \\ \cos \theta \cos 2\theta & -\sin \theta \cos 2\theta \end{pmatrix} \\ &= \frac{Gb}{2\pi(1-\nu)r} \begin{pmatrix} -\sin \theta(2 + \cos 2\theta) & \cos \theta \cos 2\theta \\ \cos \theta \cos 2\theta & -\sin \theta \cos 2\theta \end{pmatrix} \end{aligned}$$

Edge dislocation: stress field



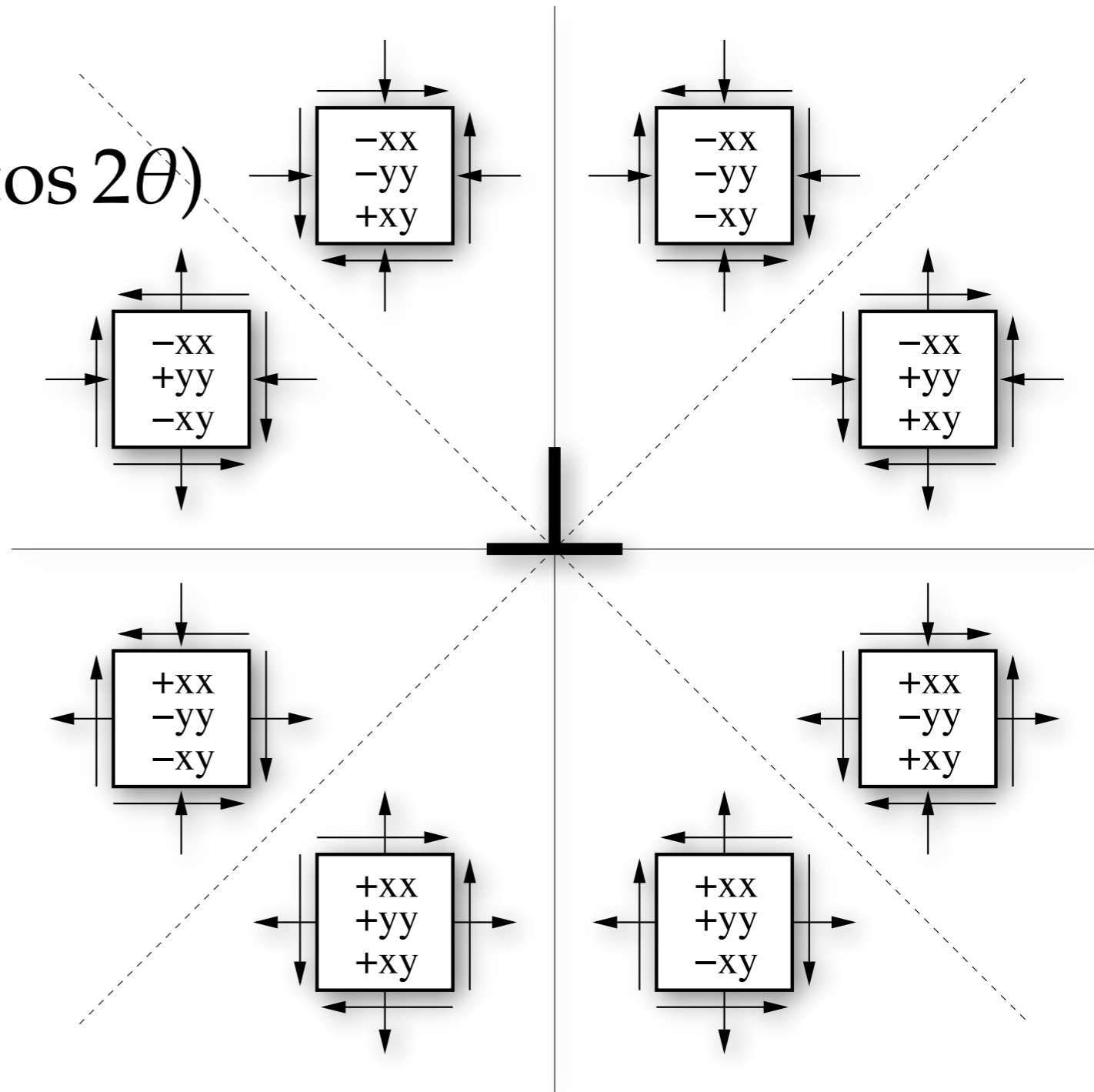
Edge dislocation: stress field

$$\sigma_{xx} = -\frac{Gb}{2\pi r(1-\nu)} \sin \theta (2 + \cos 2\theta)$$

$$\sigma_{yy} = \frac{Gb}{2\pi r(1-\nu)} \sin \theta \cos 2\theta$$

$$\sigma_{zz} = \nu(\sigma_{xx} + \sigma_{yy})$$

$$\tau_{xy} = \frac{Gb}{2\pi r(1-\nu)} \cos \theta \cos 2\theta$$



Stress and force

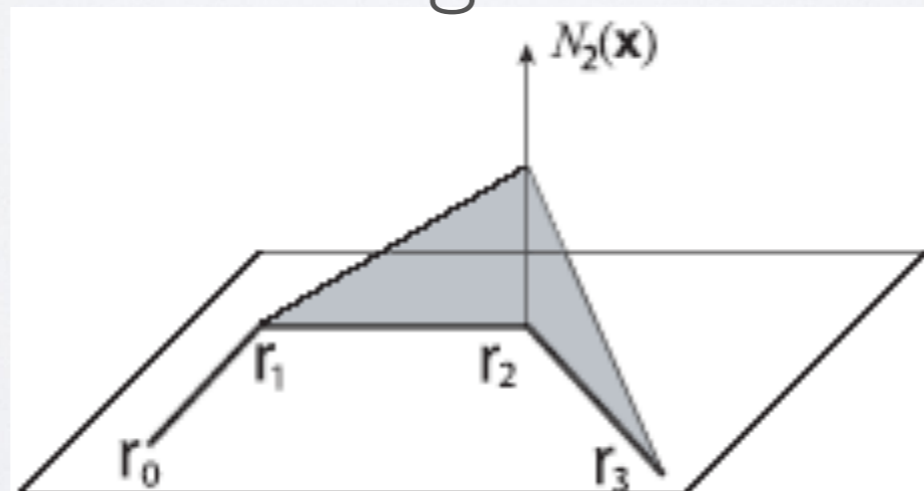
- To use, we need to be able to calculate stress at each point, *and* translate force on a dislocation segment onto the dislocation “degrees of freedom” (nodes)
- For elastically isotropic materials (with a core “cutoff”)

$$\sigma_{\alpha\beta}(\mathbf{x}) = \frac{\mu}{8\pi} \oint_C \partial_i \partial_p \partial_p R_a \left[b_m \epsilon_{im\alpha} dx'_\beta + b_m \epsilon_{im\beta} dx'_\alpha \right] \quad R = \|\mathbf{x} - \mathbf{x}'\|$$

$$R_a = \sqrt{\|\mathbf{x} - \mathbf{x}'\|^2 + a^2}$$

$$+ \frac{\mu}{4\pi(1-\nu)} \oint_C b_m \epsilon_{ink} \left[\partial_i \partial_\alpha \partial_\beta R_a - \delta_{\alpha\beta} \partial_i \partial_p \partial_p R_a \right] dx'_k$$

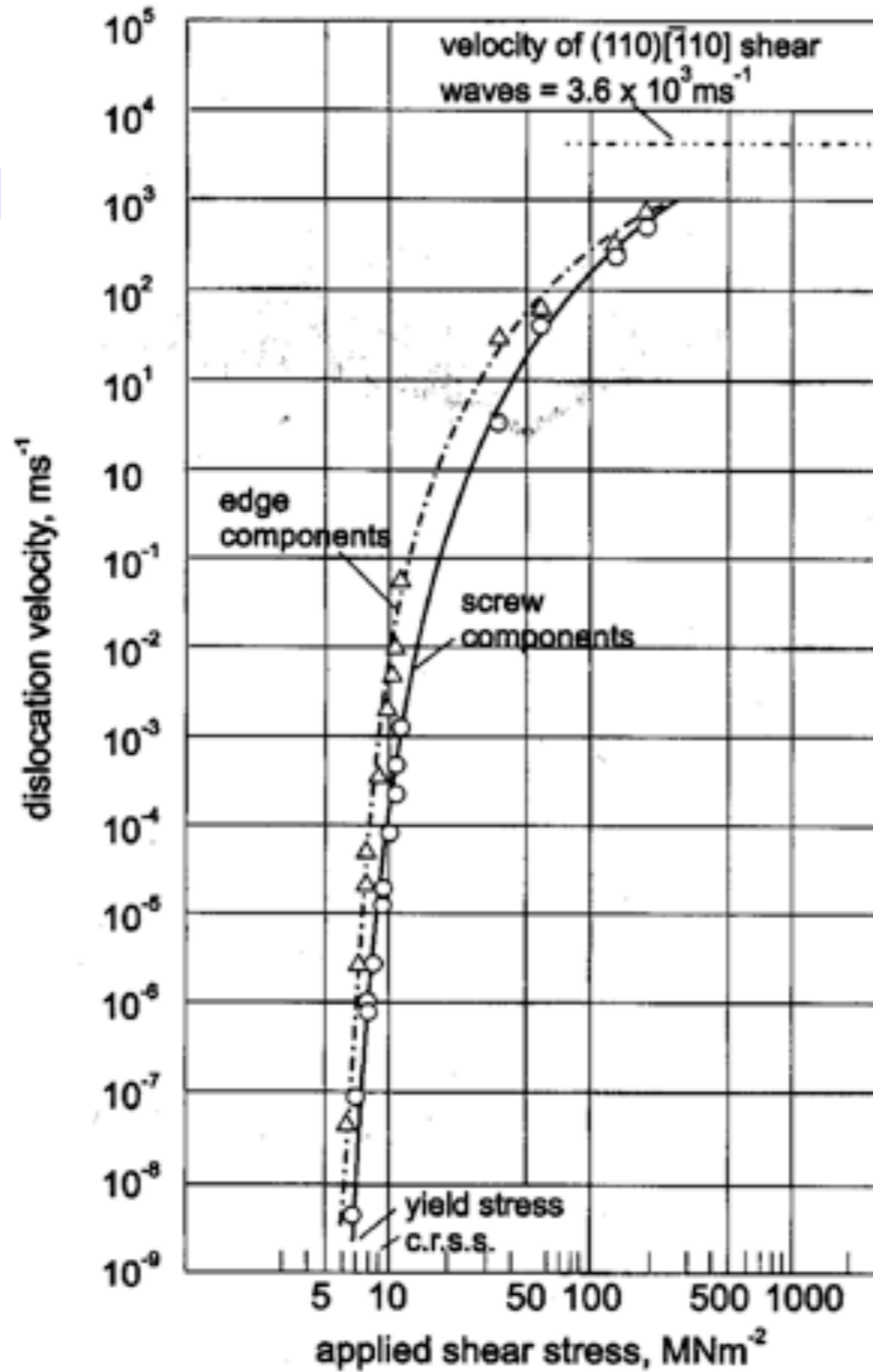
- For line segments, map force to nodes by integrating a linear “shape function” along connected segments



Ingredient 3: Evolution of dislocations

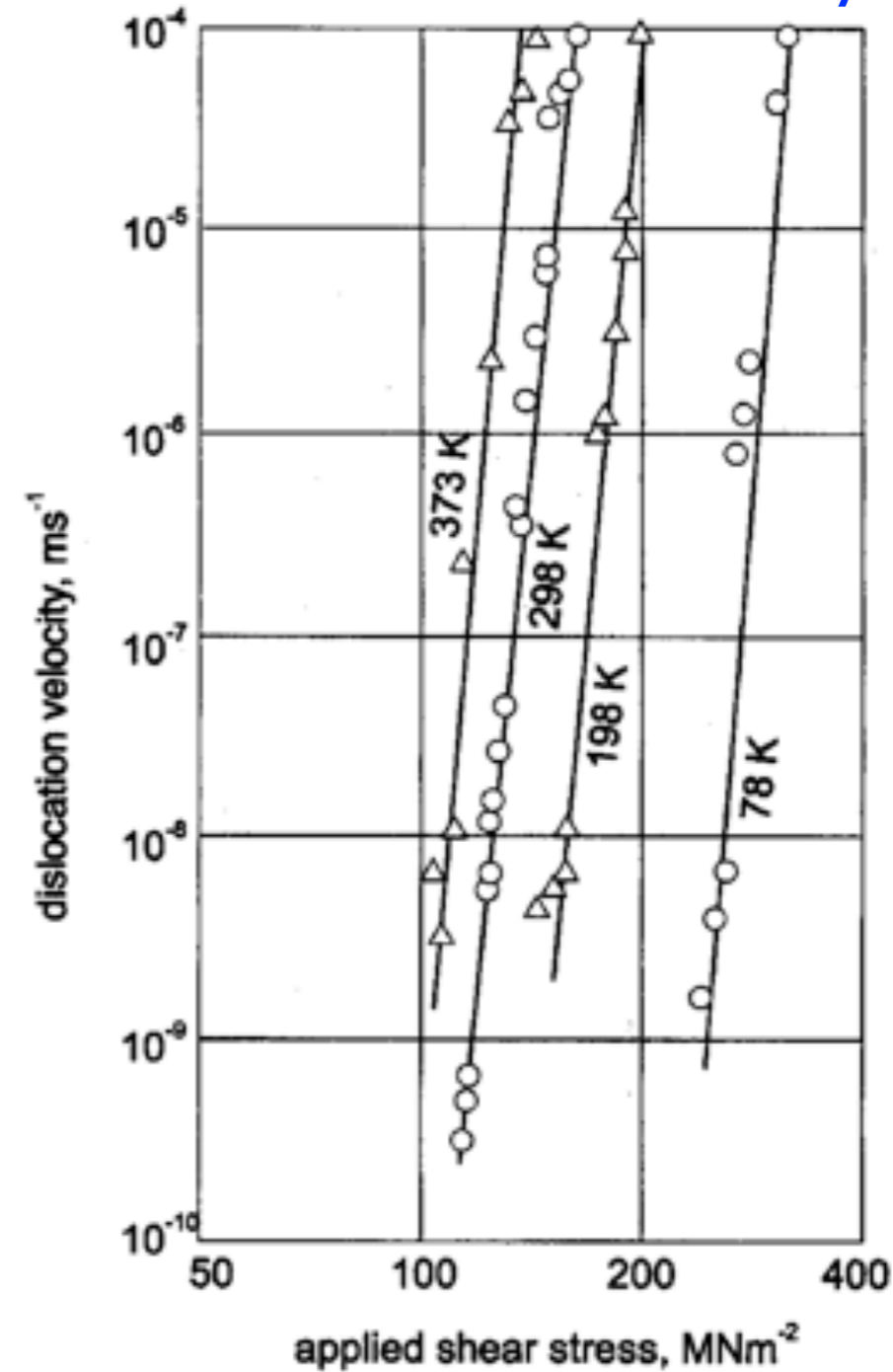
- Initial dislocation geometry + forces on dislocation:
How to evolve? $F=ma$?
- Dislocations aren't “objects” with mass; instead, they experience an *overdamped dynamics* where **velocity** is proportional to **force**: mobility.

LiF
crystal



(a)

Fe-3.25%Si
crystal



(b)

Figure 3.11 (a) Stress dependence of the velocity of edge and screw dislocations in lithium fluoride. (From Johnston and Gilman, *J. Appl. Phys.* **30**, 129, 1959.) (b) Stress dependence of the velocity of edge dislocations in 3.25 per cent silicon iron at four temperatures. (After Stein and Low, *J. Appl. Phys.* **31**, 362, 1960.)

Ingredient 3: Evolution of dislocations

- Initial dislocation geometry + forces on dislocation:
How to evolve? $F=ma$?
- Dislocations aren't “objects” with mass; instead, they experience an *overdamped dynamics* where **velocity** is proportional to **force**: mobility.
- Mobility can be *anisotropic*
 - In plane: edge dislocations vs. screw dislocations
 - Out-of-plane: cross-slip (screw) vs. climb (edge)

Ingredient 3: Evolution of dislocations

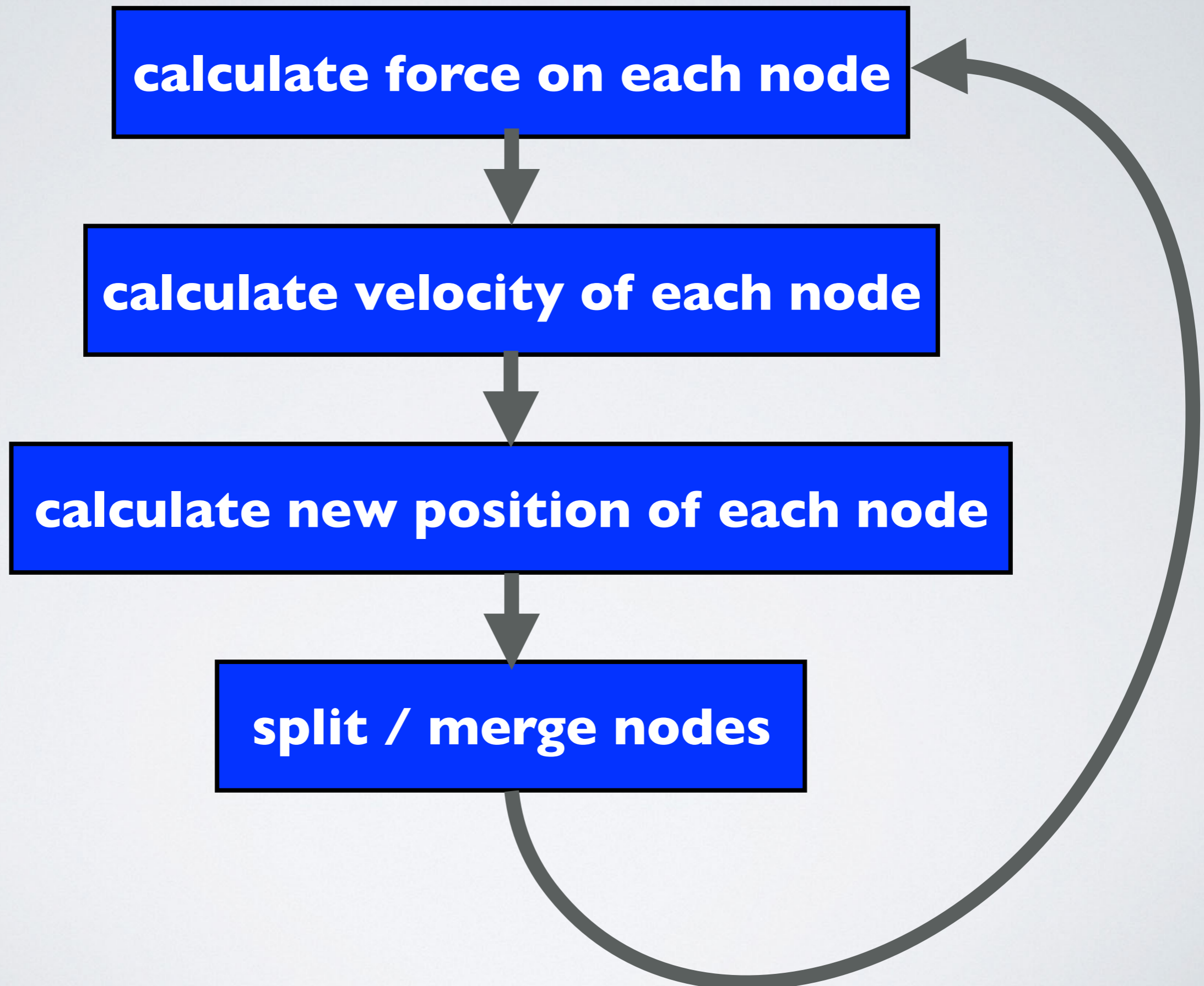
- Given velocity on nodes, integrate nodes forward in time
- Euler forward / backward schemes

$$\text{Forward: } \mathbf{x}(t + \Delta t) = \mathbf{x}(t) + \Delta t \dot{\mathbf{x}}(t)$$

$$\text{Backward: } \mathbf{x}(t + \Delta t) = \mathbf{x}(t) + \Delta t \dot{\mathbf{x}}(t + \Delta t)$$

- Newton-Krylov, other sophisticated integrators
- Remeshing:
 - Too few nodes to represent the curvature?
Introduce new nodes (**split**)
 - Too many nodes in a small area of space?
Eliminate extra nodes (**merge**)
- Dislocation annihilation reactions: elimination of dislocation segments with opposite Burgers vector

Dislocation dynamics flowchart



IV. Advanced Topics

Boundary conditions

- Interaction between dislocations (calculated via stress) depends on the boundary conditions:
 - Most simulations assume *periodic boundary conditions*
 - Simulations of finite systems involves *image stresses*
- Interaction calculation can be accelerated with Fast Multipole Methods (similar to charged interactions)

V. Dislocation Dynamics Packages

DDD software



<http://paradis.stanford.edu/>