

- Diffusion capacitance.

The concentration of excess carriers diffusing in the quasi-neutral regions can be obtained from Eq. (148):

$$\begin{cases} \delta p_n(x) = p_n(x) - p_{n0} = p_{n0} \left[ e^{eV_a/(k_B T)} - 1 \right] e^{-(x-x_{n0})/L_p} & (x > x_{n0}) \\ \delta n_p(x) = n_p(x) - n_{p0} = n_{p0} \left[ e^{eV_a/(k_B T)} - 1 \right] e^{(x+x_{p0})/L_n} & (x < -x_{p0}) \end{cases} \quad (185)$$

The charge per unit area will be (considering only holes, a similar expression will hold for electrons):

$$Q_{diff,p} = e \int_{x_{n0}}^{\infty} \delta p_n(x) dx = e L_p p_{n0} \left[ e^{eV_a/(k_B T)} - 1 \right] \quad (186)$$

Under strong forward bias,  $eV_a/(k_B T) \gg 1$ , so:

$$C_{diff,p} \approx \frac{dQ_{diff,p}}{dV_a} = \frac{e^2 L_p p_{n0}}{k_B T} e^{eV_a/(k_B T)} \approx \frac{e L_p^2}{k_B T} J_p(x_{n0}) = \frac{e}{k_B T} \tau_p J_p(x_{n0}), \quad (187)$$

where we have used Eq. (150) in the last step. Accounting now for the charge of the minority electrons diffusing into the  $p$  quasi-neutral region:

$$C_{diff} \approx \frac{dQ_{diff,p}}{dV_a} + \frac{dQ_{diff,n}}{dV_a} = \frac{e}{k_B T} [\tau_p J_p(x_{n0}) + \tau_n J_n(-x_{p0})] \quad (188)$$

Note: This is the equation found in most textbooks. However, in a text by Karl Hess (and in a brief comment on the latest edition of the Streetman-Banerjee's book) one finds instead:

$$C_{diff} \approx \frac{e}{2k_B T} [\tau_p J_p(x_{n0}) + \tau_n J_n(-x_{p0})], \quad (189)$$

where the additional factor of 1/2 is explicitly commented and the claim is made that Eq. (188) is in error. This factor can be justified in hand-waving fashion by noting that the charges  $Q_{diff,p}$  and  $Q_{diff,n}$  are like the charges in the two opposite plates of a capacitor, so that the capacitance should be given by the change w.r.t. the applied bias of the average of the electron and hole charges. A more sophisticated and rigorous explanation is given by S. E. Laux and K. Hess, IEEE Trans. Electron. Device vol. 46, no. 2 (February 1999), p. 396. Their argument is based on the observation that – rigorously speaking – the diffusion charge extends also inside the depletion region, so that the integration in Eq. (186) should extend from 0 to  $\infty$ , not from  $x_{n0}$  (and similarly for the expression for  $Q_{diff,n}$ ). Since as  $V_a$  changes charges will leave the depletion region, we will obtain a lower estimate for the charge, and so for the capacitance. In a way, this argument is equivalent to our ‘hand-waving’ argument since both reduce to accounting for the charges throughout the entire junction, not just in the quasi-neutral regions.

The junction capacitance dominates the reactance of a p-n junction under reverse bias; for forward bias, however, the charge storage capacitance  $C_s$  becomes dominant. To calculate the capacitance due to charge storage effects,

let us assume that a p<sup>+</sup>-n junction is forward biased with a steady current  $I$ . The stored charge in the injected hole distribution is:

$$Q_p = I\tau_p = qA \Delta p L_p = qAL_p p_0 e^{eV/kT} \quad \text{for } V \gg 0.0259 \text{ V} \quad (5-64)$$

The capacitance due to small changes in this stored charge is

$$C_s = \frac{dQ_p}{dV} = \frac{q^2}{kT} AL_p p_0 \frac{d(e^{eV/kT})}{dV} = \frac{q}{kT} I\tau_p \quad (5-65)$$

Similarly, we can determine the a-c conductance by allowing small changes in the current:

$$G_s = \frac{dI}{dV} = \frac{qAL_p p_0}{\tau_p} \frac{d(e^{eV/kT})}{dV} = \frac{q}{kT} I \quad (5-66)$$

Thus the a-c component of current is

$$i(a-c) = G_s v(a-c) + C_s \frac{dv(a-c)}{dt} \quad (5-67)$$

where

$$G_s = \frac{q}{kT} I(d-c) \quad \text{and} \quad C_s = G_s \tau_p$$

The charge storage capacitance can be a serious limitation for forward-biased p-n junctions in high-frequency circuits. As in the case of the switching performance discussed in the two preceding sections, the high-frequency a-c response of a junction can be improved by reducing the carrier lifetime. Since  $C_s$  is proportional to  $\tau_p$ , a short hole lifetime can make the forward-biased capacitance of a p<sup>+</sup>-n junction acceptably small for many applications.

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